

# “Flavour in the era of the LHC”

(or, more honestly: a subset of highlights from LHCb)



**Workshop on High Energy Physics Phenomenology**

**12-21st December, 2013, Puri, India**

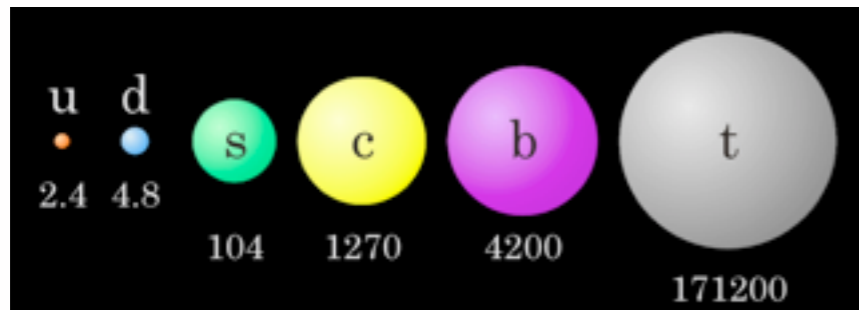
Wouter Hulsbergen (Nikhef, Amsterdam)  
on behalf of the LHCb collaboration

# Motivation

- SM explains fermion masses and quark flavour transitions via Yukawa couplings

$$-\mathcal{L}_{\text{Yukawa}}^{\text{SM}} = Y_d^{ij} \bar{Q}_L^i \phi D_R^j + Y_u^{ij} \bar{Q}_L^i \tilde{\phi} U_R^j + Y_e^{ij} \bar{L}_L^i \phi E_R^j + \text{h.c.}$$

quark (and lepton) masses



CKM matrix

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \sim \begin{pmatrix} \blacksquare & \blacksquare & \cdot \\ \blacksquare & \blacksquare & \blacksquare \\ \cdot & \blacksquare & \blacksquare \end{pmatrix}$$

- many parameters (10 out of 19 SM parameters are in quark sector!)
- rich phenomenology
  - ▶ meta-stable baryons with many decay modes, some very rare ...
  - ▶ neutral meson mixing, CP violation, ....

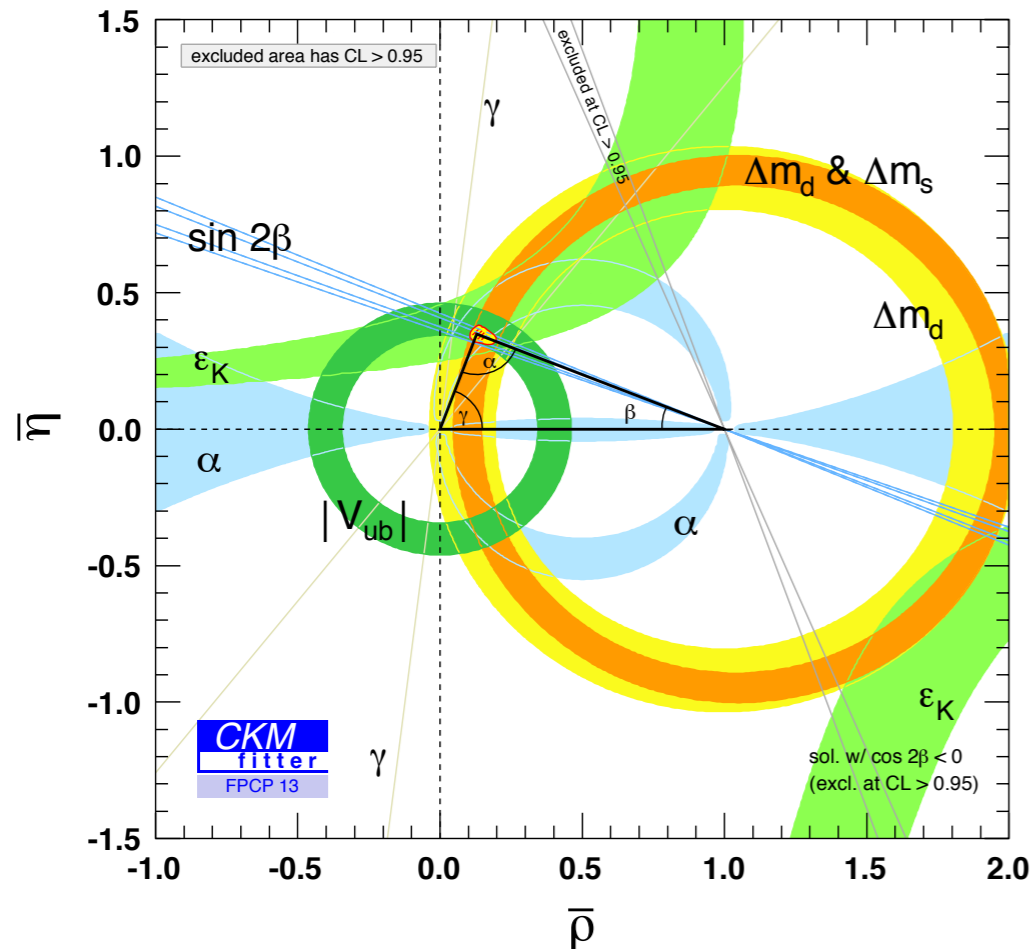


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- extremely successful ...



... but puzzling

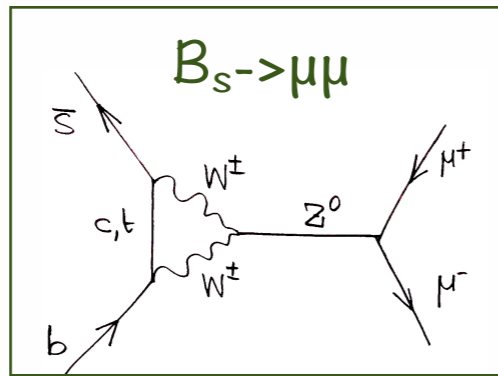
- why the large hierarchy in masses and couplings?
- if due to symmetry, then what breaks it?
- if there is NP at the TeV scale, then why don't we see it?

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

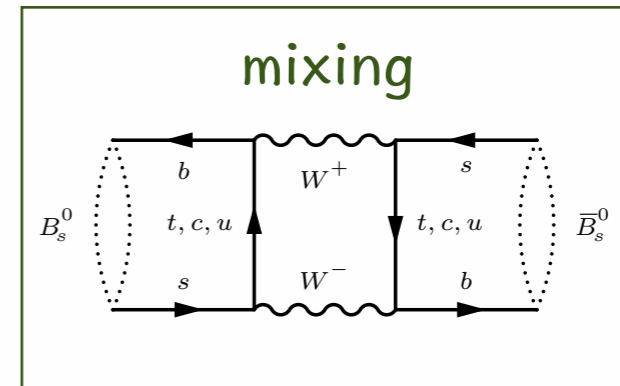
# The NP flavour problem

- flavour transitions probes high mass scales in quantum loops, e.g. in FCNC

$$\Delta F = 1$$



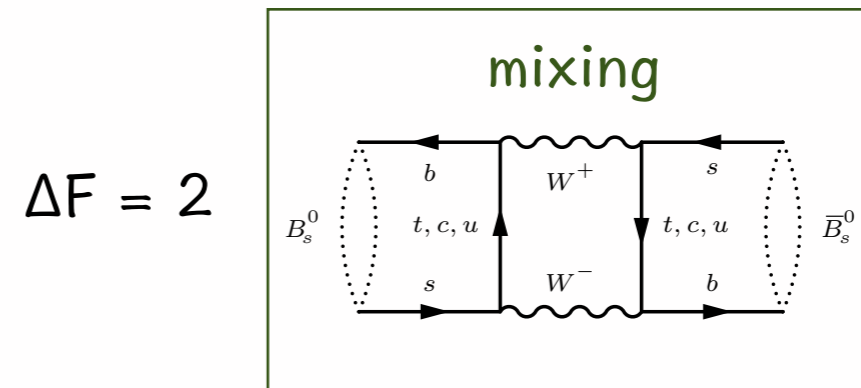
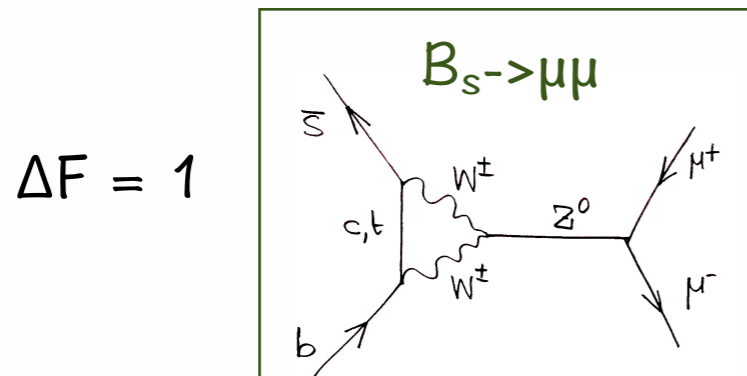
$$\Delta F = 2$$





# The NP flavour problem

- flavour transitions probes high mass scales in quantum loops, e.g. in FCNC



- parameterize in terms of operators, couplings and mass scales

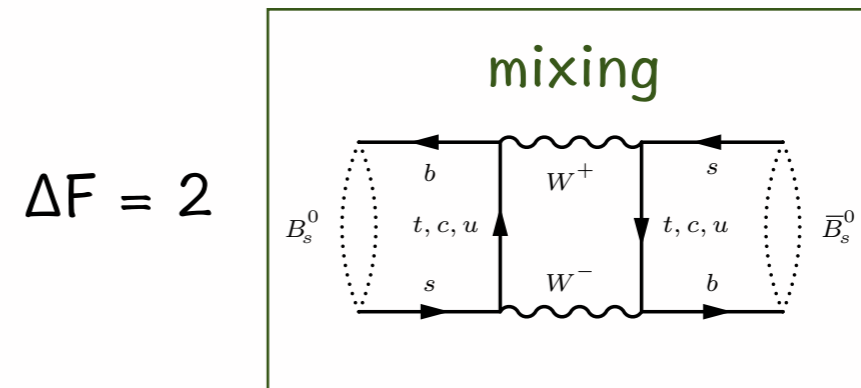
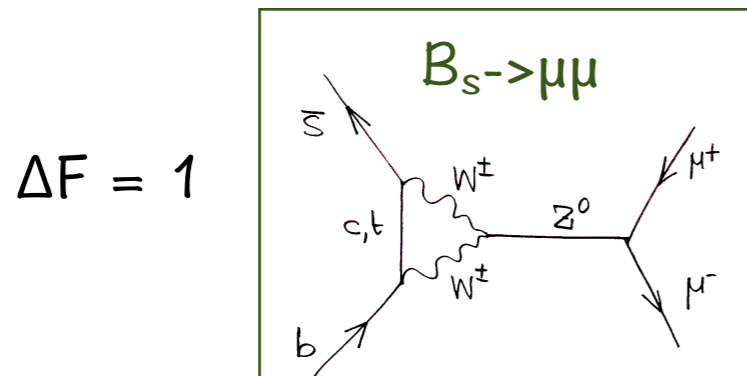
$$\mathcal{A}(f_i \rightarrow f_j + X) = \mathcal{A}_0 \left[ \frac{c_{SM}}{M_W^2} + \frac{c_{NP}}{\Lambda^2} \right]$$

Operator	Bounds on $\Lambda$ in TeV ( $c_{NP} = 1$ )		Bounds on $c_{NP}$ ( $\Lambda = 1$ TeV)		Observables
	Re	Im	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	$9.8 \times 10^2$	$1.6 \times 10^4$	$9.0 \times 10^{-7}$	$3.4 \times 10^{-9}$	$\Delta m_K; \epsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	$1.8 \times 10^4$	$3.2 \times 10^5$	$6.9 \times 10^{-9}$	$2.6 \times 10^{-11}$	$\Delta m_K; \epsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	$1.2 \times 10^3$	$2.9 \times 10^3$	$5.6 \times 10^{-7}$	$1.0 \times 10^{-7}$	$\Delta m_D;  q/p , \phi_D$
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	$6.2 \times 10^3$	$1.5 \times 10^4$	$5.7 \times 10^{-8}$	$1.1 \times 10^{-8}$	$\Delta m_D;  q/p , \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	$6.6 \times 10^2$	$9.3 \times 10^2$	$2.3 \times 10^{-6}$	$1.1 \times 10^{-6}$	$\Delta m_{B_d}; S_{\psi K_S}$
$(\bar{b}_R d_L)(\bar{b}_L d_R)$	$2.5 \times 10^3$	$3.6 \times 10^3$	$3.9 \times 10^{-7}$	$1.9 \times 10^{-7}$	$\Delta m_{B_d}; S_{\psi K_S}$
$(\bar{b}_L \gamma^\mu s_L)^2$	$1.4 \times 10^2$	$2.5 \times 10^2$	$5.0 \times 10^{-5}$	$1.7 \times 10^{-5}$	$\Delta m_{B_s}; S_{\psi\phi}$
$(\bar{b}_R s_L)(\bar{b}_L s_R)$	$4.8 \times 10^2$	$8.3 \times 10^2$	$8.8 \times 10^{-6}$	$2.9 \times 10^{-6}$	$\Delta m_{B_s}; S_{\psi\phi}$

**Table 1.1:** Bounds on representative dimension-six  $\Delta F = 2$  operators, assuming an effective coupling  $c_{NP}/\Lambda^2$ . The bounds are quoted on  $\Lambda$ , setting  $|c_{NP}| = 1$ , or on  $c_{NP}$ , setting  $\Lambda = 1$  TeV. The right column denotes the main observables used to derive these bounds (see next chapter for more details).

# The NP flavour problem

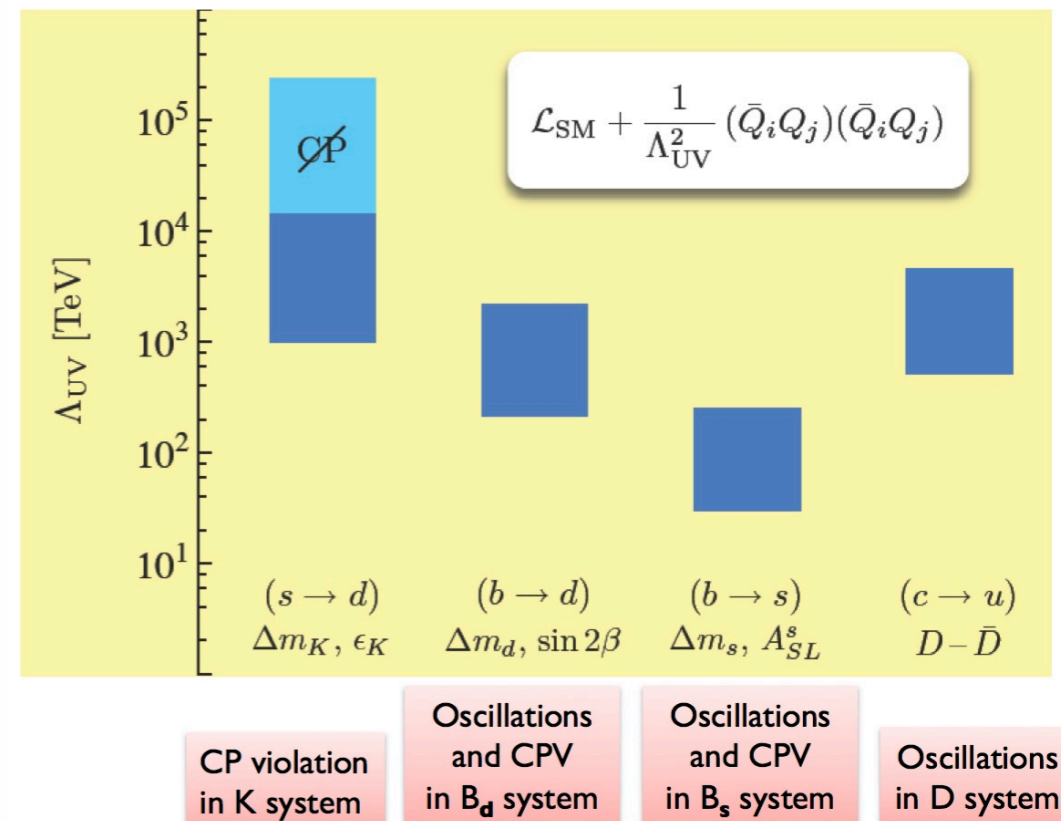
- flavour transitions probes high mass scales in quantum loops, e.g. in FCNC



- parameterize in terms of operators, couplings and mass scales

$$\mathcal{A}(f_i \rightarrow f_j + X) = \mathcal{A}_0 \left[ \frac{C_{SM}}{M_W^2} + \frac{C_{NP}}{\Lambda^2} \right]$$

- NP models with generic flavour structure at 1 TeV are ruled out
- either NP scale is much higher, or couplings are small. why?

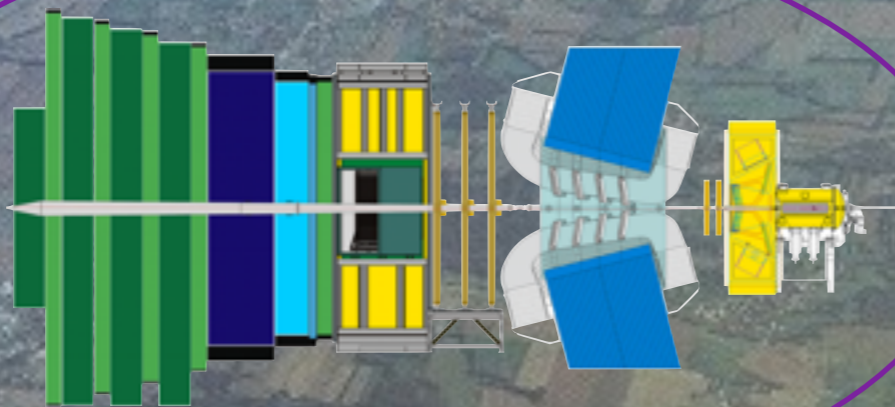


(Fig. Neubert, EPS-HEP'11)

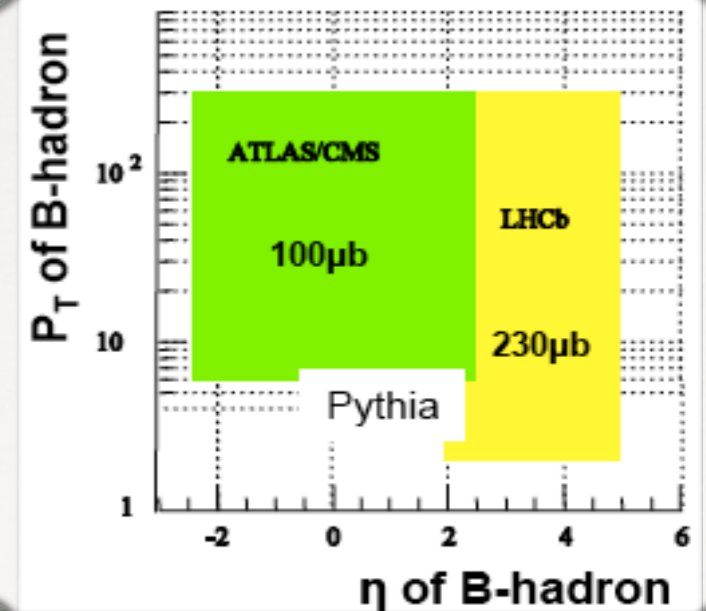
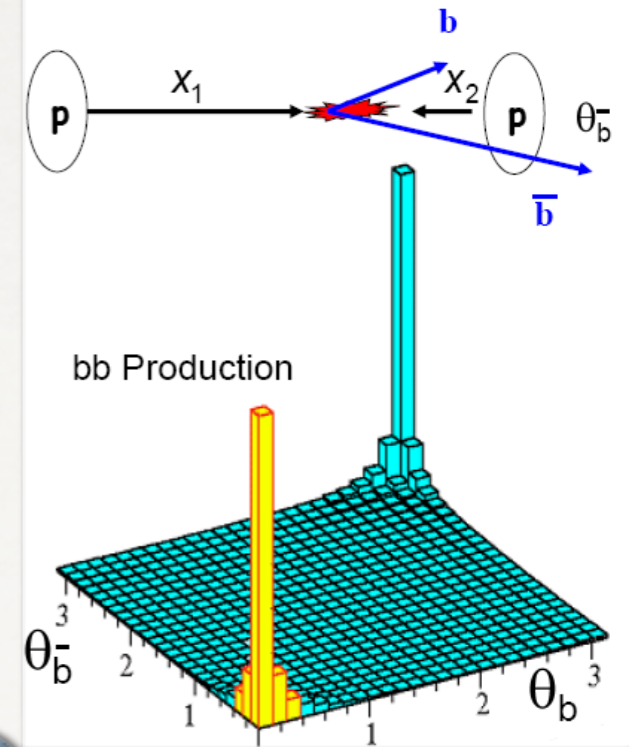


# The LHC: a hadronic beauty and charm factory

$$\begin{aligned}\sigma(pp \rightarrow \text{inel}) &= 60 \text{ mb} \\ \sigma(pp \rightarrow c\bar{c}) &= 6 \\ \sigma(pp \rightarrow b\bar{b}) &= 0.3\end{aligned}$$



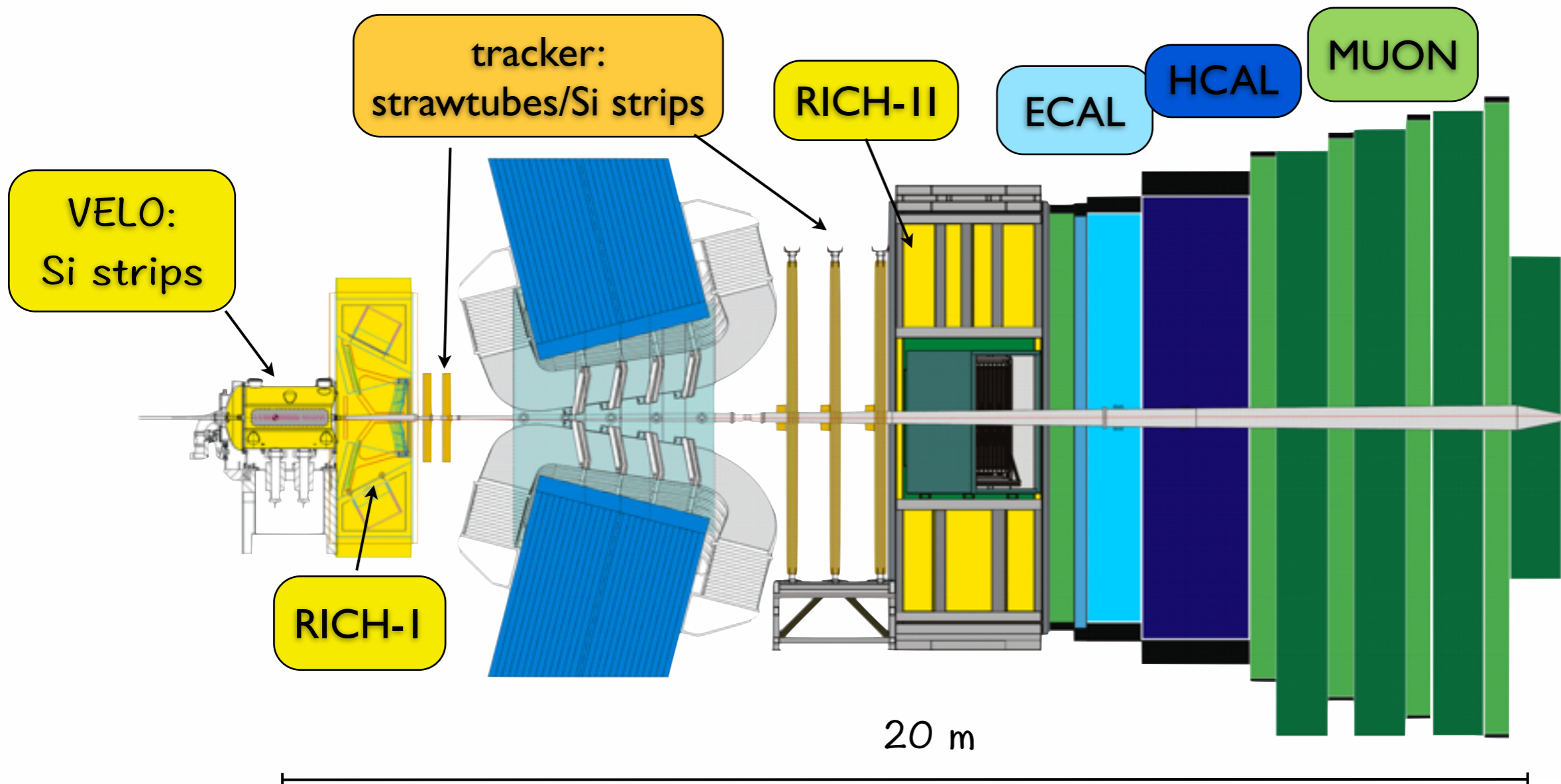
Gluon-Gluon-Fusion:





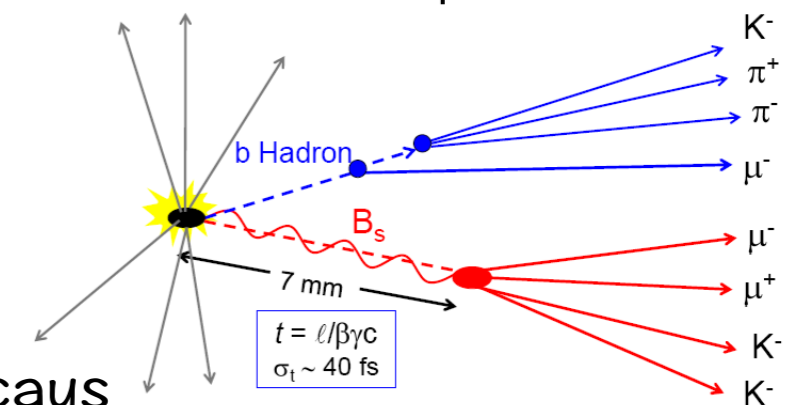
# The LHCb experiment

The LHCb Detector  
JINST 3 (2008) S08005



Requirements dictated by flavour physics:

- ▶ vertex resolution: resolve  $B_s$  oscillations
- ▶ excellent  $p/\pi/K/\mu/e$  separation: flavour tagging; rare decays
- ▶ momentum resolution: resolve  $B/B_s, D/D_s$ , etc





## Challenge:

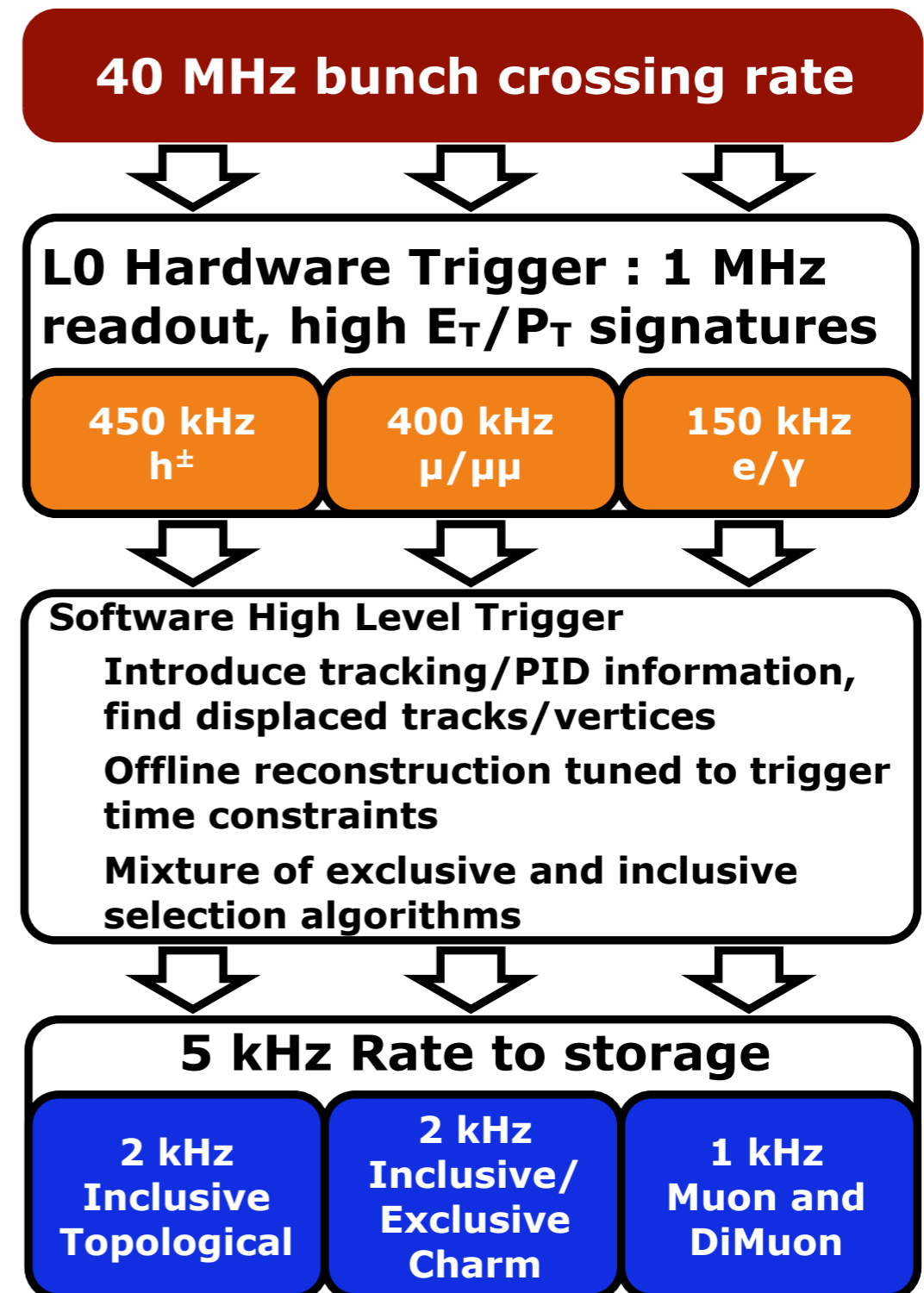
- high efficiency for manageable data rates

## Main backgrounds:

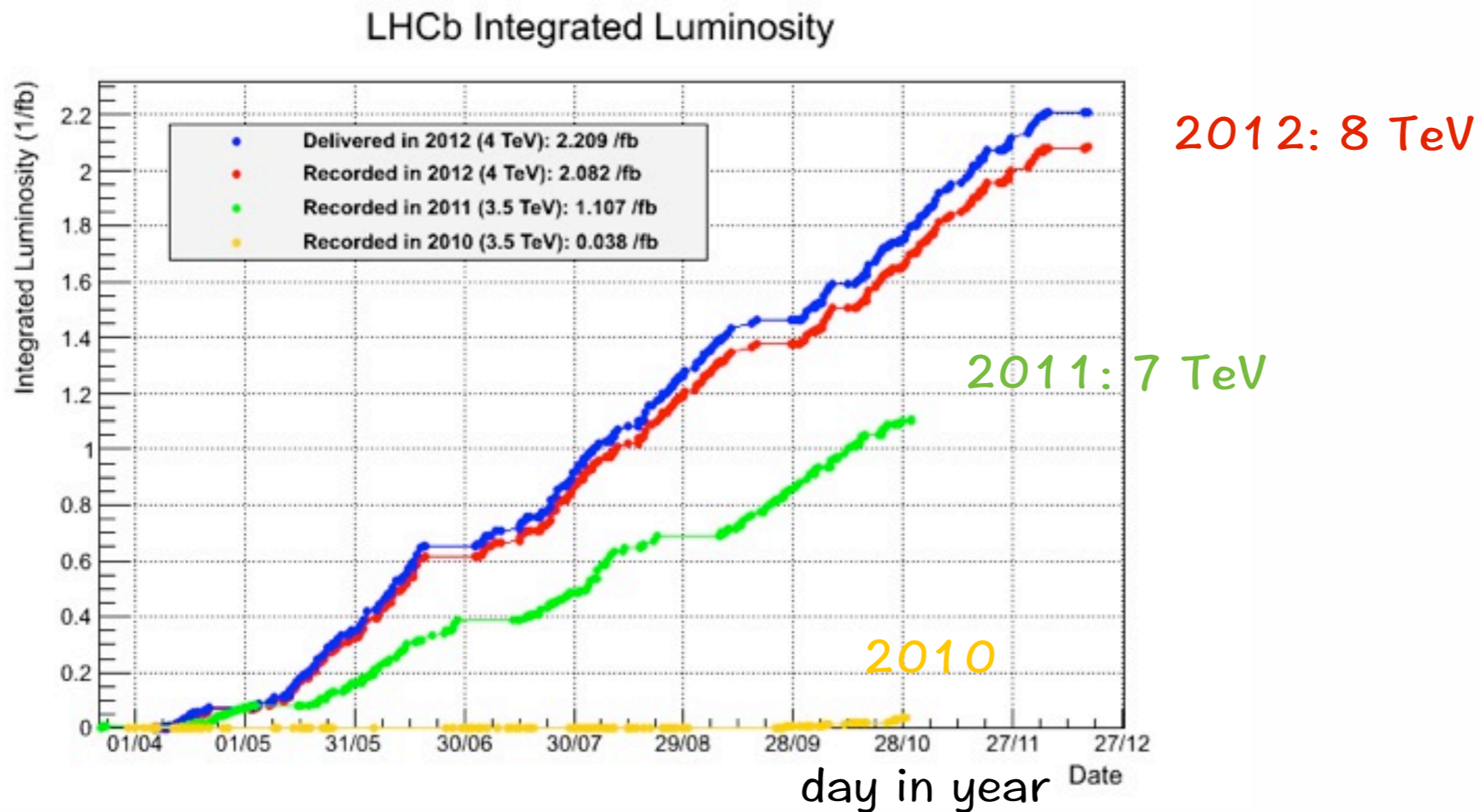
- minimum bias
- less interesting charm and beauty decays

## Handles:

- high  $p_T$  signals (hadrons, electrons, muons)
- displaced vertices



# The LHCb data set



LHC is currently shut down to prepare for 13 TeV running

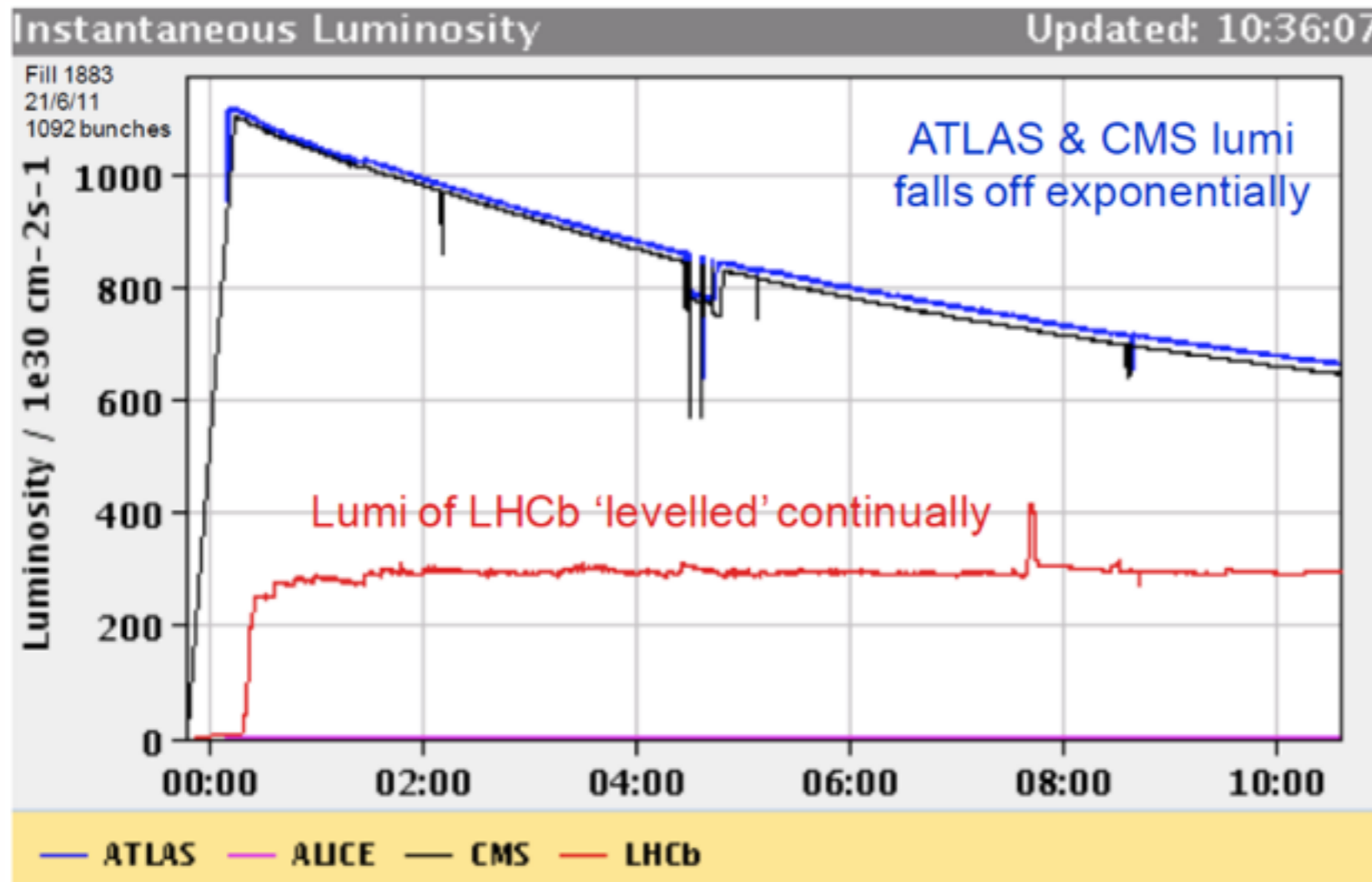
most analyses reported in this talk used 2011 dataset (1/fb) only

... many 2012 (2/fb) results will come in spring



# running strategy: reduce pile-up

- LHCb tracking/PID performance sensitive to pile-up
- beams at LHCb displaced from head-on to reduce instantaneous luminosity



optimal  $L_{\text{inst}}$ :

$$2-4 \times 10^{32} / \text{cm}^2/\text{s}$$

➡ LHCb collects less luminosity than CMS and ATLAS, but at much lower pile-up

# LHCb Physics

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LHCb physics menu [\(link\)](#)

- charm and beauty: cross-sections, spectroscopy, lifetimes, CP violation, mixing, rare decays, states, ...
- electroweak production: Z, W, jets, ...
- heavy ions: e.g. J/ψ in pA collisions
- ...

Selected results for this talk

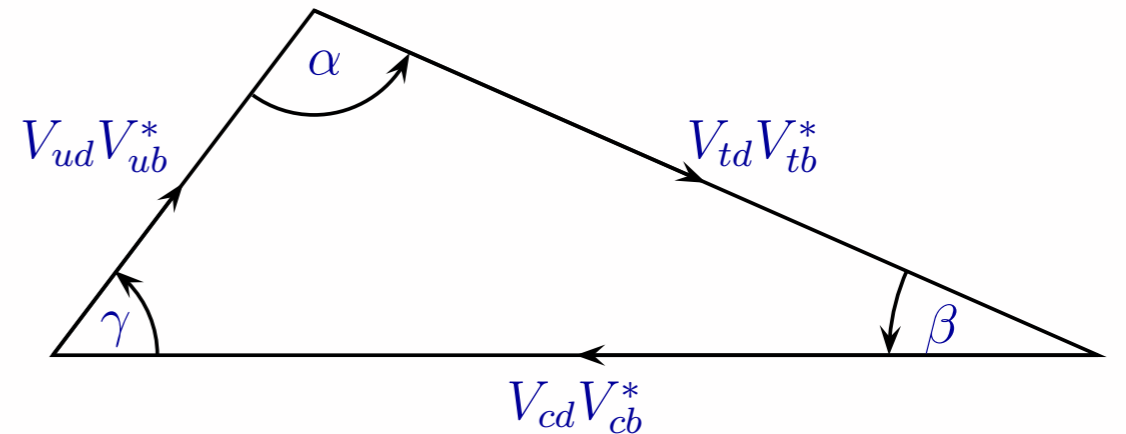
- gamma: the least well known angle of the CKM matrix
- mixing and time-dependent CP violation in the  $B_s$  system
- rare decays:  $B_s \rightarrow \mu^+ \mu^-$  and  $B_s \rightarrow K^{*0} \mu^+ \mu^-$



# The unitarity triangle

- CKM matrix is unitary and has only 4 independent physical parameters
- unitarity constraints expressed as 'triangles' in complex plane, e.g.

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0 \quad \Rightarrow$$



$$\alpha = \arg \left[ -\frac{V_{td} V_{tb}^*}{V_{ud} V_{ub}^*} \right] \quad \beta = \arg \left[ -\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*} \right] \quad \gamma = \arg \left[ -\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} \right]$$

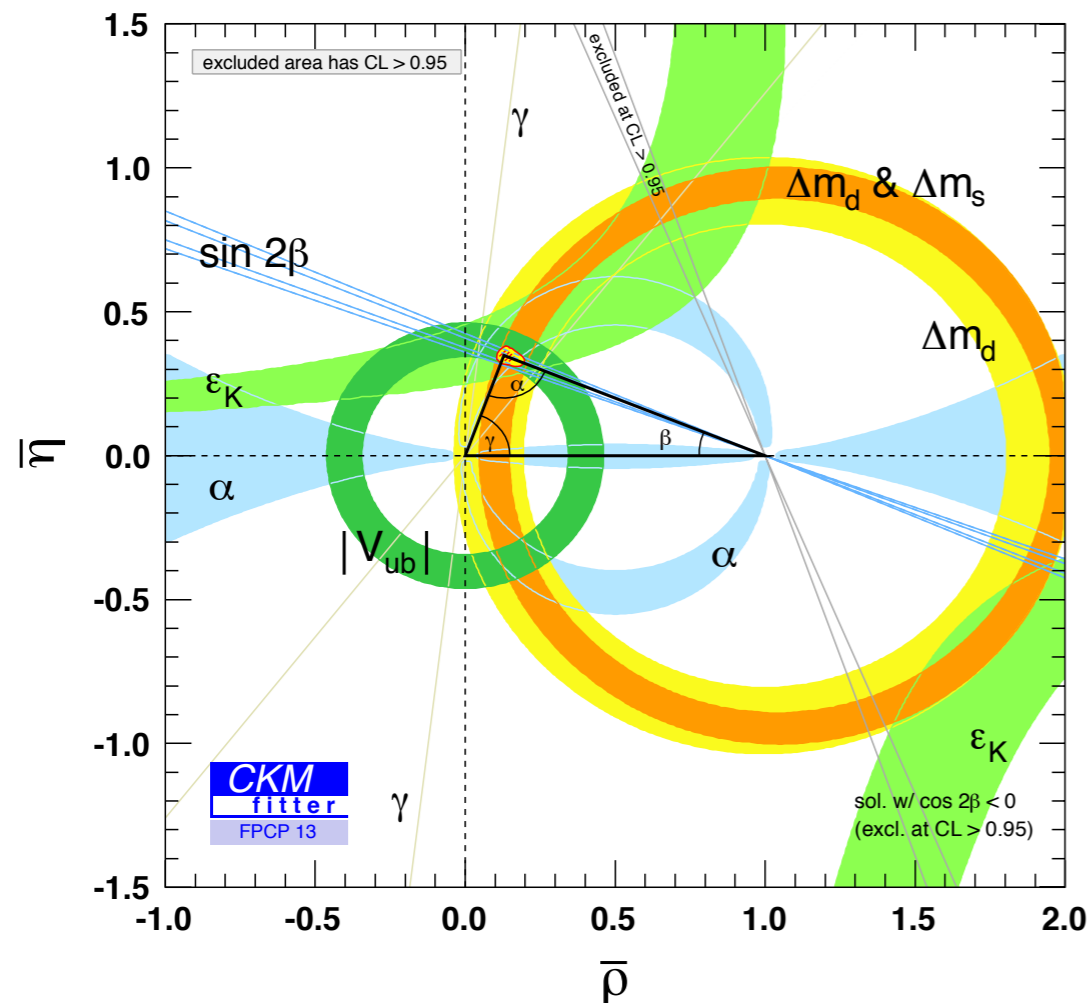
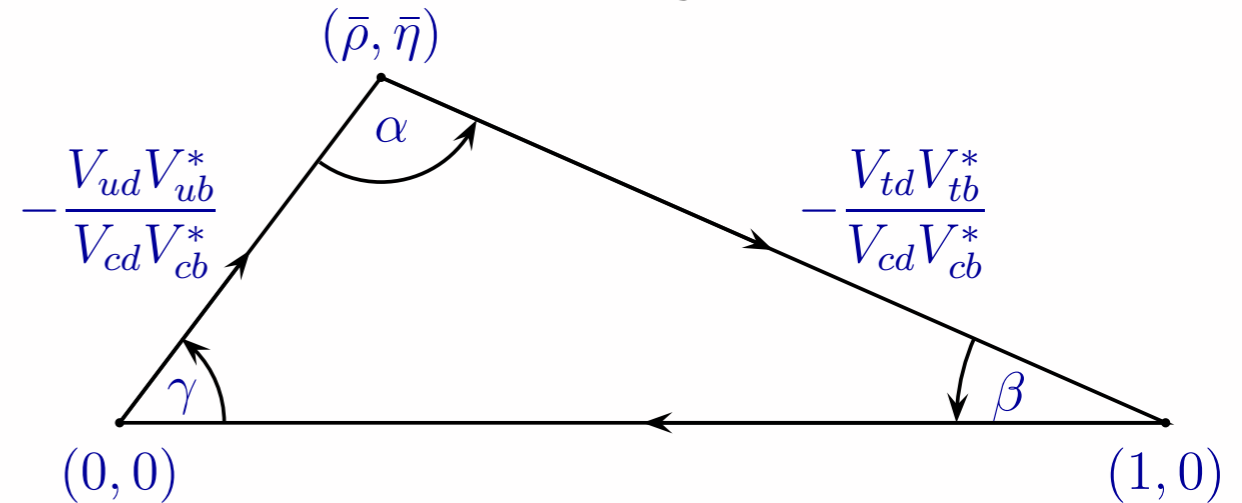
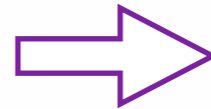
In Wolfenstein phase convention:

$$V = \begin{pmatrix} \blacksquare & \blacksquare & e^{-i\gamma} \\ -\blacksquare & \blacksquare & \blacksquare \\ e^{-i\beta} & -e^{i\beta} & \blacksquare \end{pmatrix}$$

# The unitarity triangle

- CKM matrix is unitary and has only 4 independent physical parameters
- unitarity constraints expressed as 'triangles' in complex plane, e.g.

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

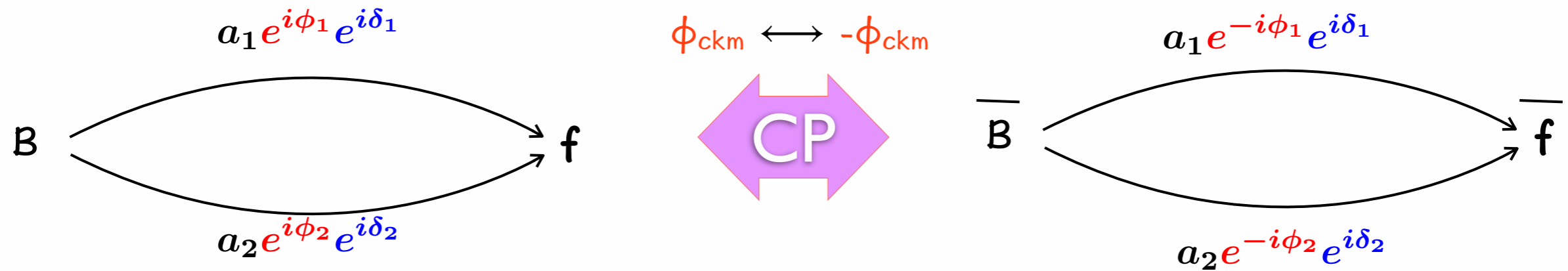


- angle  $\gamma$  has largest exp. uncertainty
- can be constrained with  $B \rightarrow DK$  decays

# Observing CP violation

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- CPV arises due to interfering amplitudes with different CKM phases

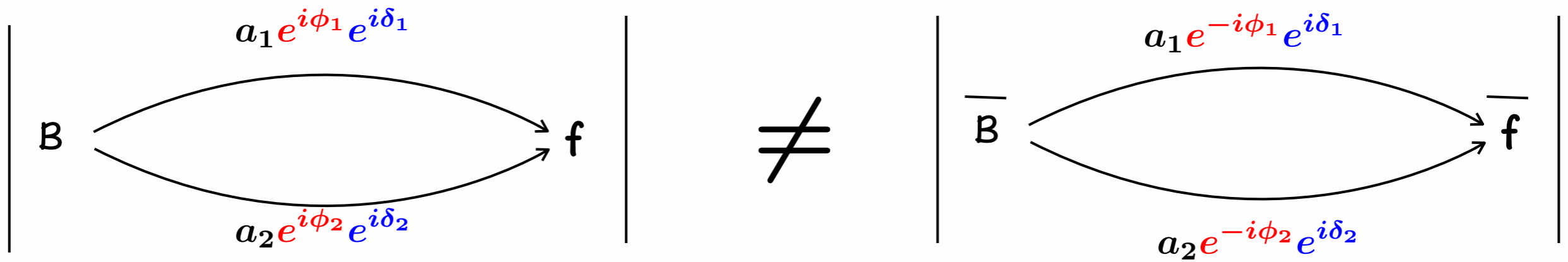




# Observing CP violation

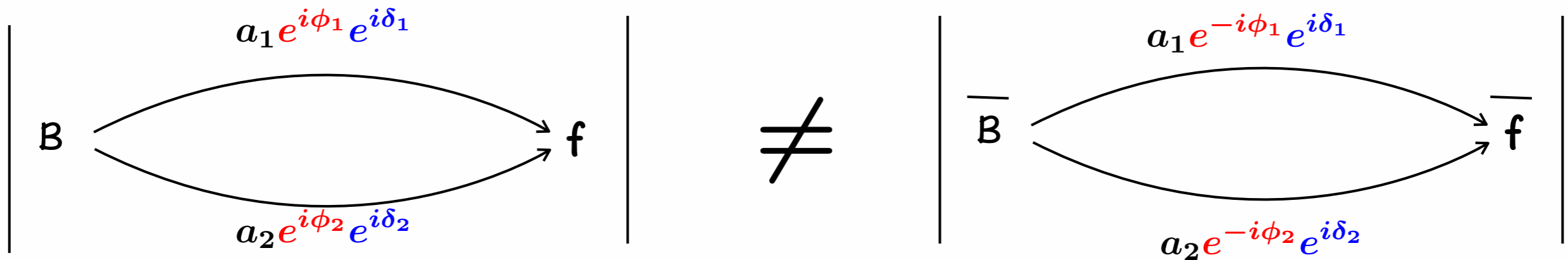
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- CPV arises due to interfering amplitudes with different CKM phases



# Observing CP violation

- CPV arises due to interfering amplitudes with different CKM phases



- two types of phenomena

- charge asymmetry (direct CPV)

$$A_f^{\text{ch}} \equiv \frac{N(B \rightarrow f) - N(\bar{B} \rightarrow \bar{f})}{N(B \rightarrow f) + N(\bar{B} \rightarrow \bar{f})}$$

- time-dependent asymmetry (indirect CPV)

arises if  $B^0$  and  $\bar{B}^0$  have common final state

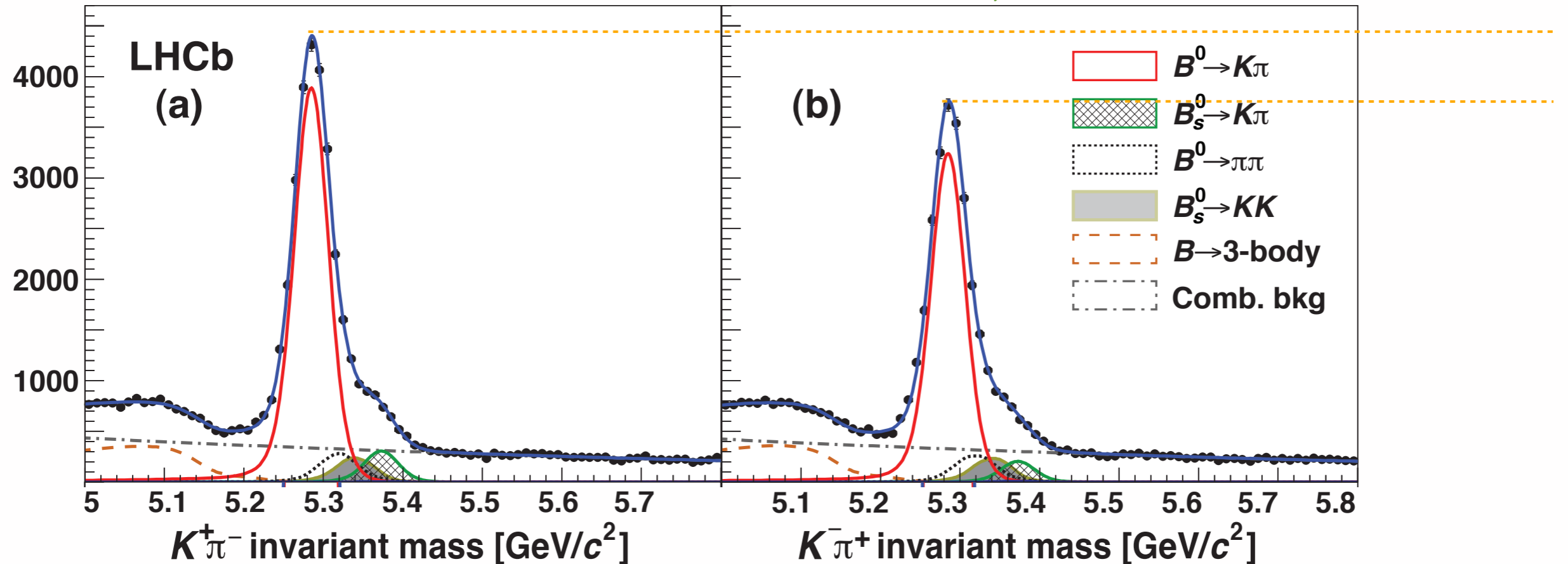
$$A(t) \equiv \frac{N(\bar{B}^0(t) \rightarrow f) - N(B^0(t) \rightarrow f)}{N(\bar{B}^0(t) \rightarrow f) + N(B^0(t) \rightarrow f)}$$

# Large direct CP violation exists in ...

... two-body decays of  $B^0$  mesons

arXiv:1304.6173,PRL.110, 221601(2013)

41000  $B_d \rightarrow K\pi$  candidates in 1/fb



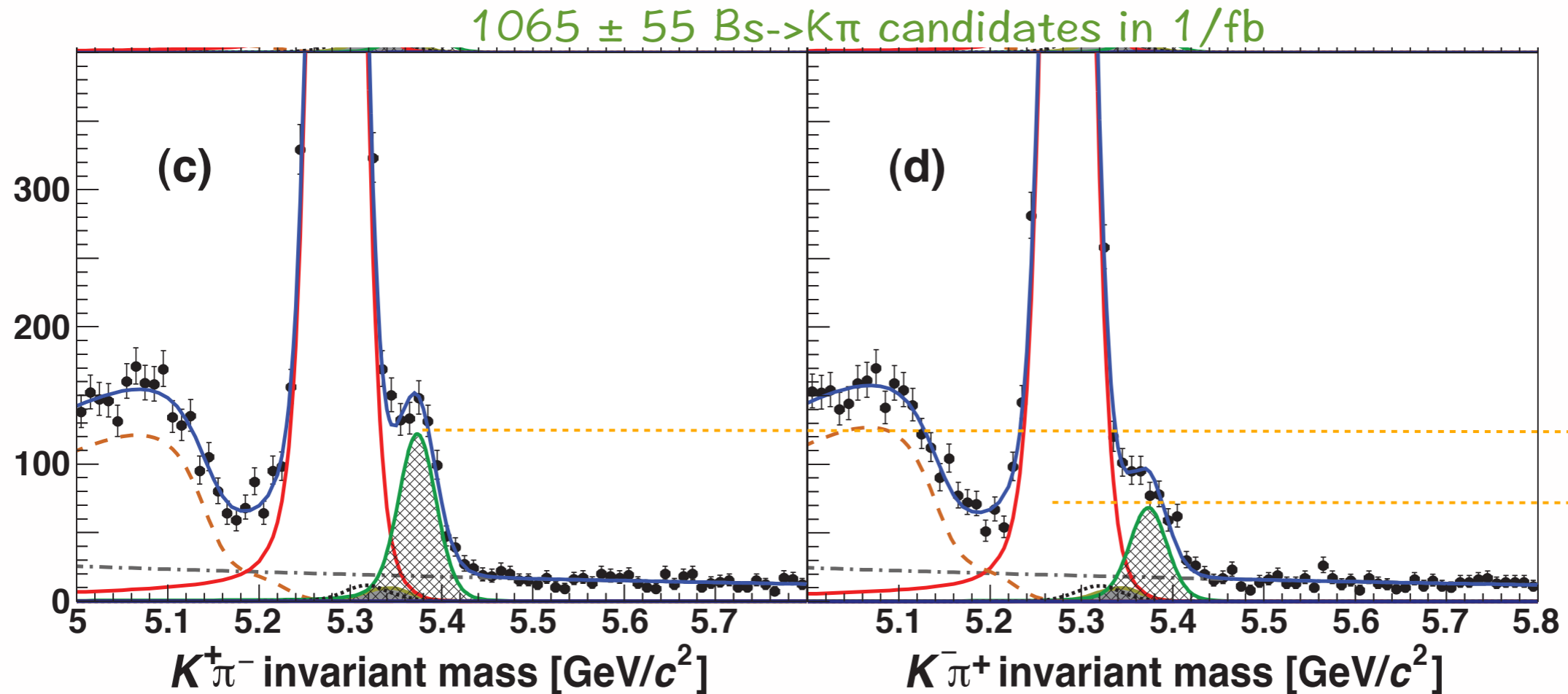
$$A_{CP}(B^0 \rightarrow K^+\pi^-) = -0.080 \pm 0.007 (\text{stat}) \pm 0.003 (\text{syst})$$

(in agreement with Belle and Babar, but with smaller uncertainty)

# Large direct CP violation exists in ...

... two-body decays of  $B_s$  mesons

arXiv:1304.6173,PRL.110, 221601(2013)



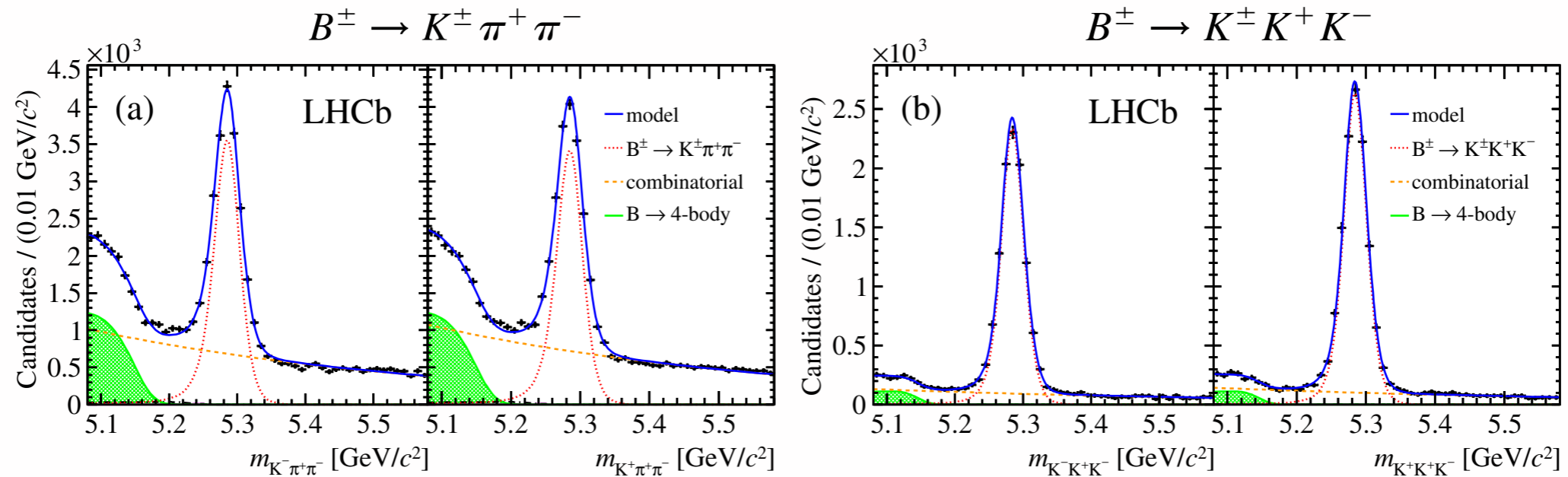
$$A_{CP}(B_s^0 \rightarrow K^- \pi^+) = 0.27 \pm 0.04 (\text{stat}) \pm 0.01 (\text{syst}).$$

(first observation of direct CPV in  $B_s$  decays)

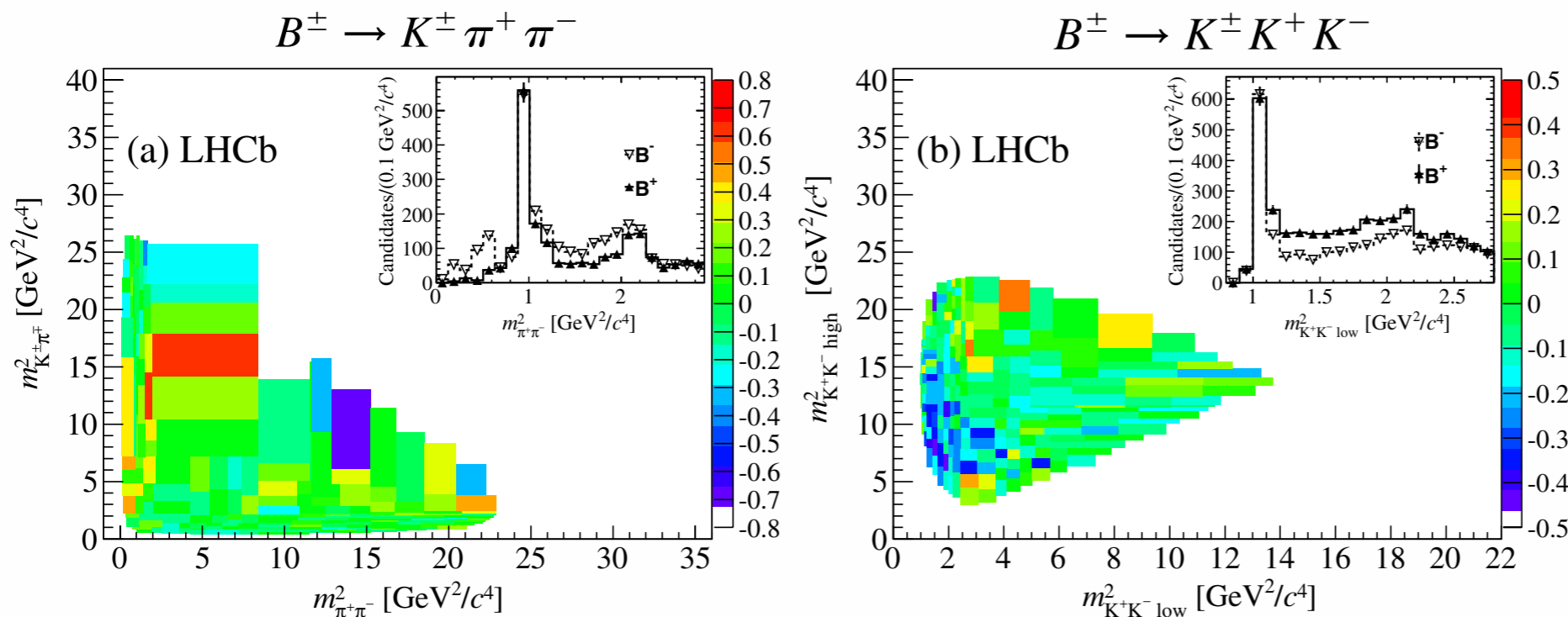
# Large direct CP violation exists ...

arXiv:1306.1246, PRL 111, 101801 (2013)

... and in 3-body  $B^+$  decays



over 50% in some parts of the Dalitz plane?!

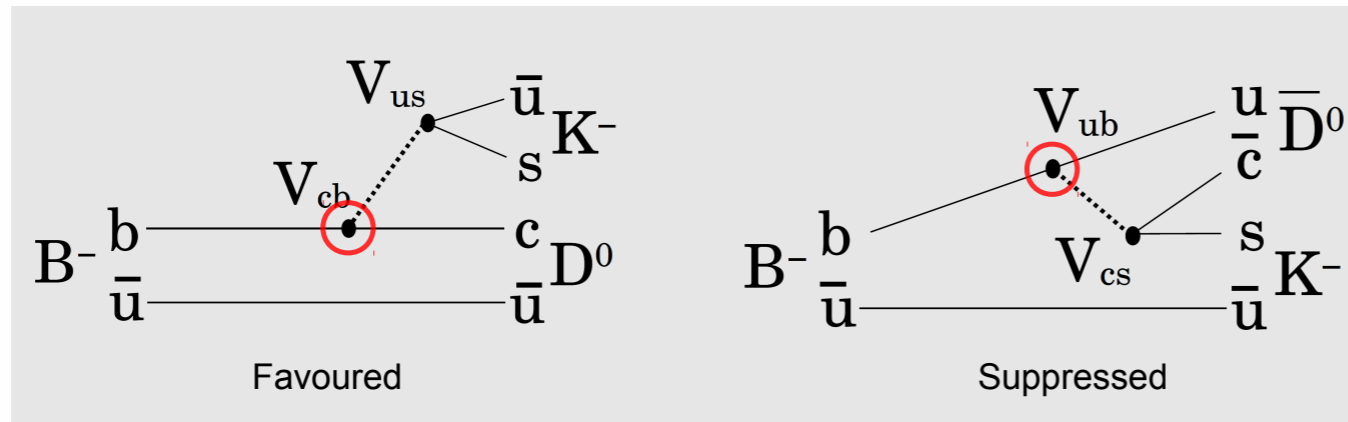


Evidence for  
“compound” CPV?  
(Cheng, Chua and Soni,  
PRD 71, 014030 (2005))

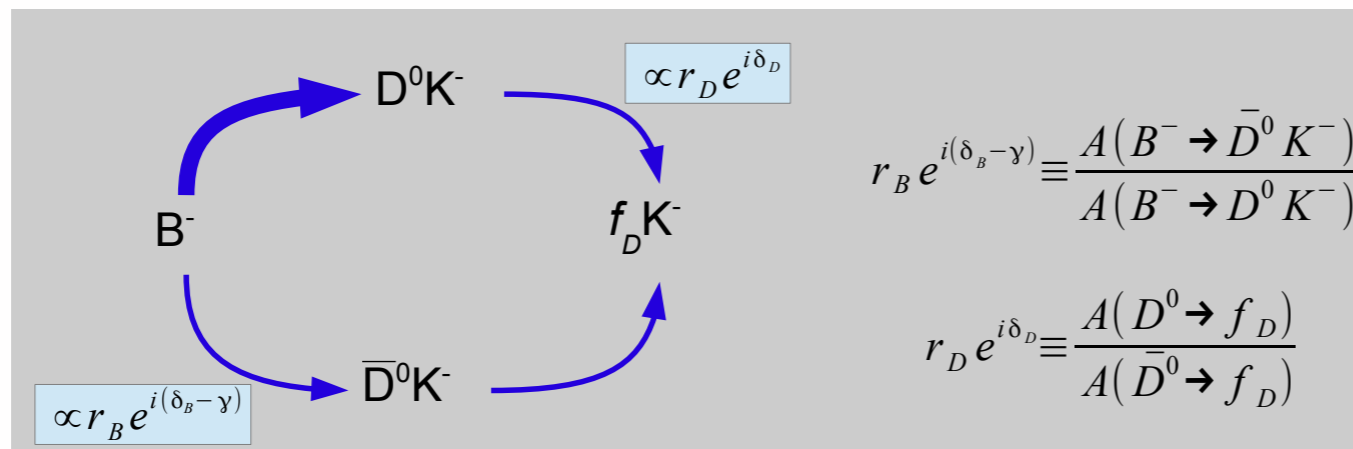


# Measurement of 'gamma'

- CP violation in  $B^\pm \rightarrow D^0 K^\pm$  occurs via 'tree-level' diagrams: small theoretical uncertainty, no 'background' from NP
- interfering amplitudes



'hadronic' parameters are a nuisance but can be extracted from the data



- need common final state for  $D^0$  and anti- $D^0$ 
  - GLW: CP eigenstate ( $D^0 \rightarrow \pi\pi, KK$ )
  - ADS: suppressed  $D^0$  with favoured anti- $D^0$  decay
  - GGSZ: Dalitz analysis of  $K_s^0 K^+ K^-$  or  $K_s^0 \pi^+ \pi^-$

[Gronau-London-Wyler] PLB 253,483(1991), PLB 265,172 1991)  
 [Atwood-Dunietz-Soni] PRL 78,257(1997), PRD 63,036005(2001)  
 [Giri-Grossman-Soffer-Zupan] PRD 68,054018(2003)

# ADS/GWL in $B \rightarrow [\pi K] K/\pi$ and $B \rightarrow [\pi K \pi \pi] K/\pi$

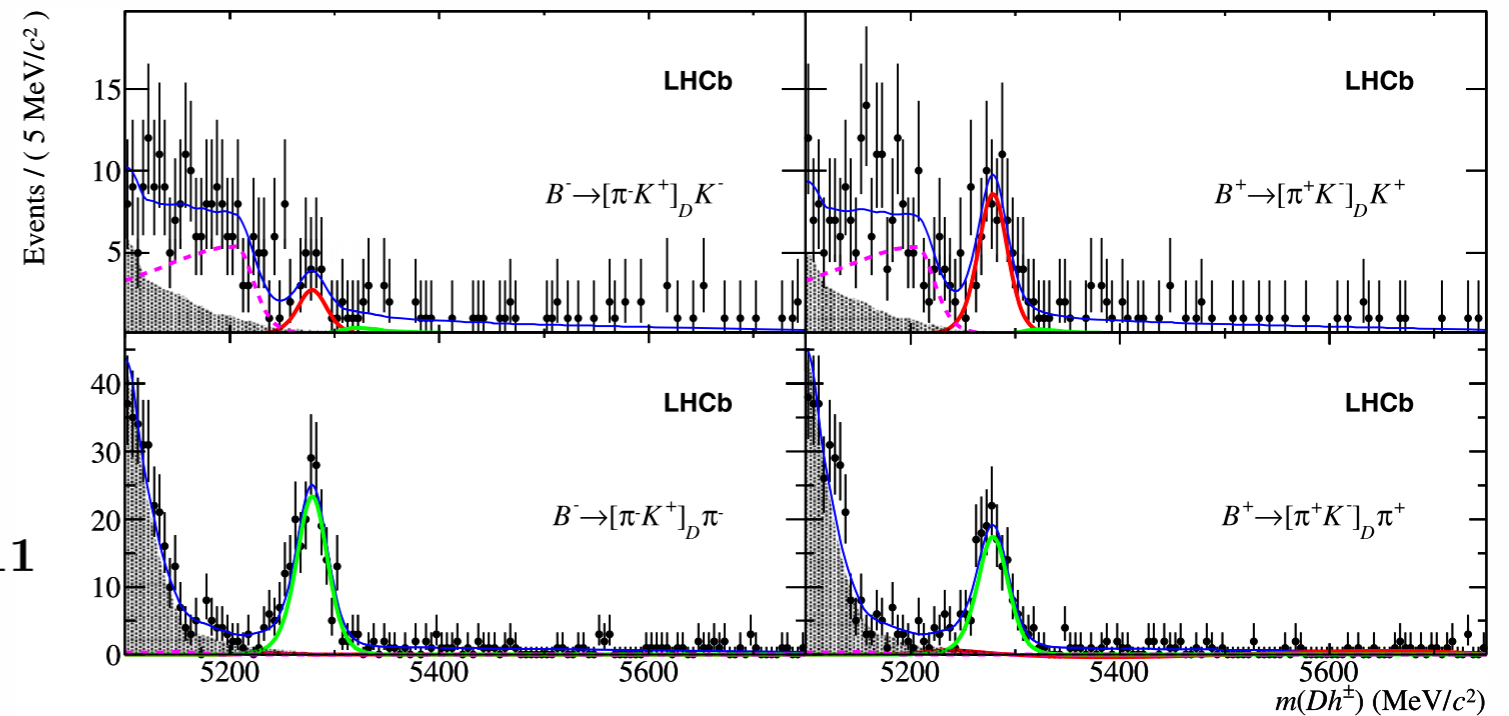
Several modes explored ....

1/fb: arXiv:1203.3662, PLB 712 (2012), 203

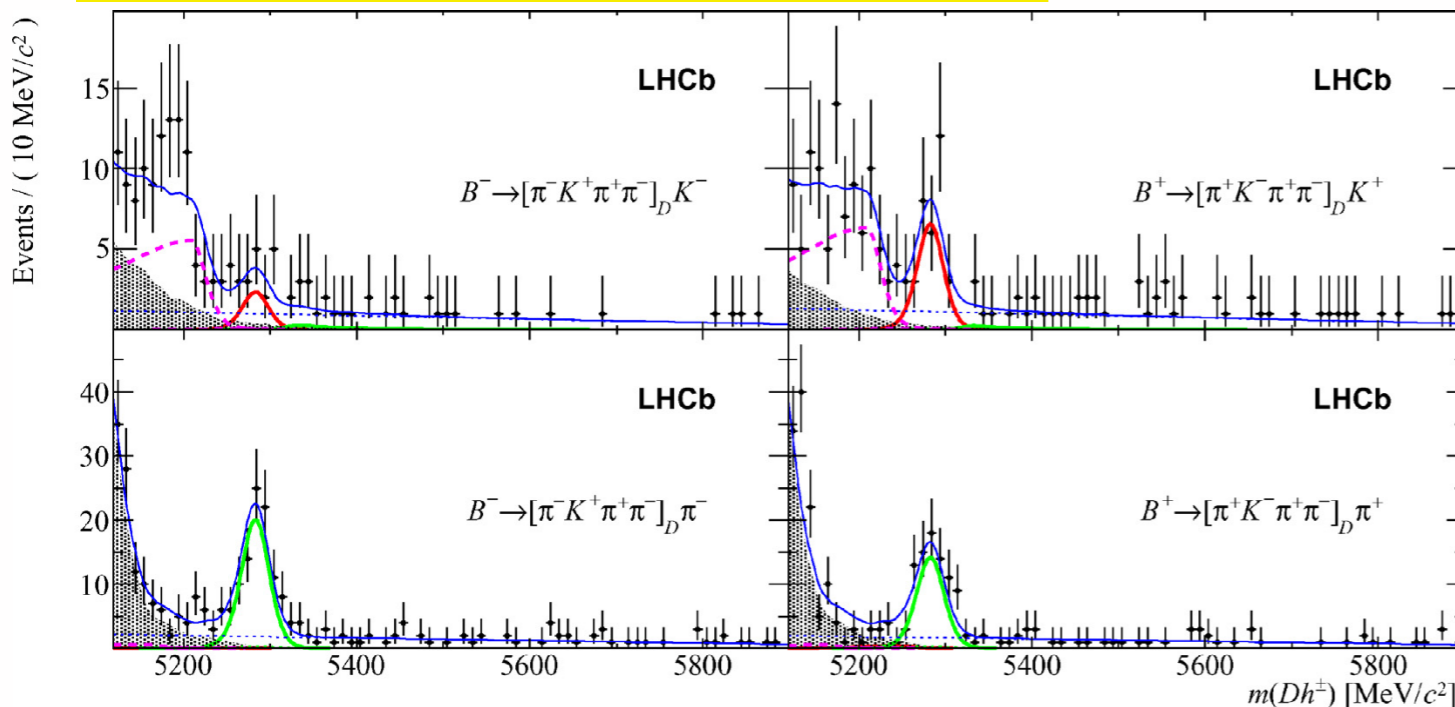
- ▶ First observation of ADS mode  $B^+ \rightarrow D^0(\pi^+ K^-) K^+$  ( $10\sigma$ )
- ▶ clear CP asymmetries:

$$A_{\text{ADS}}(D^0(\pi^- K^+) K^-) = -0.52 \pm 0.15 \pm 0.02$$

$$A_{\text{ADS}}(D^0(\pi^- K^+) \pi^-) = -0.143 \pm 0.062 \pm 0.11$$



1/fb: arXiv:1303.4646, PLB 723 (2013) 44



- ▶ First observation of ADS modes with  $D^0 \rightarrow \pi K \pi \pi$ , both in  $B \rightarrow DK$  and  $B \rightarrow \pi K$

# GGSZ: $B \rightarrow D\pi/K$ with $D \rightarrow K_S K K$ or $D \rightarrow K_S \pi\pi$

1/fb: PLB 718(2012)43

2/fb: LHCb-CONF-2013-004

- CP conjugate three body decay: extract gamma from asymmetry of  $B^+$  and  $B^-$  yield across Dalitz plane
- strong phase changes across  $D^0$  Dalitz plane: use CLEO data ([PRD82\(2010\)112006](#)) to constrain the phase

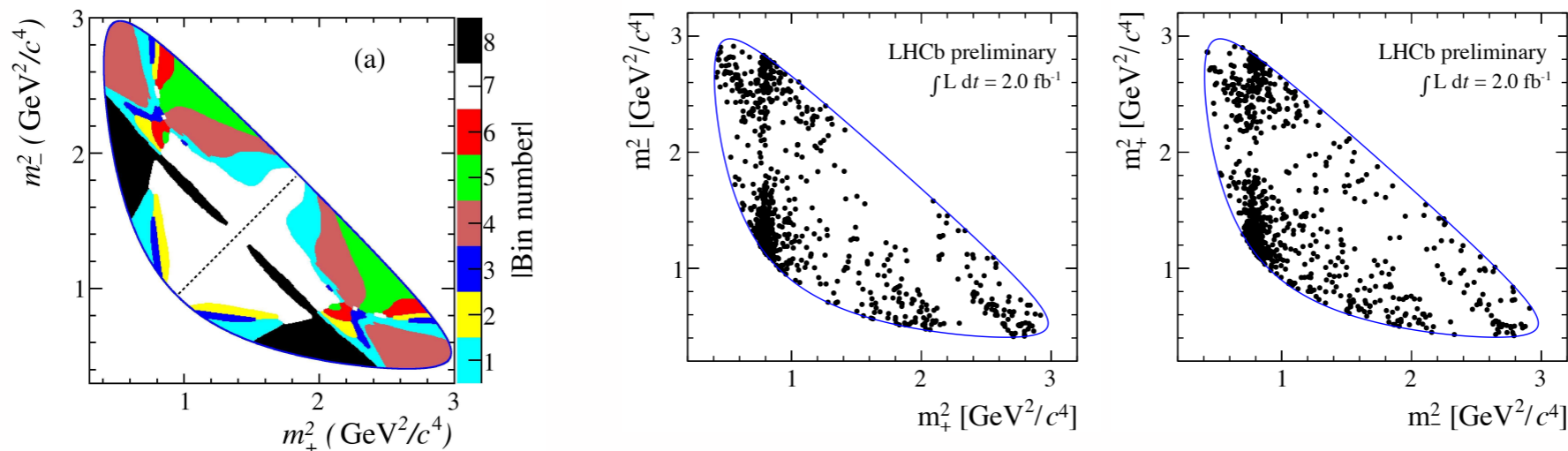


Figure 4: Dalitz plots for  $B^\pm \rightarrow (K_S^0 \pi^+ \pi^-)_D K^\pm$  decays; (left)  $B^+$ , (right)  $B^-$ .

- extract two parameters that are function of gamma:

$$x_\pm \equiv r_B \cos(\delta_B \pm \gamma), \quad y_\pm \equiv r_B \sin(\delta_B \pm \gamma).$$

- result (2012, 2/fb):

$$x_+ = (-8.7 \pm 3.1 \pm 1.6 \pm 0.6) \times 10^{-2}$$

$$x_- = (5.3 \pm 3.2 \pm 0.9 \pm 0.9) \times 10^{-2}$$

$$y_+ = (0.1 \pm 3.6 \pm 1.4 \pm 1.9) \times 10^{-2}$$

$$y_- = (9.9 \pm 3.6 \pm 2.2 \pm 1.6) \times 10^{-2}$$

$\swarrow$  stat.     $\uparrow$  LHC-syst.     $\nwarrow$  CLEO-syst.



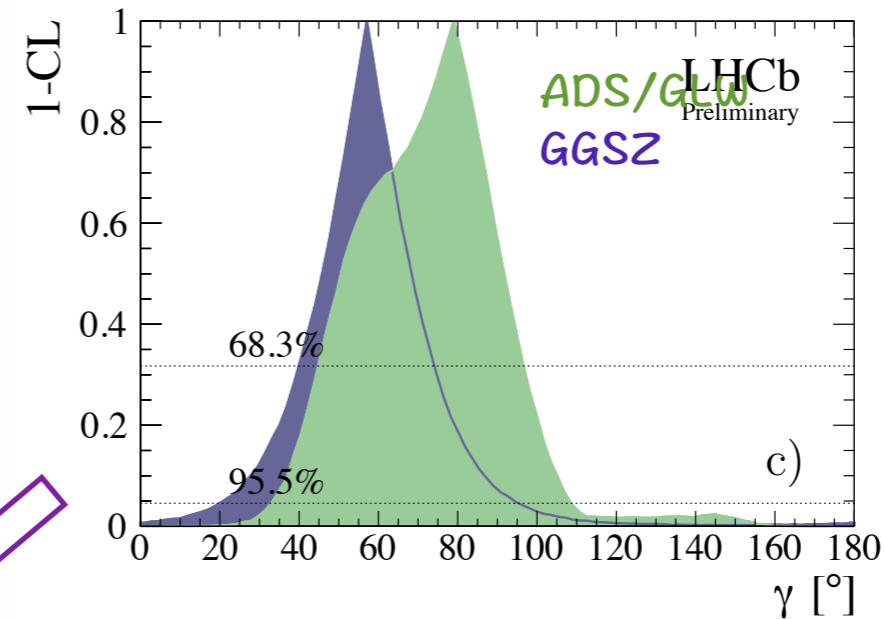
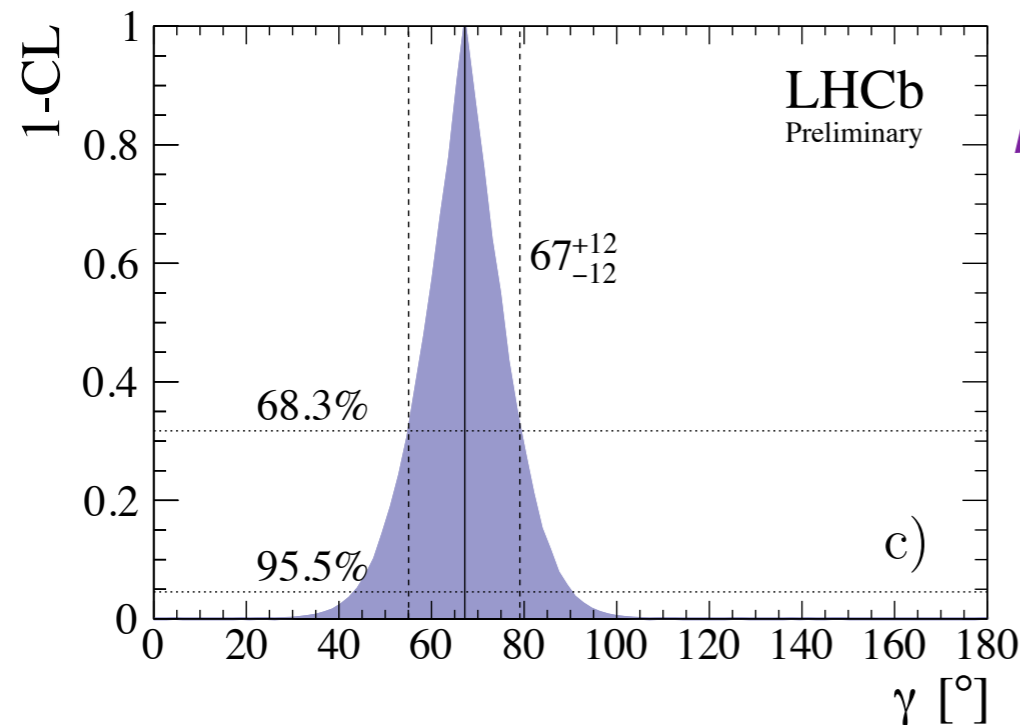
# gamma: combination of LHCb measurements

- LHCb B->DK results combined to extract single confidence interval

Combination uses

- ▶ ADS/GLW: 1/fb (2011)
- ▶ GGSZ: 3/fb (2011 + 2012)

LHCb-CONF-2013-006 (prel.)



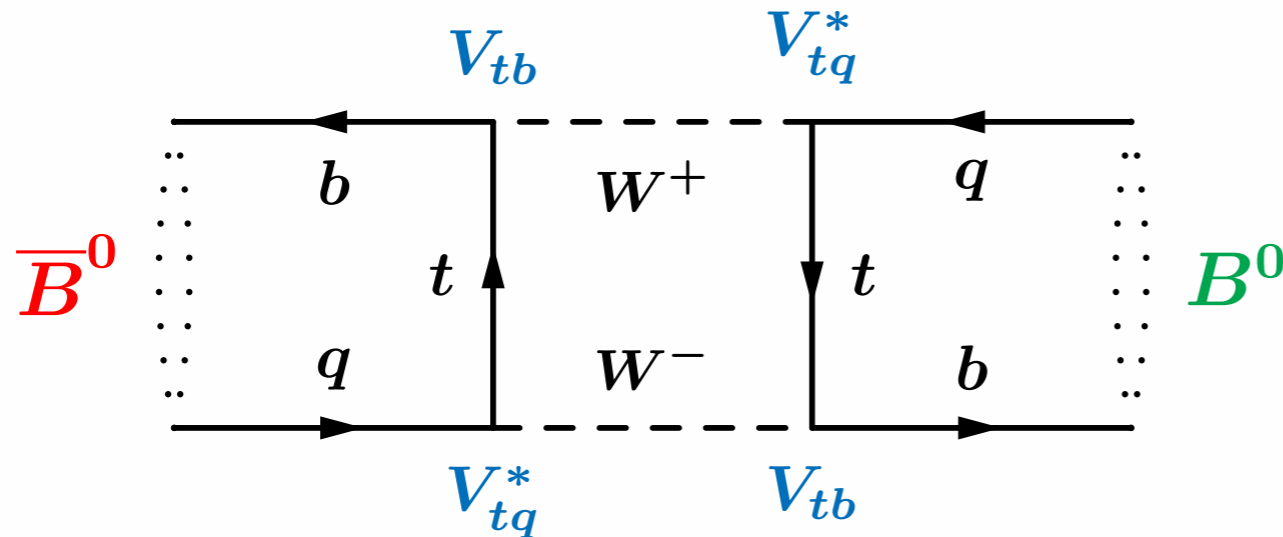
- ❖ result:  $\gamma = (67 \pm 12)^\circ$  preliminary
- ❖ most precise value from single experiment
- ❖ in good agreement with Belle and Babar

- sensitivity to gamma will improve further when ADS/GLW extended to 3/fb
- combination with B->D $\pi$  exists, but only for 2011 data

arXiv:1305.2050,PLB726(2013)151

# Neutral meson mixing

- neutral meson flavour states mix via this weak interaction loop diagram



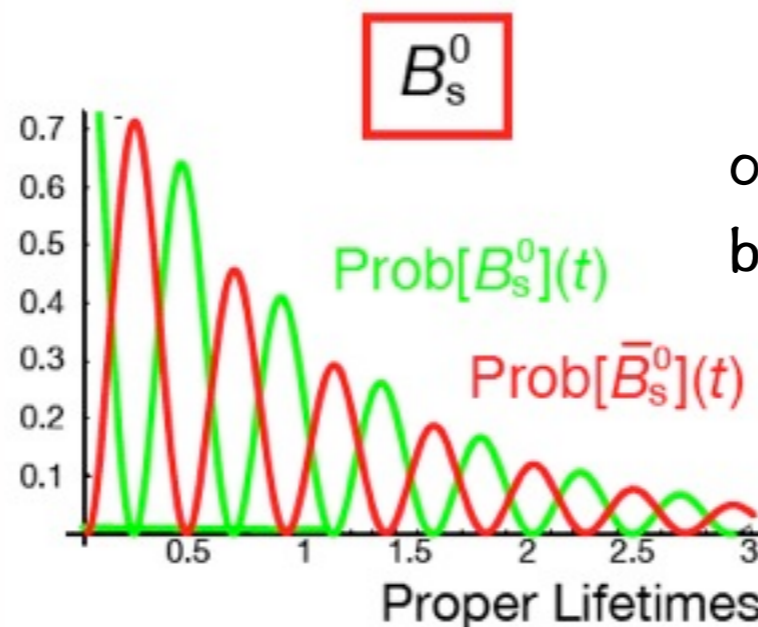
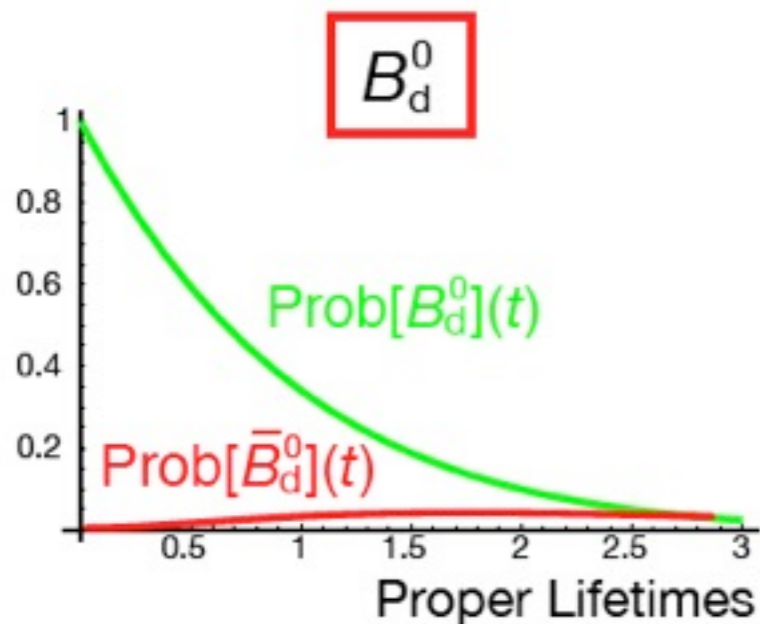
2nd order weak process

very sensitive to BSM contributions

- heavy (H) and light (L) mass eigenstates are mixture of flavour eigenstates

$$|B_L\rangle = p |B^0\rangle + q |\bar{B}^0\rangle$$

$$|B_H\rangle = p |B^0\rangle - q |\bar{B}^0\rangle$$



oscillation frequency governed by mass difference

# Neutral meson mixing

- mixing and decay described by Schroedinger equation

$$i \frac{d}{dt} \begin{pmatrix} \langle B^0 | B(t) \rangle \\ \langle \bar{B}^0 | B(t) \rangle \end{pmatrix} = \begin{pmatrix} M_{11} - \frac{i}{2} \Gamma_{11} & M_{12} - \frac{i}{2} \Gamma_{12} \\ M_{21} - \frac{i}{2} \Gamma_{21} & M_{22} - \frac{i}{2} \Gamma_{22} \end{pmatrix} \begin{pmatrix} \langle B^0 | B(t) \rangle \\ \langle \bar{B}^0 | B(t) \rangle \end{pmatrix}$$

- phenomenology described in terms of 5 observables

$$m \equiv \frac{m_H + m_L}{2}$$

$$\Gamma \equiv \frac{\Gamma_H + \Gamma_L}{2}$$

$$\Delta m \equiv m_H - m_L \simeq 2|M_{12}|$$

$$\Delta\Gamma \equiv \Gamma_L - \Gamma_H \simeq 2|\Gamma_{12}| \cos \phi_{12}$$

$$1 - \left| \frac{q}{p} \right|^2 = \frac{|\Gamma_{12}|}{|M_{12}|} \sin \phi_{12}$$

CP violation,  
small in SM

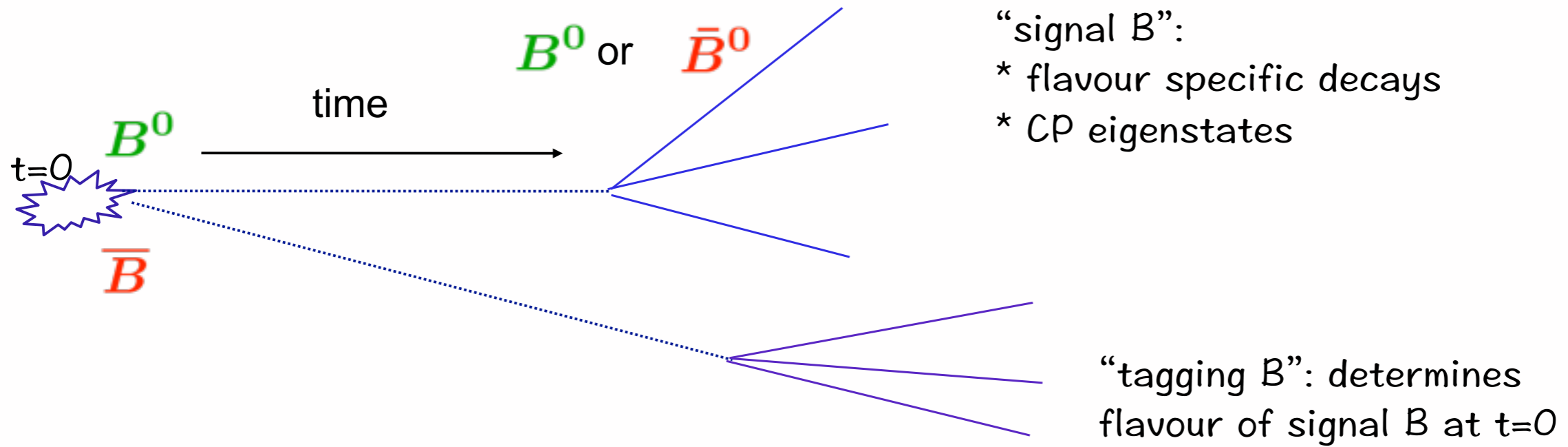
clean predictions in SM

- $M_{12}$  and  $\Gamma_{12}$ : mixing amplitudes via off-shell and on-shell states, respectively
- CP violating phase  $\phi_{12}$  is their relative phase:  $\phi_{12} \equiv \arg(-M_{12}/\Gamma_{12})$ 
  - in SM close to zero both in Bs and Bd system

# Experimental observables in B mixing

---

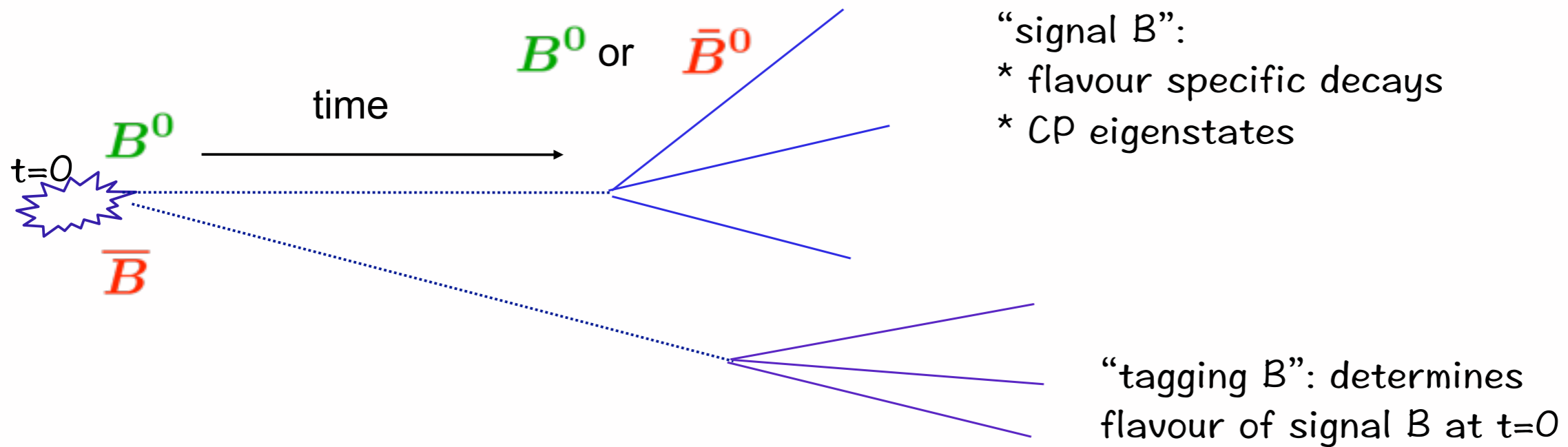
- decay time allows to observe time-evolution





# Experimental observables in B mixing

- decay time allows to observe time-evolution



- four types of observables

1. oscillations:

$$A^{\text{mix}}(t) = \frac{N^{B^0\bar{B}^0|\bar{B}^0B^0} - N^{B^0B^0|\bar{B}^0\bar{B}^0}}{N^{B^0\bar{B}^0|\bar{B}^0B^0} + N^{B^0B^0|\bar{B}^0\bar{B}^0}} = \cos(\Delta m_q t)$$

2. CPV in mixing:

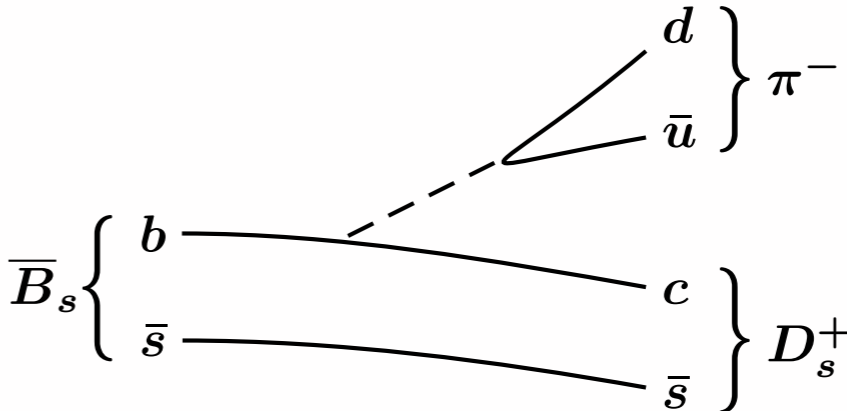
$$a_{fs} \equiv \frac{N^{B^0B^0} - N^{\bar{B}^0\bar{B}^0}}{N^{B^0B^0} + N^{\bar{B}^0\bar{B}^0}} = \frac{1 - |q/p|^4}{1 + |q/p|^4}$$

3. CPV in interference between mixing and decay (time-dependent CPV)

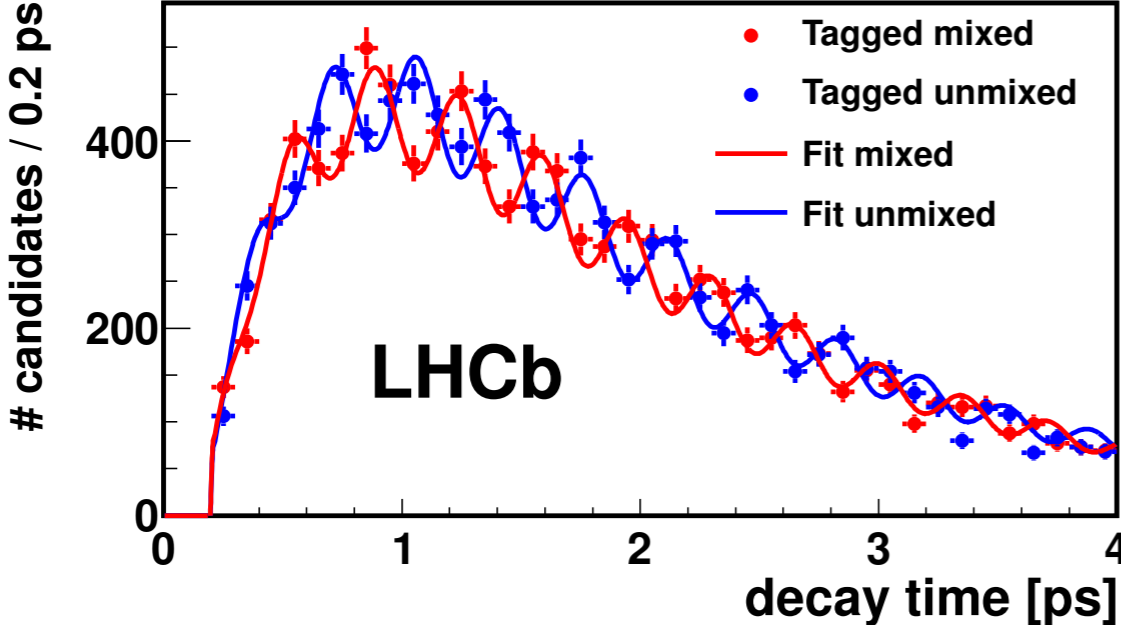
4. lifetimes

# Mixing frequency

■ most easy mode:  $B_s \rightarrow D_s^- \pi^+$



~34k signal events in 1.0/fb



opposite side tagger: Eur.Phys.J. C72(2012) 2022  
same side tagger: LHCb-CONF-2012-033

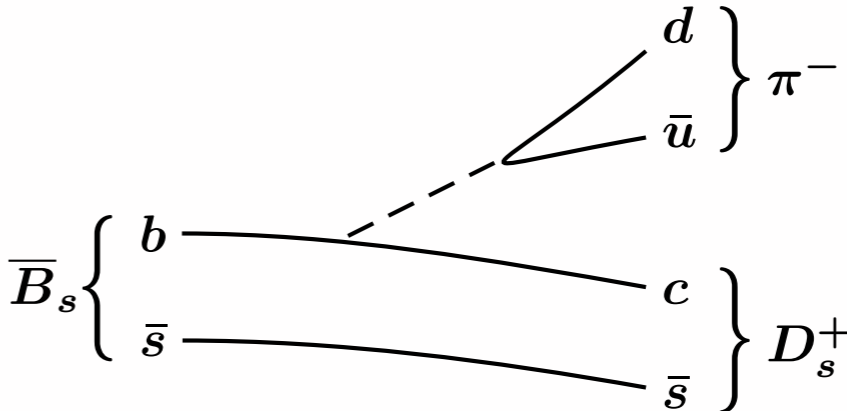
$$N_{\pm}(t) \propto \text{eff}(t) \times e^{-\Gamma t} \times \left( 1 \pm D \cos(\Delta m t) \right)$$

↑  
decay time efficiency

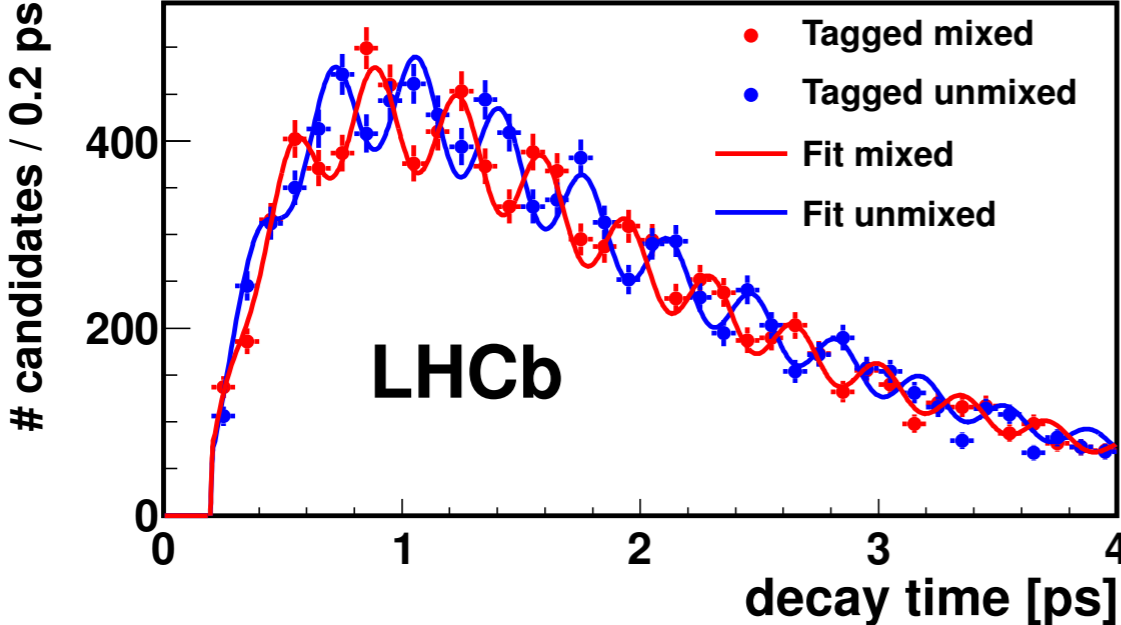
↑  
experimental 'dilution':  
- wrong flavour tagging  
- decay time resolution

# Mixing frequency

■ most easy mode:  $B_s \rightarrow D_s^- \pi^+$



~34k signal events in 1.0/fb



opposite side tagger: Eur.Phys.J. C72(2012) 2022  
 same side tagger: LHCb-CONF-2012-033

$$N_{\pm}(t) \propto \text{eff}(t) \times e^{-\Gamma t} \times \left( 1 \pm D \cos(\Delta m t) \right)$$

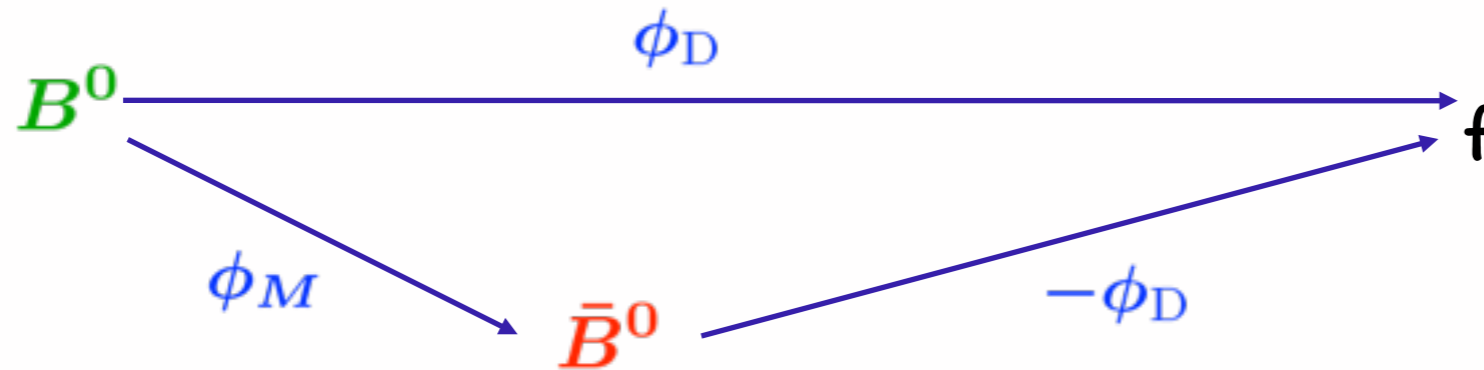
$\Delta m_s = 17.768 \pm 0.023 \text{ (stat)} \pm 0.006 \text{ (syst)}$

- result (arXiv:1304.4741):
- world's most precise measurement
- compatible with SM prediction

$$\Delta m_s^{\text{SM}} = 17.3 \pm 2.6 \quad (\text{Lenz and Nierste, 2011})$$

# Time-dependent CP violation

- common final state  $f$  : mixing induced CPV



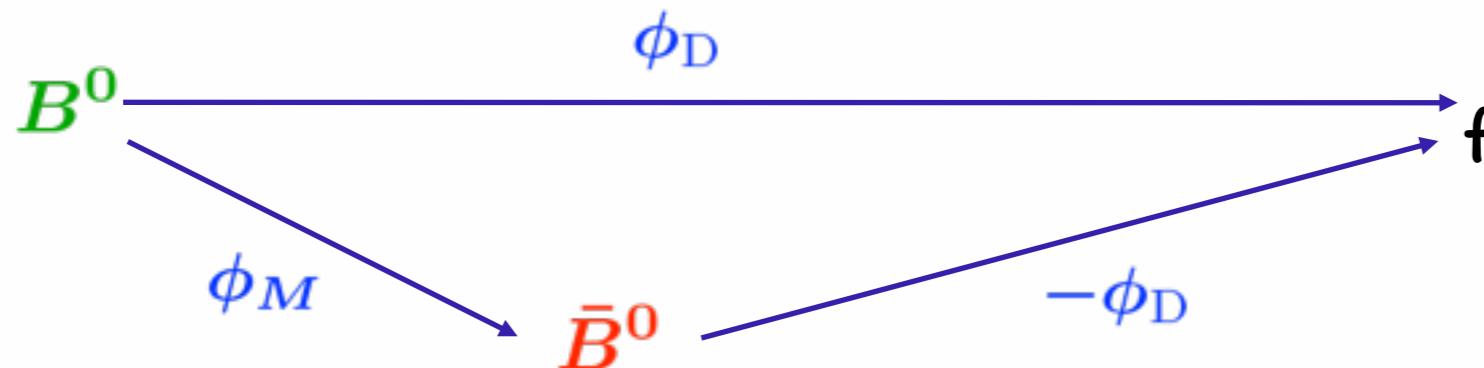
- if  $f$  is CP eigenstate, time dependent CP violation with pattern

$$A_{\text{CP}}(t) \equiv \frac{N(\bar{B} \rightarrow f) - N(B \rightarrow f)}{N(\bar{B} \rightarrow f) + N(B \rightarrow f)} = D \eta_f^{\text{CP}} \sin \phi_f \sin(\Delta m t)$$

$\eta_f^{\text{CP}} = \pm 1$   
 $\phi_f = \phi_M - 2\phi_D$   
 experimental dilution  
 flavour 'tag' at  $t=0$

# Time-dependent CP violation

- common final state  $f$  : mixing induced CPV



- if  $f$  is CP eigenstate, time dependent CP violation with pattern

$$A_{\text{CP}}(t) \equiv \frac{N(\bar{B} \rightarrow f) - N(B \rightarrow f)}{N(\bar{B} \rightarrow f) + N(B \rightarrow f)} = D \eta_f^{\text{CP}} \sin \phi_f \sin(\Delta m t)$$

$\phi_f = \phi_M - 2\phi_D$

- 'golden' modes:

$$B_d \rightarrow \psi K_s : \quad \phi_{J/\psi K_s} = 2\beta \quad \stackrel{\text{SM}}{=} 50^\circ - 55^\circ$$

$$B_s \rightarrow \psi \phi : \quad \phi_{J/\psi \phi} = -2\beta_s \quad \stackrel{\text{SM}}{=} -2.1^\circ \pm 0.1^\circ$$

(SM predictions from CKMFitter, PRD83,036004 and UFTit 2010 )

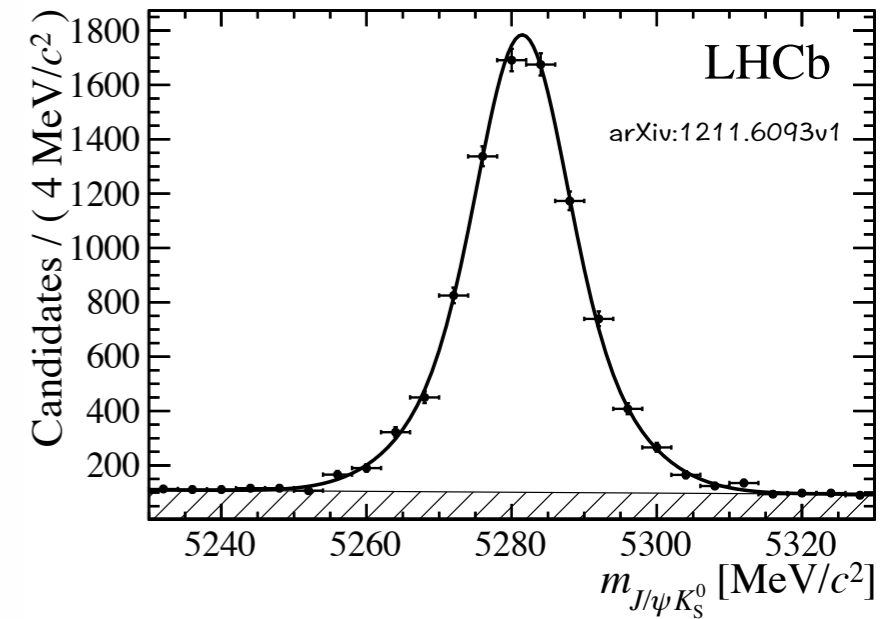
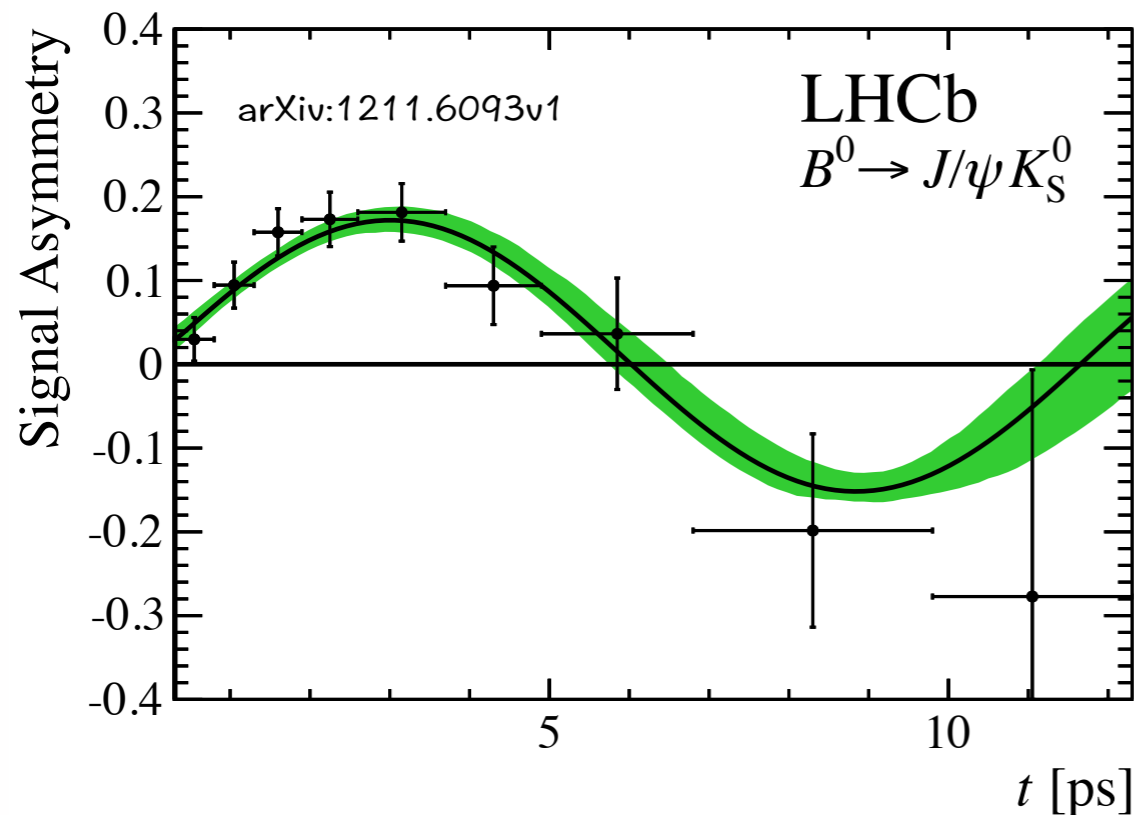
- new phase in mixing  $\Rightarrow \phi_f = \phi_f^{\text{SM}} + \delta\phi_M^{\text{NP}}$



# $B_d \rightarrow J/\psi K_s$ at LHCb

arXiv:1211.6093, PLB 721 (2013) 24-31

- important test: can we reproduce B-factory results?
- LHCb has  $\sim 8200$   $B_d \rightarrow J/\psi K_s$  candidates in 1.0/fb



$$A_{CP}(t) = D \sin 2\beta \sin(\Delta m t)$$

$$D_{\text{tag}} = 0.270 \pm 0.015$$

$$\epsilon_{\text{tag}} = (2.4 \pm 0.3)\%$$

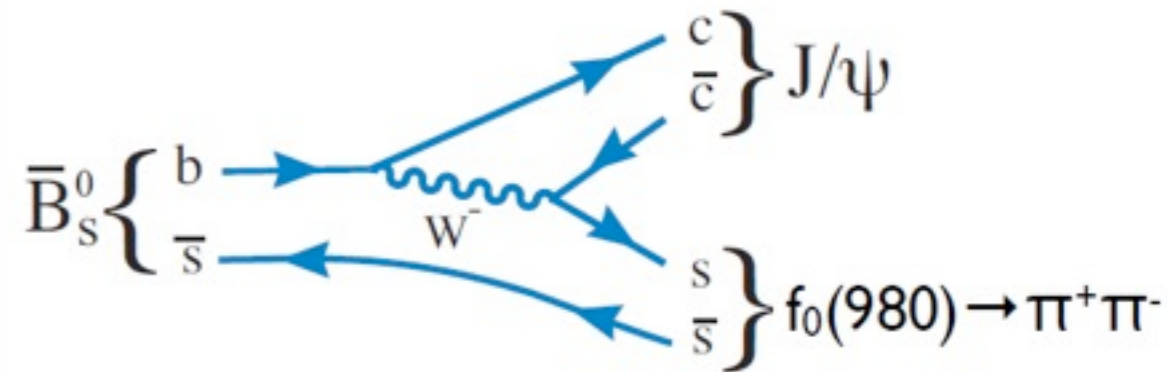
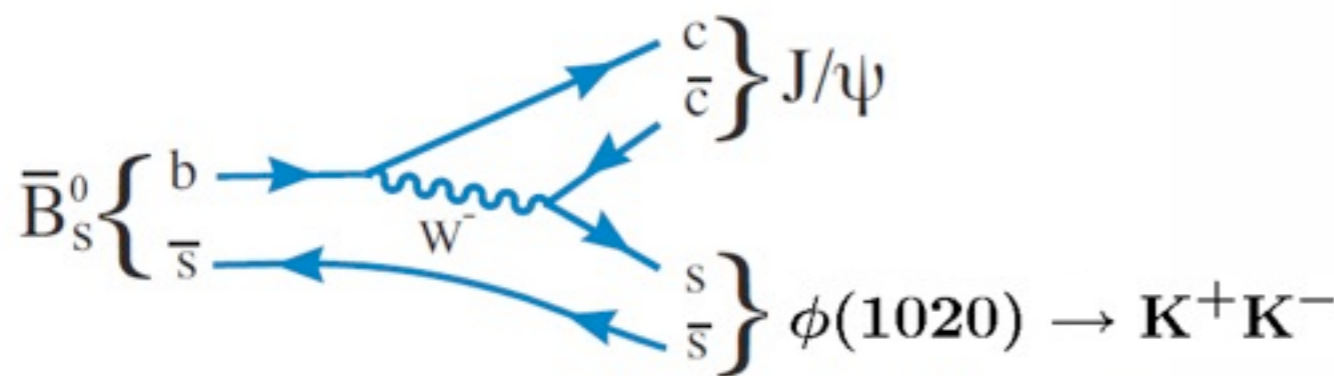
- result (1.0/fb, arXiv:1211.6093):

$$S_{J/\psi K_S^0} = 0.73 \pm 0.07 (\text{stat}) \pm 0.04 (\text{syst}),$$

- in good agreement with WA:  $\sin 2\beta^{\text{WA}} = 0.679 \pm 0.020$

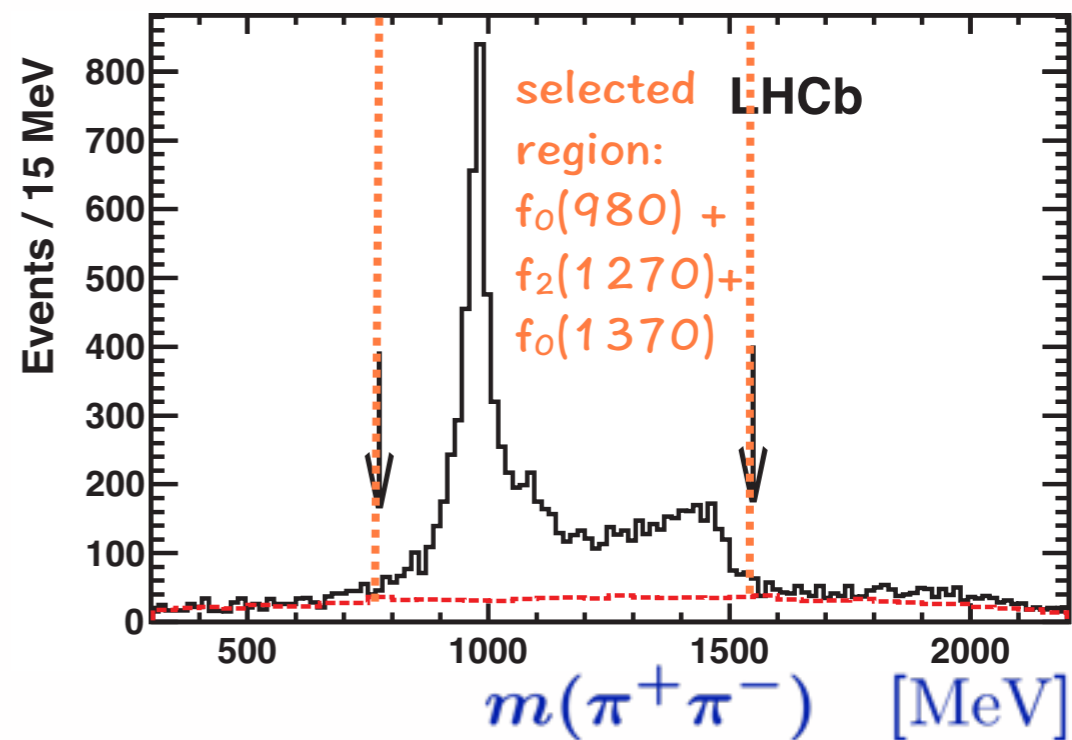
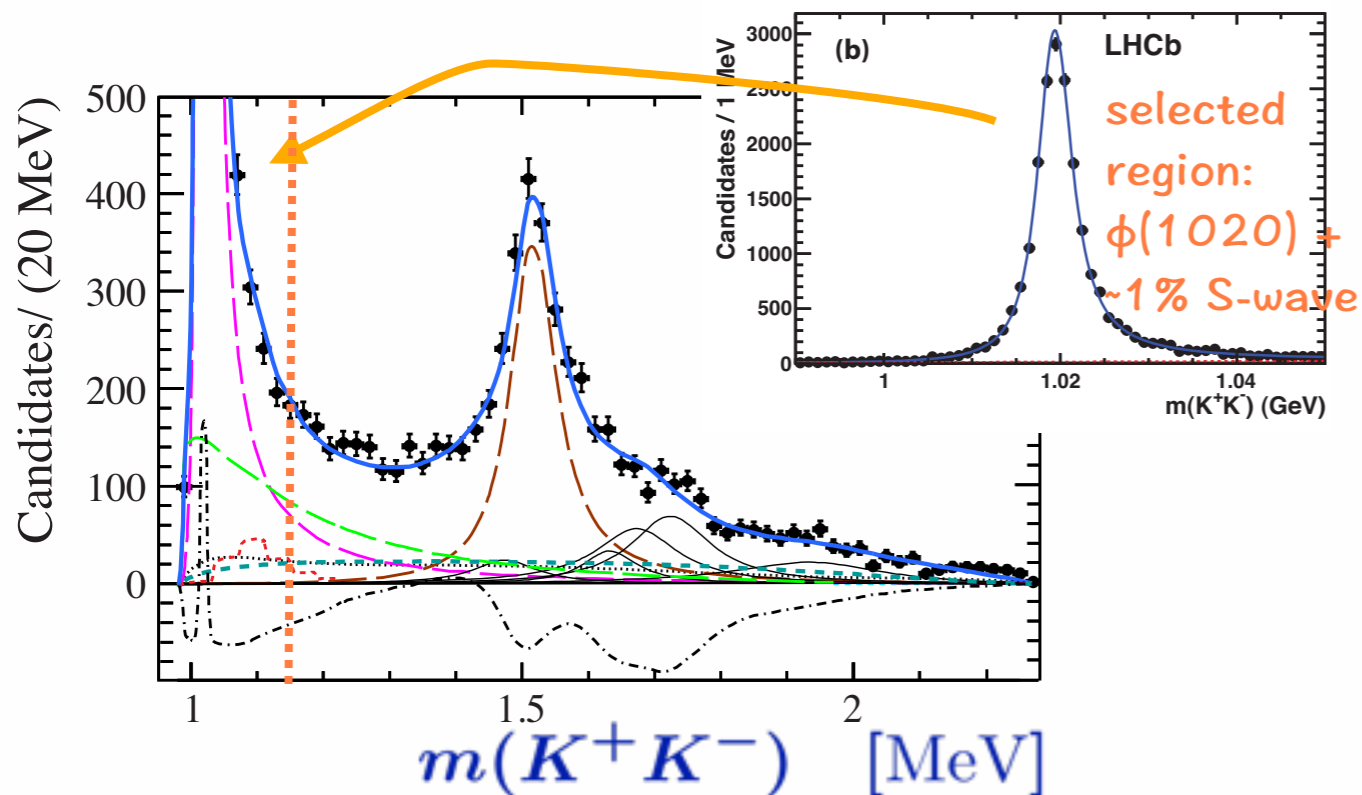
# $\phi_s$ at LHCb

- two most interesting modes:



- narrow  $\phi$  resonance --> clean
- vector-vector final state

- yield about 1/3 of  $B_s \rightarrow J/\psi \phi$
- vector-pseudo-scalar final state

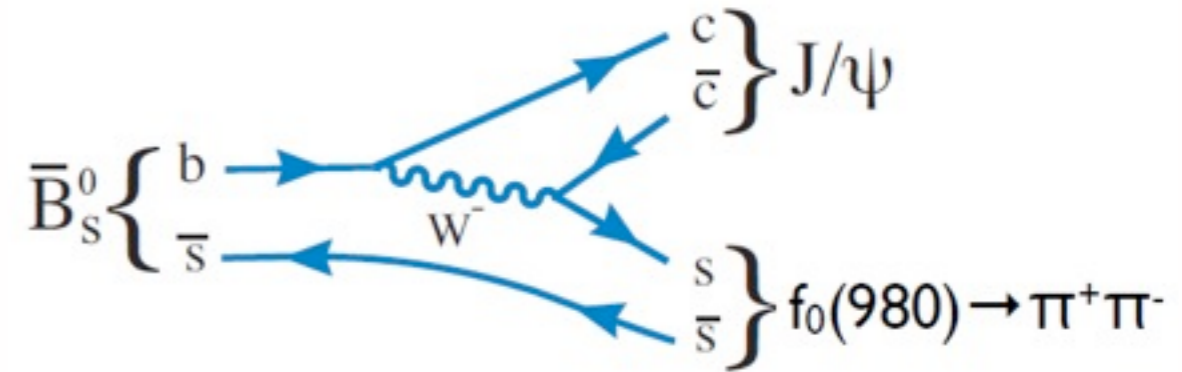
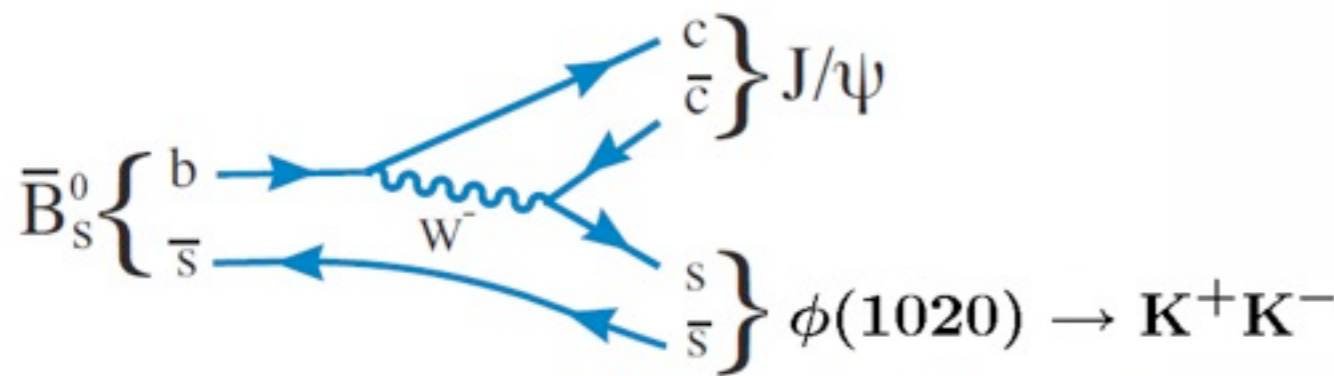


Dalitz analysis: LHCb-PAPER-2012-040,  
arXiv:1302.1213 (accepted by PRD)

Dalitz analysis: LHCb-PAPER-2012-005,  
PRD86(2012)052006

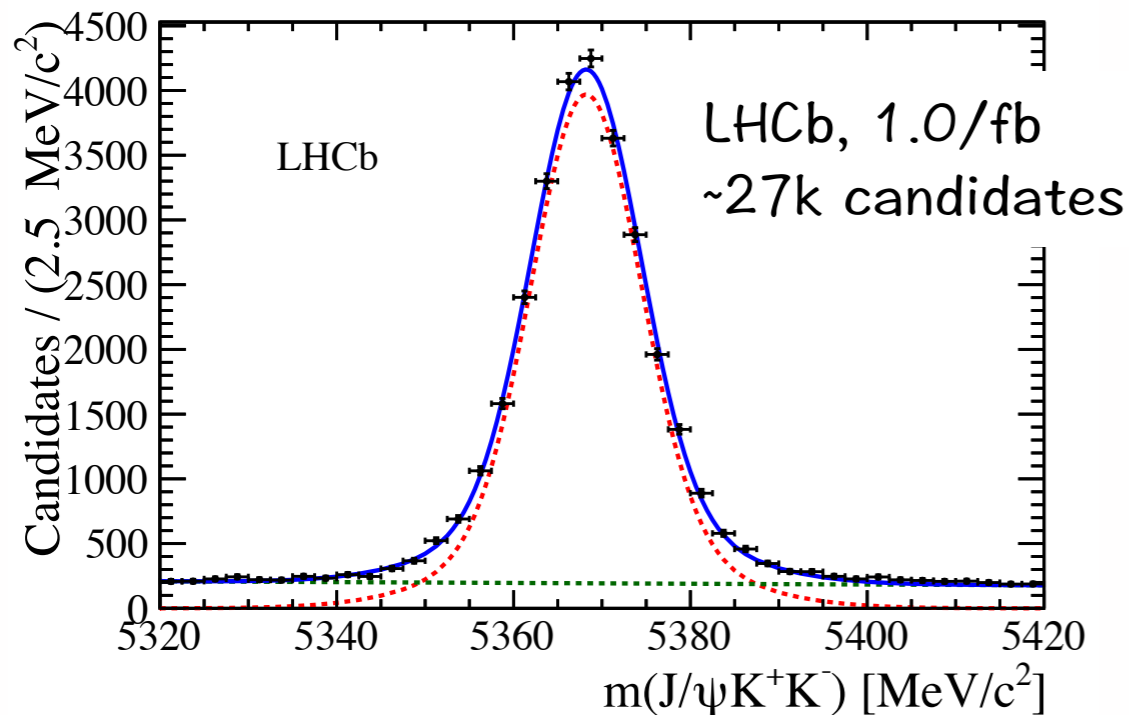
# $\phi_s$ at LHCb

- two most interesting modes:

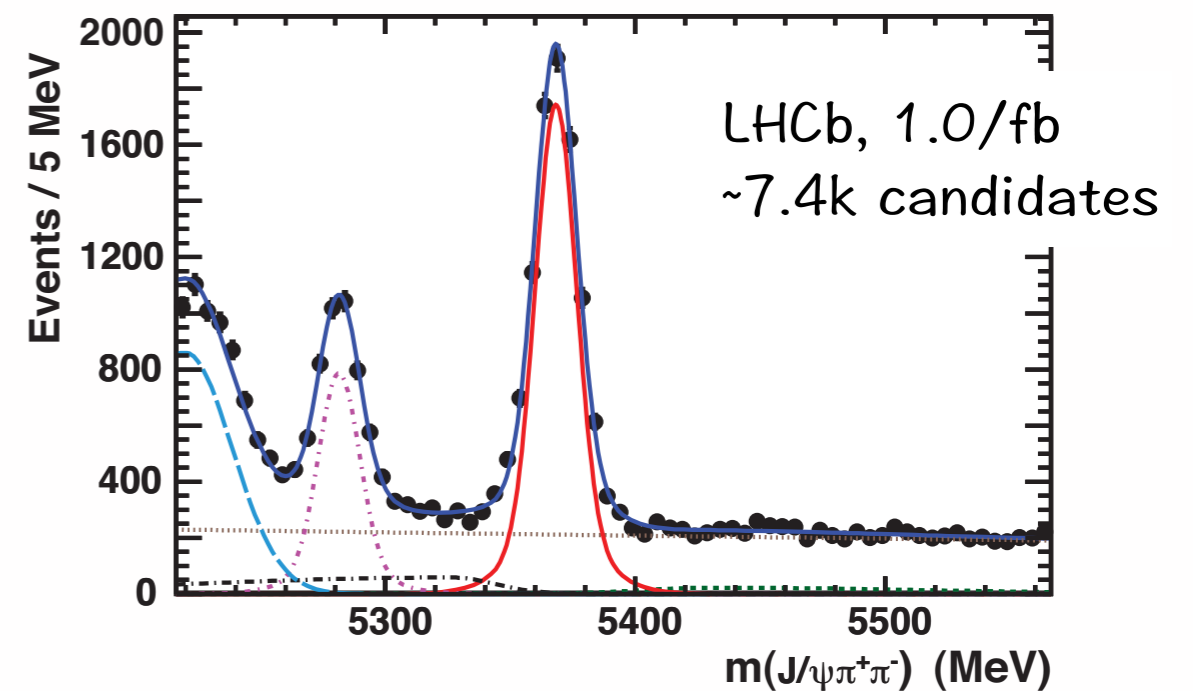


- narrow  $\phi$  resonance --> clean
- vector-vector final state

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- vector-pseudo-scalar final state



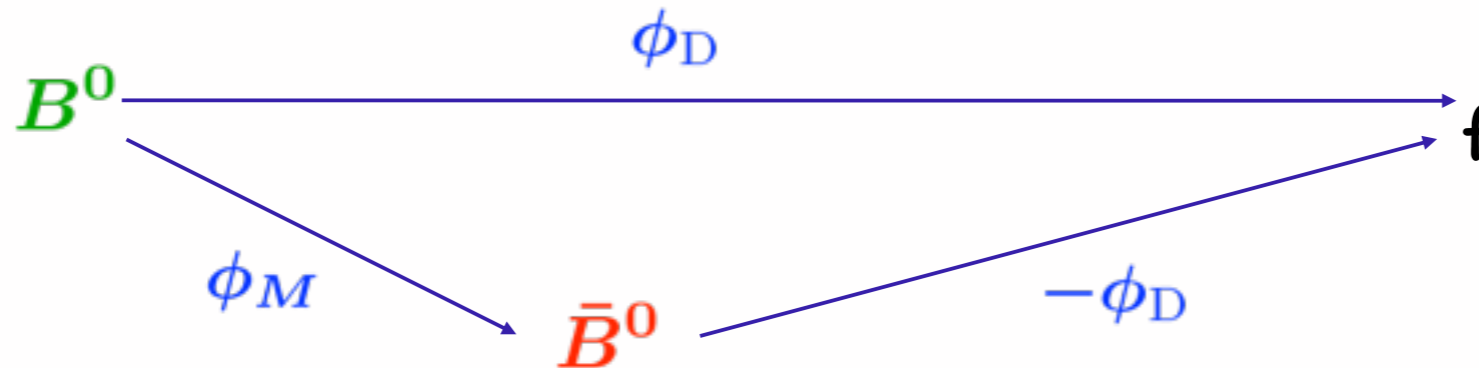
LHCb-PAPER-2013-002



LHCb-PAPER-2012-006,  
Phys. Lett. B713 (2012) 378

# Time-dependent CPV

- common final state  $f$  : mixing induced CPV



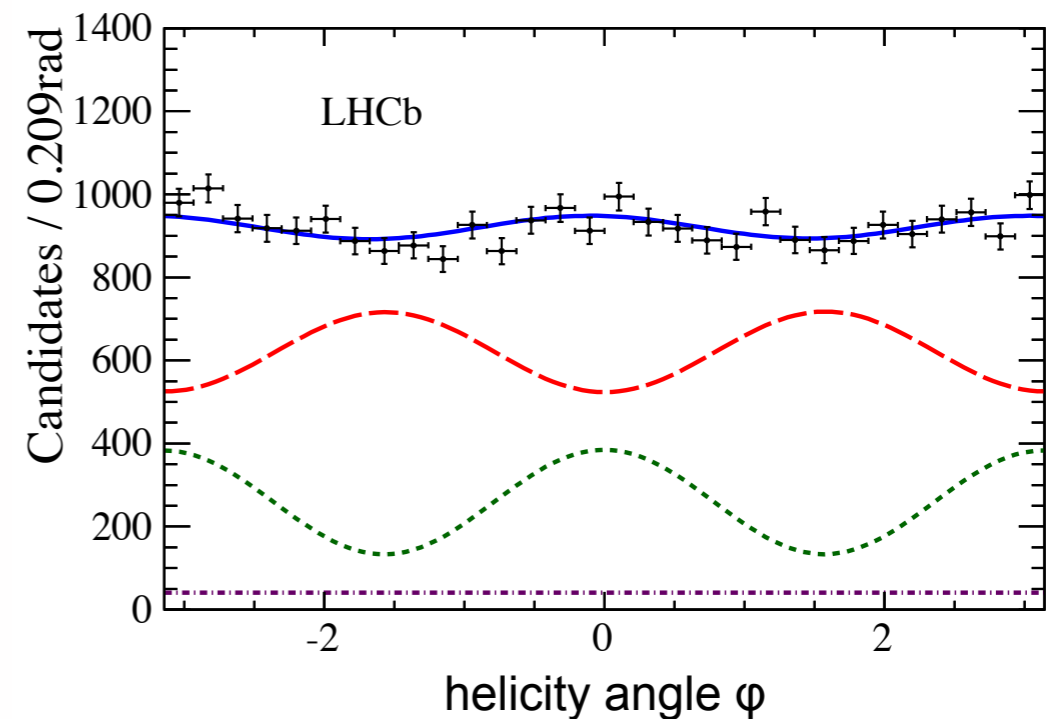
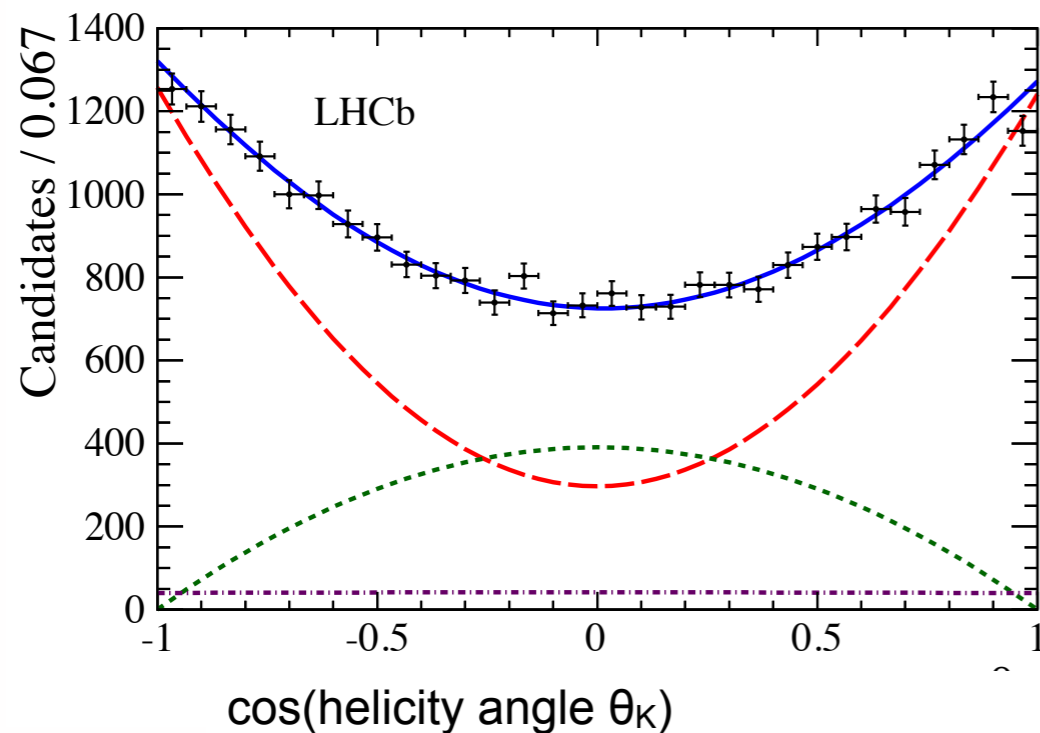
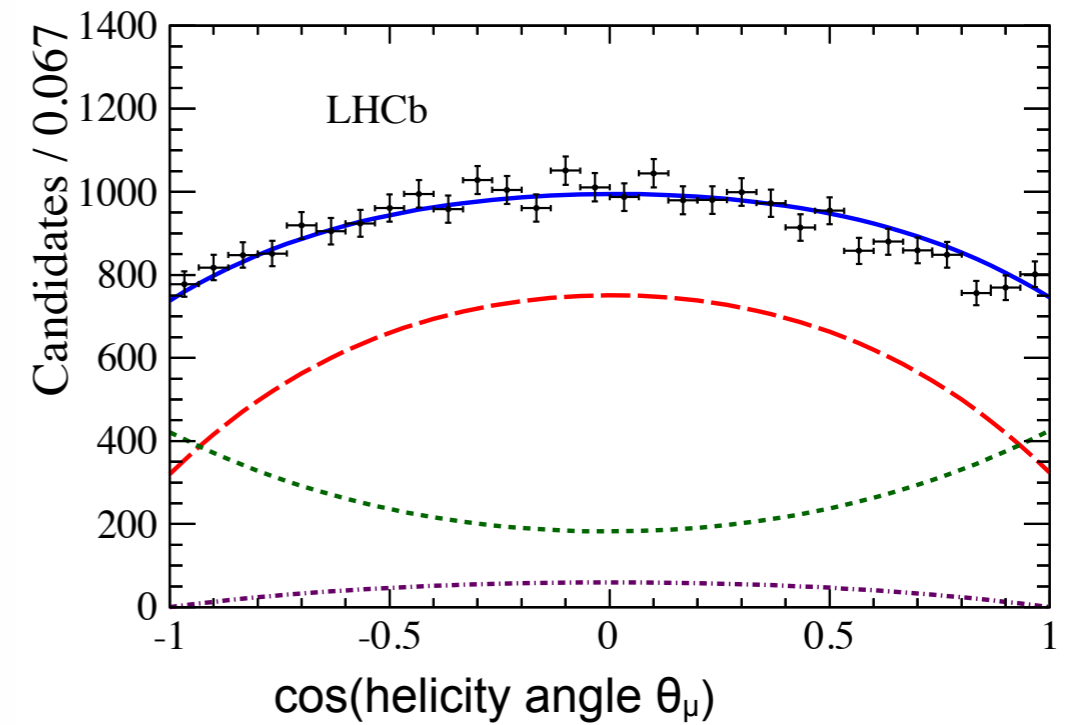
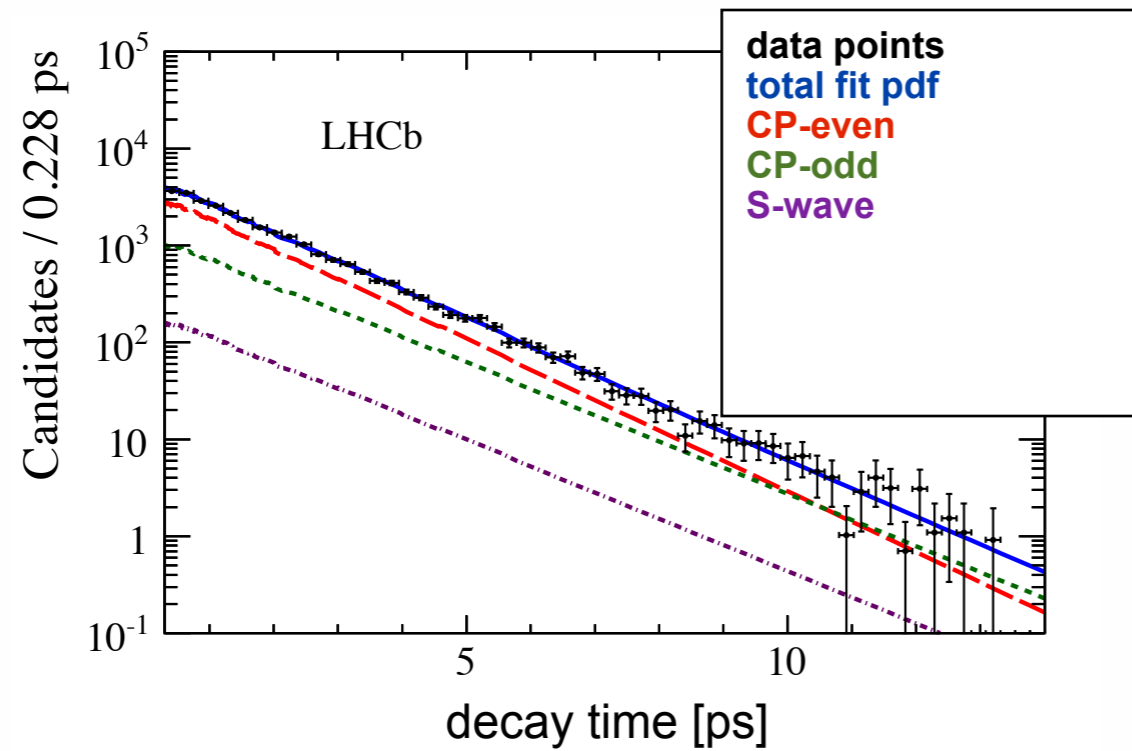
- if  $f$  is CP eigenstate, time dependent CP violation with pattern

$$A_{\text{CP}}(t) \equiv \frac{N(\bar{B} \rightarrow f) - N(B \rightarrow f)}{N(\bar{B} \rightarrow f) + N(B \rightarrow f)} = D \eta_f^{\text{CP}} \sin \phi_f \sin(\Delta m t)$$

- CP eigenvalue:  $\eta_f^{\text{CP}} = (-1)^L$ 
  - $B_s \rightarrow J/\psi f_{0,2}$ : CP-odd (see also LHCb-PAPER-2012-005)
  - $B_s \rightarrow J/\psi \phi$  : mixture of CP-odd and CP-even
    - needs 'angular analysis': include 4-body decay-angles in fit

- ambiguity:  $(\phi_s, \Delta\Gamma_s) \longleftrightarrow (\pi - \phi_s, -\Delta\Gamma_s)$

# MLL fit for $B_s \rightarrow J/\psi \phi$

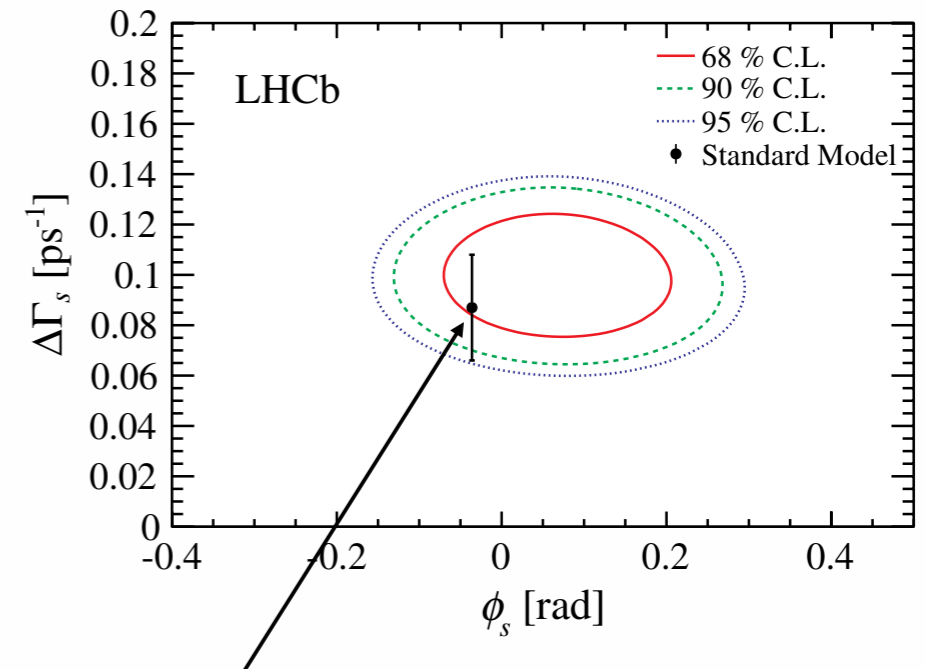




■ result for mixing parameters of J/psiKK fit

$$\begin{aligned} \phi_s &= 0.07 \pm 0.09 \text{ (stat)} \pm 0.01 \text{ (syst)} \text{ rad} \\ \Gamma_s &= 0.663 \pm 0.005 \text{ (stat)} \pm 0.006 \text{ (syst)} \text{ ps}^{-1} \\ \Delta\Gamma_s &= 0.100 \pm 0.016 \text{ (stat)} \pm 0.003 \text{ (syst)} \text{ ps}^{-1} \end{aligned}$$

dominant systematics: angular and decay time acceptance



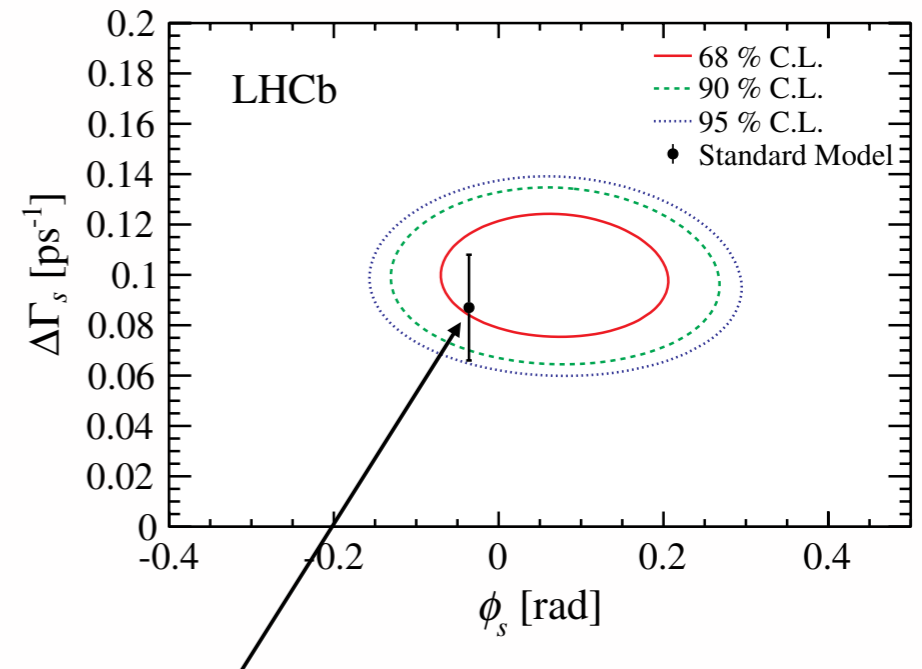
SM: Lenz and Nierste, arXiv: 1102.4274,  
CKM fitter, arXiv:1106.4041

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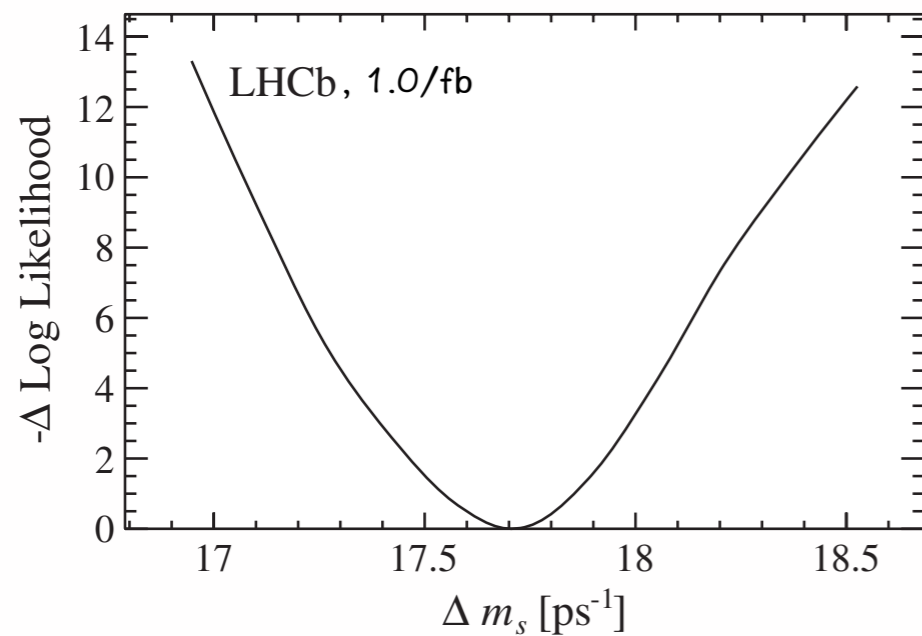
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dominant systematics: angular and decay time acceptance

- data now contains enough information to constrain also oscillation frequency



SM: Lenz and Nierste, arXiv: 1102.4274, CKM fitter, arXiv:1106.4041



$$\Delta m_s = 17.70 \pm 0.10 \text{ (stat)} \pm 0.01 \text{ (syst)} \text{ ps}^{-1}$$

in perfect agreement with  $B_s \rightarrow D_s \pi$

(and as good as  $\Delta m_s$  WA before start of LHC!)

## ■ result for mixing parameters of $J/\psi K K$ fit

$$\begin{aligned}\phi_s &= 0.07 \pm 0.09 \text{ (stat)} \pm 0.01 \text{ (syst)} \text{ rad} \\ \Gamma_s &= 0.663 \pm 0.005 \text{ (stat)} \pm 0.006 \text{ (syst)} \text{ ps}^{-1} \\ \Delta\Gamma_s &= 0.100 \pm 0.016 \text{ (stat)} \pm 0.003 \text{ (syst)} \text{ ps}^{-1}\end{aligned}$$

dominant systematics: angular and  
decay time acceptance

## ■ result for $J/\psi \pi^+ \pi^-$ fit

$$\phi_s = -0.14_{-0.16}^{+0.17} \pm 0.01 \text{ rad}$$

*Phys. Lett. B*713 (2012) 378  
+ update in PRD87,112010 (2013)

## ■ combination

$$\begin{aligned}\phi_s &= 0.01 \pm 0.07 \text{ (stat)} \pm 0.01 \text{ (syst)} \text{ rad}, \\ \Gamma_s &= 0.661 \pm 0.004 \text{ (stat)} \pm 0.006 \text{ (syst)} \text{ ps}^{-1}, \\ \Delta\Gamma_s &= 0.106 \pm 0.011 \text{ (stat)} \pm 0.007 \text{ (syst)} \text{ ps}^{-1}.\end{aligned}$$

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$$\begin{aligned} \phi_s &= 0.07 \pm 0.09 \text{ (stat)} \pm 0.01 \text{ (syst)} \text{ rad} \\ \Gamma_s &= 0.663 \pm 0.005 \text{ (stat)} \pm 0.006 \text{ (syst)} \text{ ps}^{-1} \\ \Delta\Gamma_s &= 0.100 \pm 0.016 \text{ (stat)} \pm 0.003 \text{ (syst)} \text{ ps}^{-1} \end{aligned}$$

dominant systematics: angular and decay time acceptance

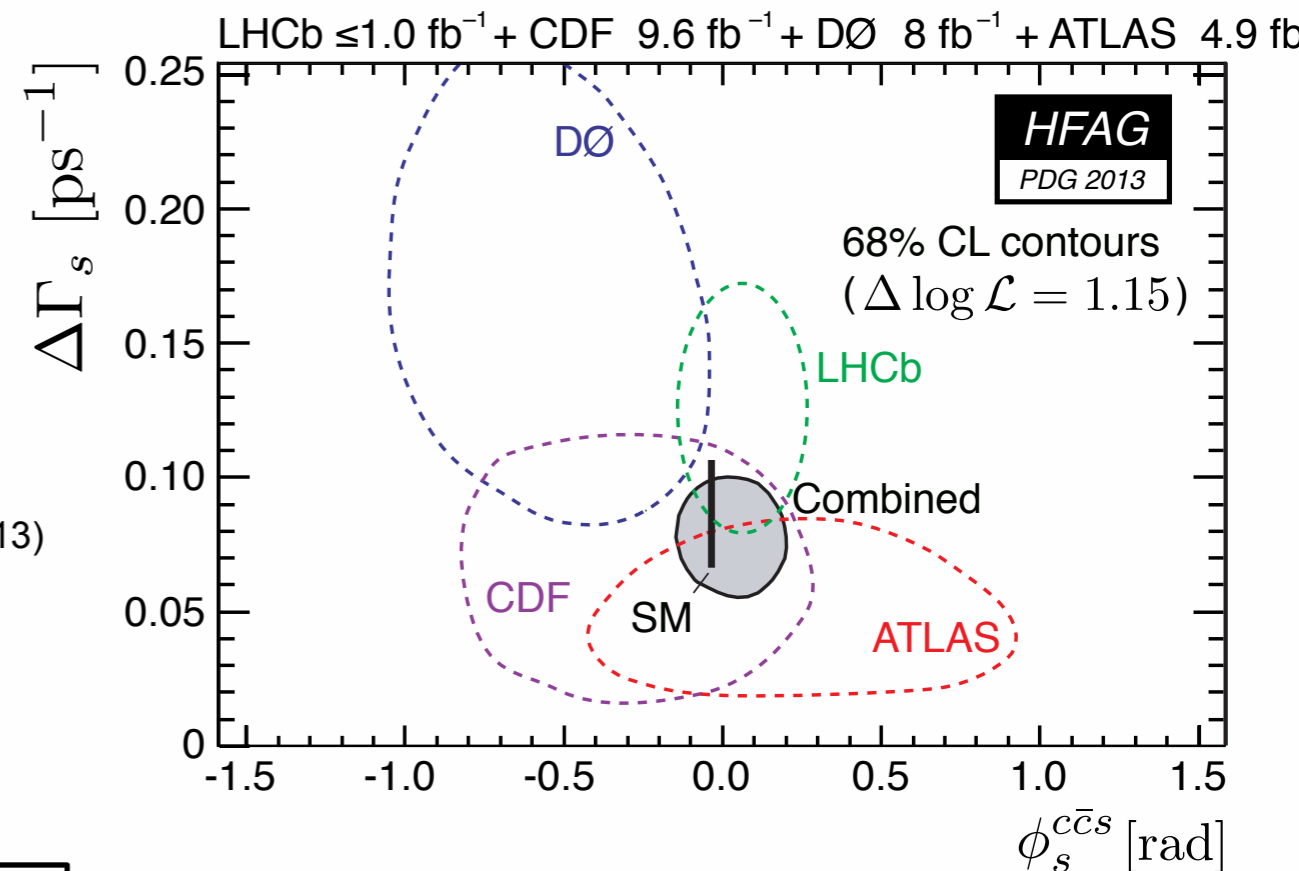
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*Phys. Lett. B*713 (2012) 378  
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# A closer look at lifetimes

---

- entangled states, with different lifetimes? you bet!

$$\tau_{KK}^{\text{eff}} = 1.455 \pm 0.046 \text{ (stat)} \pm 0.006 \text{ (syst)} \text{ ps}$$

arXiv:1207.5993, PLB 716 (2012) 393-400

$$\tau_{D_s^- D_s^+}^{\text{eff}} = 1.379 \pm 0.026 \text{ (stat)} \pm 0.017 \text{ (syst)} \text{ ps}$$

arXiv:1312.1217, subm. to PRL

$$\tau_{J/\psi \pi^+ \pi^-}^{\text{eff}} = 1.652 \pm 0.024 \text{ (stat)} \pm 0.024 \text{ (syst)} \text{ ps}$$

arXiv:1304.2600, PRD 87, 112010 (2013)

$$\tau_{J/\psi K_s^0}^{\text{eff}} = 1.75 \pm 0.12 \text{ (stat)} \pm 0.07 \text{ (syst)} \text{ ps}$$

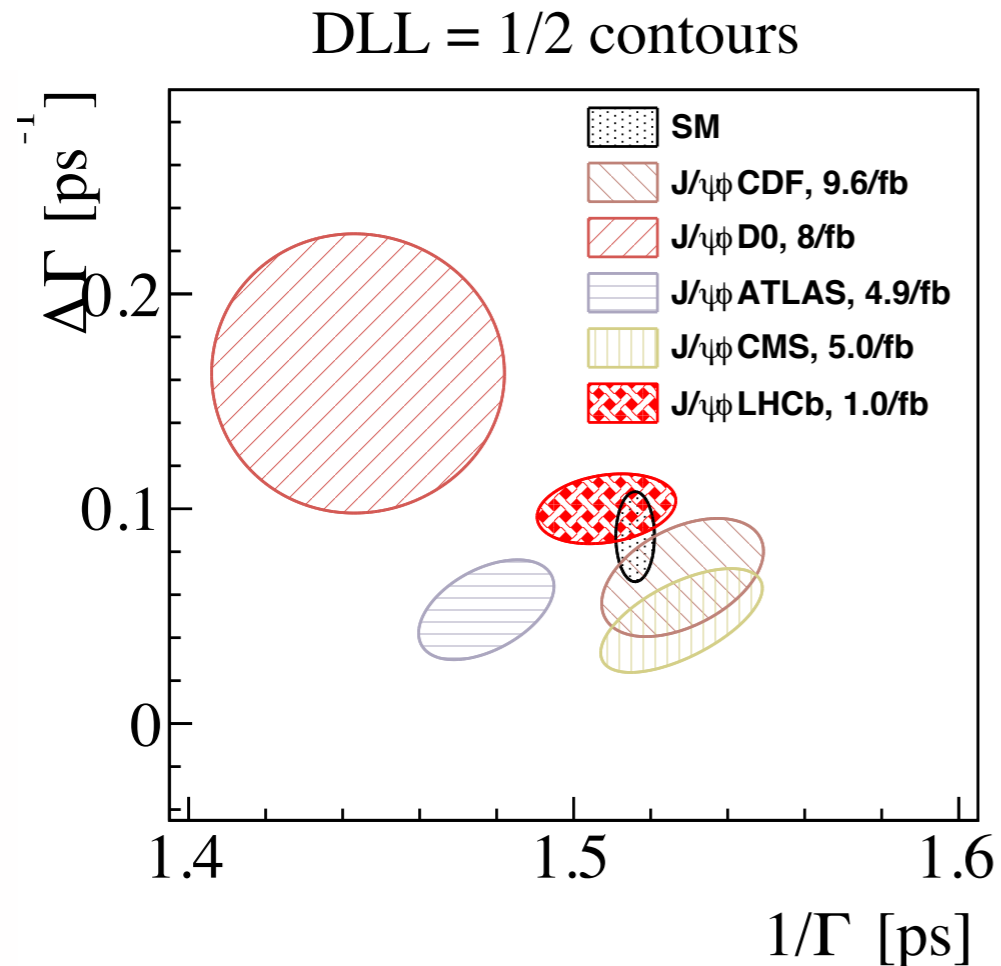
arXiv:1304.4500, Nucl. Phys. B 873 (2013) 275

- states with different CP-content are sensitive to different combinations of  $\Gamma$  and  $\Delta\Gamma$  (and  $\phi_s$ )

(see also Fleischer and Knegjens, Eur.Phys.J. C71 (2011) 1789)



# A closer look at lifetimes



CDF: Phys. Rev. Lett 109, 171802 (2012)

D0: Phys. Rev. D **85**, 032006 (2012),

ATLAS: [JHEP 12 \(2012\) 072](#)

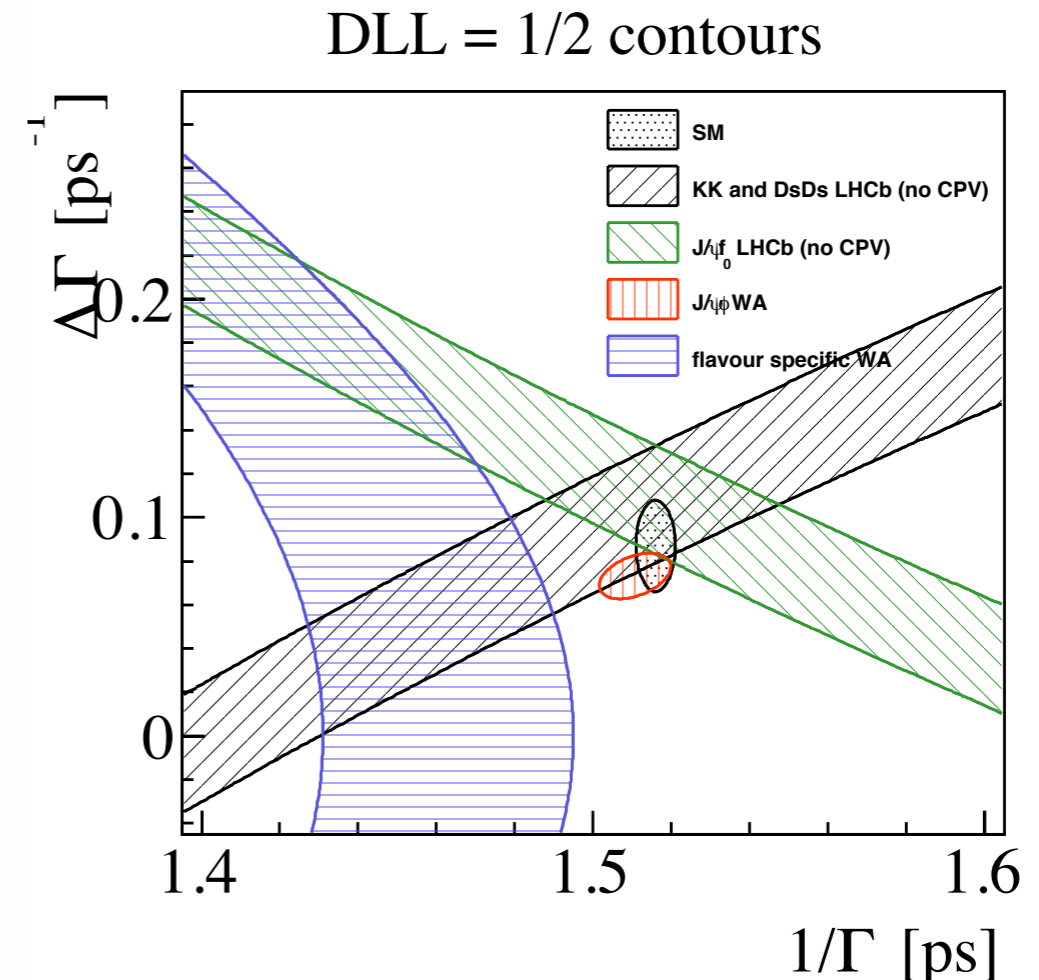
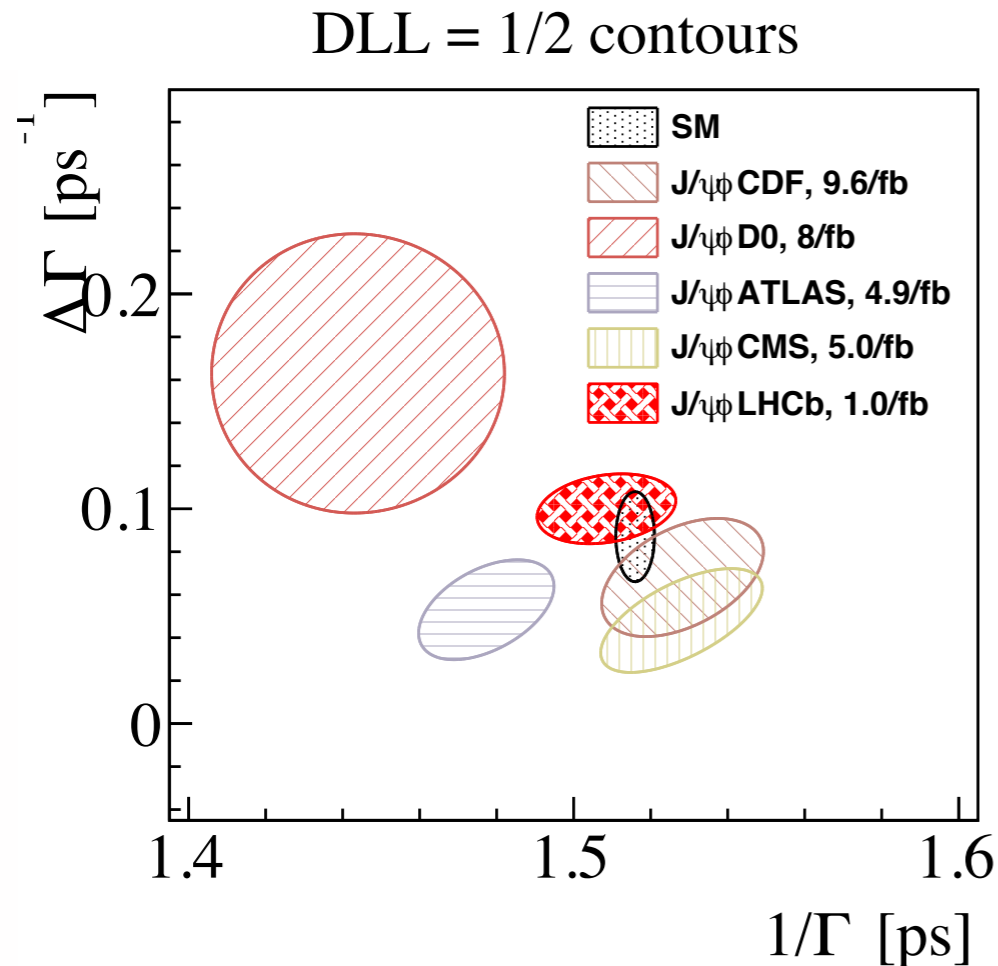
CMS: CMS-PAS-BPH-11-006

latest SM prediction: Lenz and Nierste, arXiv: 1102.4274

(  $\Gamma_s$  prediction uses  $\tau_{B_d}^{\text{WA}}$  and  $-4 \cdot 10^{-4} < \frac{\tau_{B_s}}{\tau_{B_d}} - 1 < 0$  )

- everything in agreement with mixing model and SM+HQE predictions

# A closer look at lifetimes



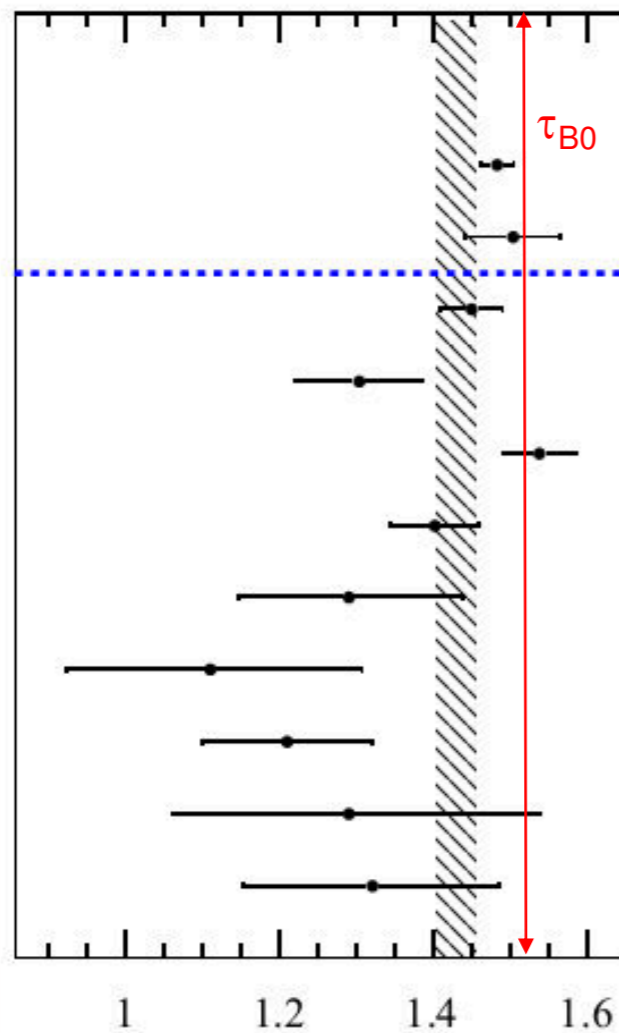
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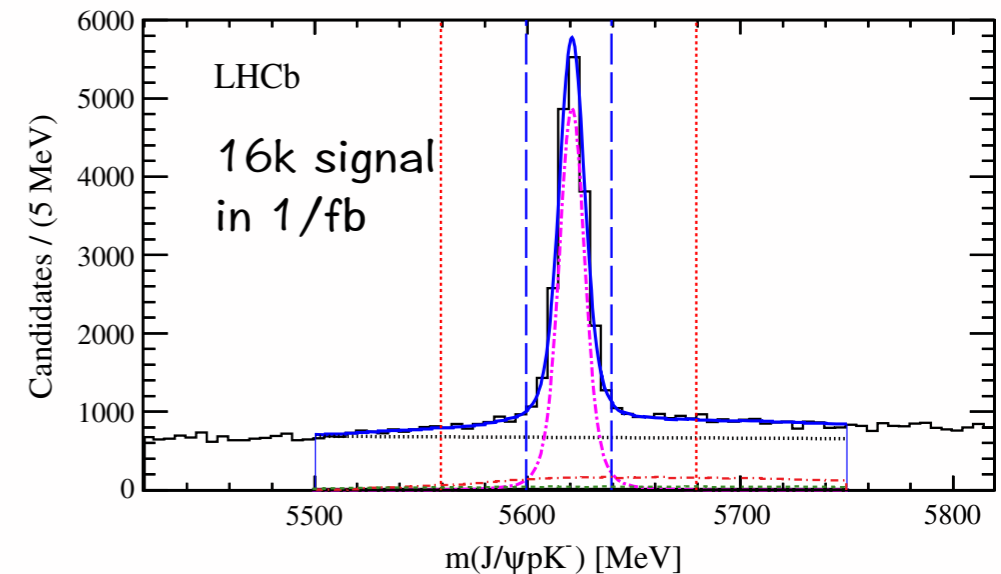
# $\Lambda_b$ baryon lifetime

- clean prediction from HQE:  $\tau(\Lambda_b) / \tau(B^0) = 1 \pm \text{few } \%$
- long standing controversy: experiments measured 'too small'  $\tau(\Lambda_b)$



Experiment

- LHCb (1/fb) (2013) [ $J/\psi p K$ ]
- CMS (2012) [ $J/\psi \Lambda$ ]
- ATLAS (2012) [ $J/\psi \Lambda$ ]
- D0 (2012) [ $J/\psi \Lambda$ ]
- CDF (2011) [ $J/\psi \Lambda$ ]
- CDF (2010) [ $\Lambda_c^+ \pi$ ]
- D0 (2007) [ $J/\psi \Lambda$ ]
- DLPH (1999) [Semileptonic decay]
- ALEP (1998) [Semileptonic decay]
- OPAL (1998) [Semileptonic decay]
- CDF (1996) [Semileptonic decay]



- most recent LHCb measurement, using 16k  $\Lambda_b \rightarrow J/\psi p K$  in 1/fb:
- ... in good agreement with expectation

$$\frac{\tau_{\Lambda_b^0}}{\tau_{B^0}} = 0.976 \pm 0.012 \pm 0.006$$

# CPV in mixing

■ CPV in mixing:  $a_{sl} \equiv 1 - \left| \frac{q}{p} \right|^2 = \frac{|\Gamma_{12}|}{|M_{12}|} \sin \phi_{12}$

■ very small in SM:  $a_{sl}^{\text{SM}}(B_d) = (-4.1 \pm 0.6) \cdot 10^{-4}$  and  $a_{sl}^{\text{SM}}(B_s) = (1.9 \pm 0.3) \cdot 10^{-5}$   
(Lenz and Nierste, 2011)

■ LHCb: measure time-integrated rate asymmetry in  $B_s \rightarrow D_s^- X \mu^+ \nu$

$$A = \frac{N(D_s^- \mu^+) - N(D_s^+ \mu^-)}{N(D_s^- \mu^+) + N(D_s^+ \mu^-)} = \frac{1}{2} a_{sl} - (a_{\text{prod}} - \frac{1}{2} a_{sl}) \frac{\int_0^\infty dt e^{-\Gamma_s t} \cos(\Delta m_s t)}{\int_0^\infty dt e^{-\Gamma_s t} \cosh(\frac{1}{2} \Delta \Gamma_s t)}$$

production asymmetry: ~0.01

dilution factor from oscillations ~0.2%

# CPV in mixing

■ CPV in mixing:  $a_{sl} \equiv 1 - \left| \frac{q}{p} \right|^2 = \frac{|\Gamma_{12}|}{|M_{12}|} \sin \phi_{12}$

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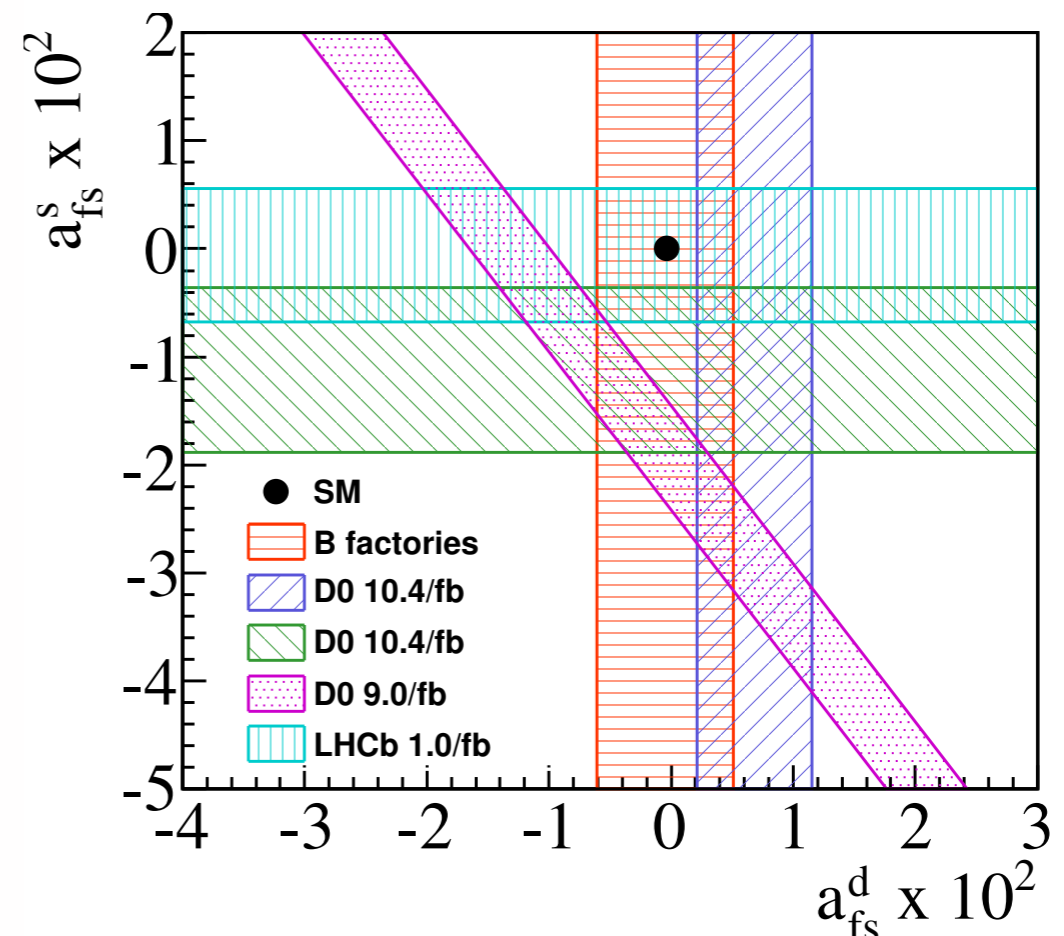
$$A = \frac{N(D_s^- \mu^+) - N(D_s^+ \mu^-)}{N(D_s^- \mu^+) + N(D_s^+ \mu^-)} = \frac{1}{2} a_{sl} - (a_{\text{prod}} - \frac{1}{2} a_{sl}) \frac{\int_0^\infty dt e^{-\Gamma_s t} \cos(\Delta m_s t)}{\int_0^\infty dt e^{-\Gamma_s t} \cosh(\frac{1}{2} \Delta \Gamma_s t)}$$

■ result (1.0/fb, [arxiv:1308:1048](https://arxiv.org/abs/1308.1048), subm. to PLB)

$$a_{sl}^s = (-0.06 \pm 0.50 \pm 0.36)\%$$

current systematics dominated by statistical uncertainty on efficiency ratio

DLL = 1/2 contours





# $B_s \rightarrow \mu\mu$ and $B_d \rightarrow \mu\mu$

Very rare decays in SM:

- no tree-level FCNC
- helicity suppression
- CKM suppression

... none of which is necessarily true in NP!

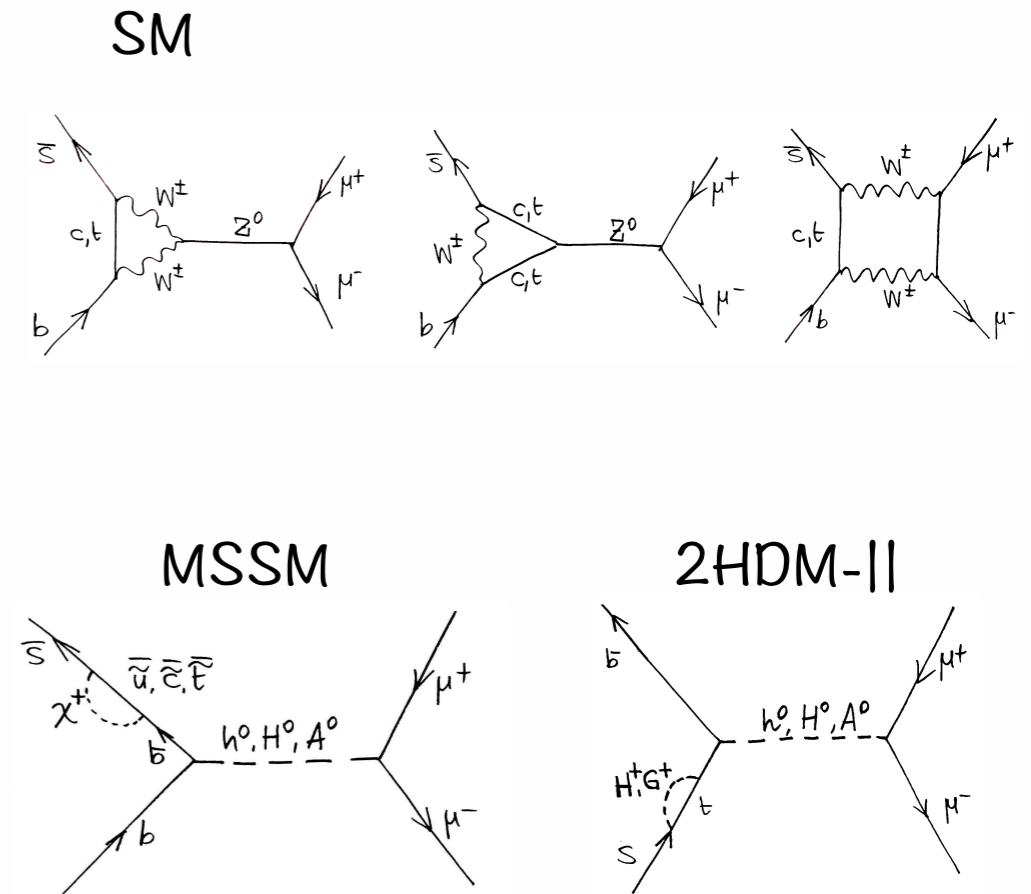
Precise SM predictions

- $\text{Br}(B_d \rightarrow \mu\mu) = (1.1 \pm 0.2) 10^{-10}$
- $\text{Br}(B_s \rightarrow \mu\mu) = (3.5 \pm 0.2) 10^{-9}$

Buras et al, EPJ C72 (2012) 2172; see also PRL109 (2012) 041801

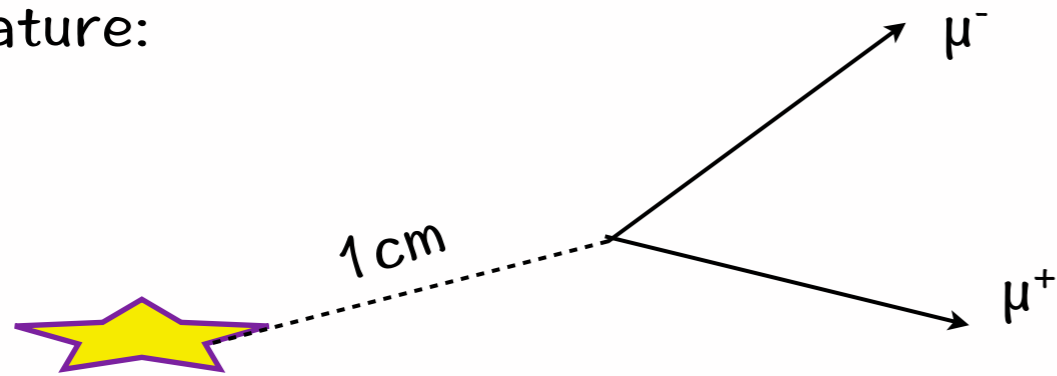
Strong enhancement in e.g. MSSM:  $\text{Br} \propto \tan^6\beta$

Previously known as the “golden channel” for NP discovery ... more recently named the “SUSY killer”!

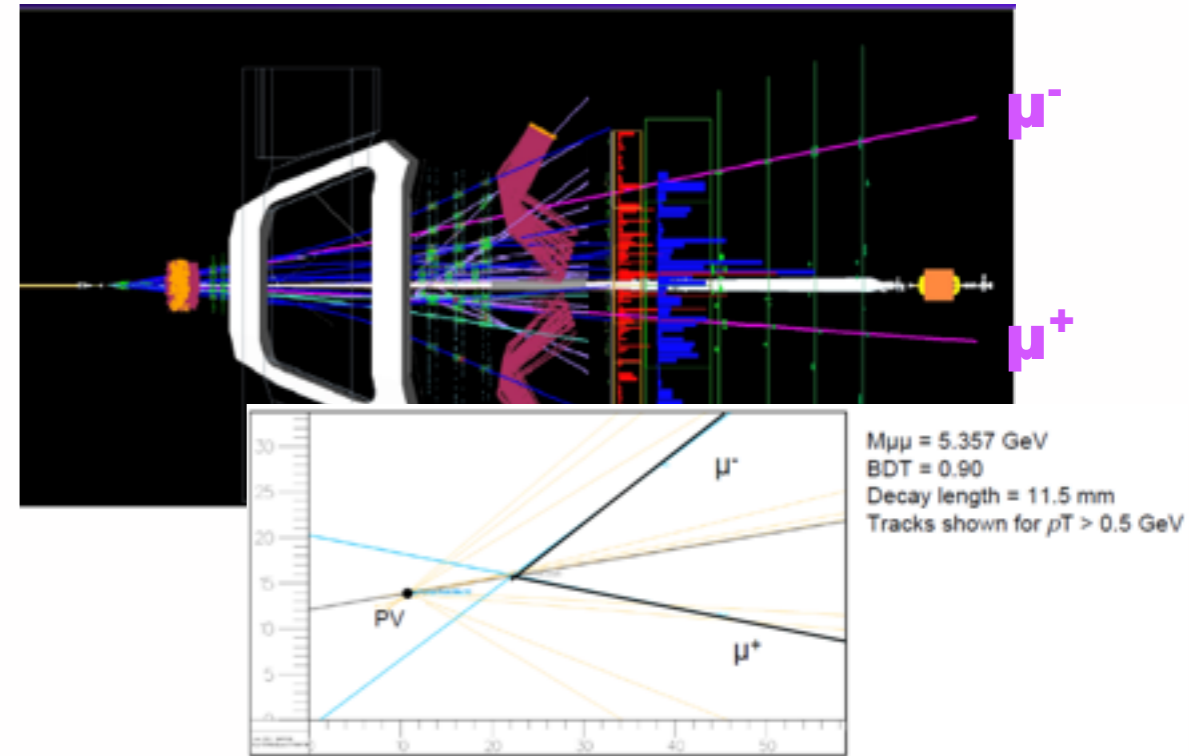


# The quest for $B_q \rightarrow \mu\mu$

Signature:

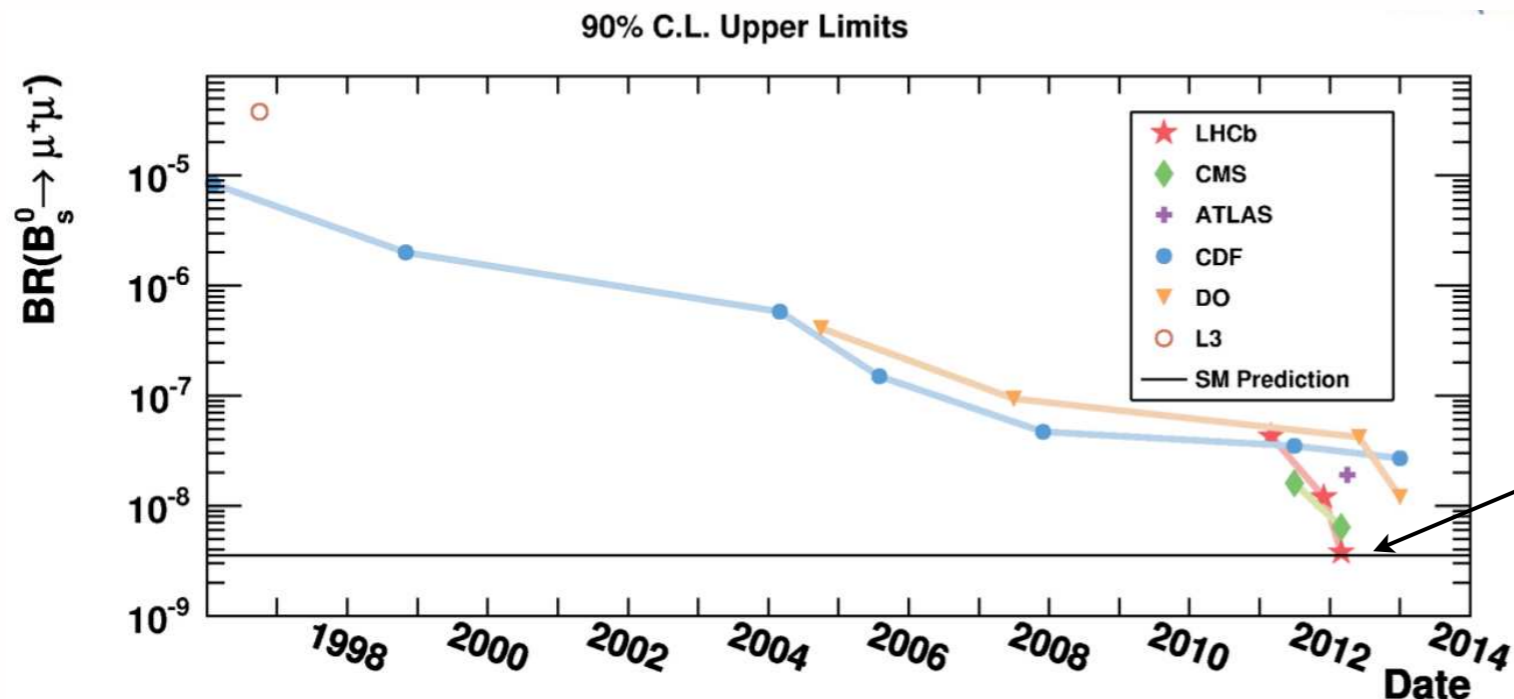


Actual candidate in LHCb data



Complicated measurement: large backgrounds from  $B \rightarrow hh$  and double  $B \rightarrow X\mu$

Searches started as early as 1984 (CLEO)!  
Evolution of limit as function of time:

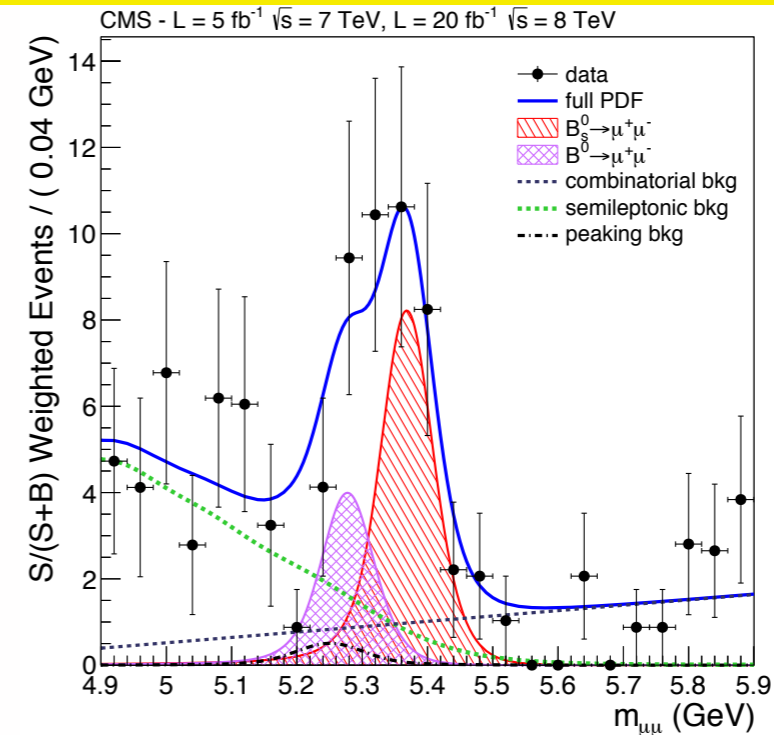
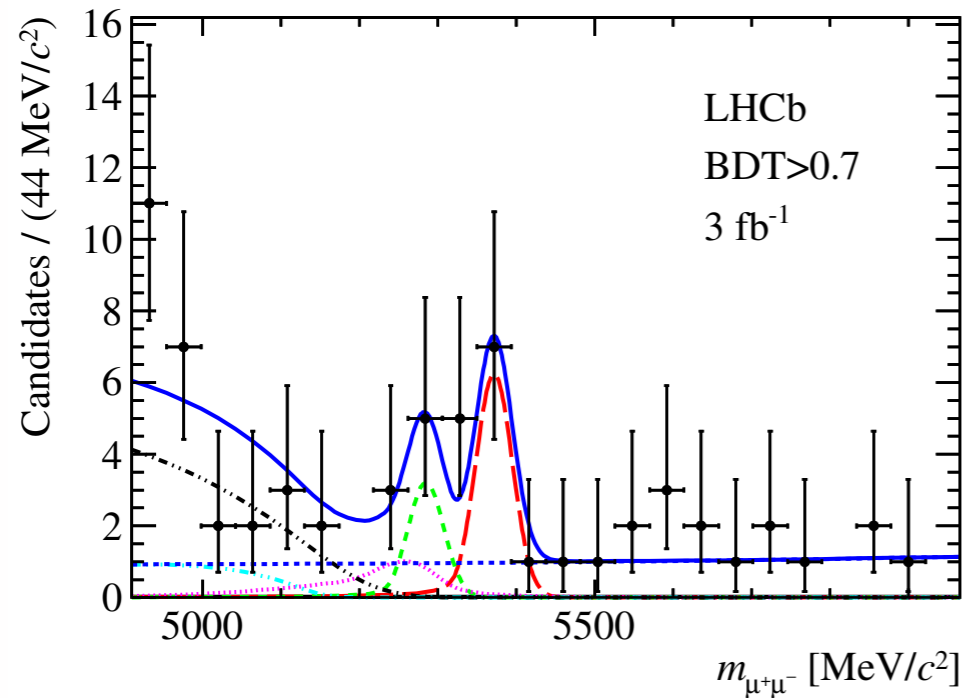


first hints by LHCb and CMS

# EPS2013: LHCb/CMS joint discovery of $B_s \rightarrow \mu^+ \mu^-$

LHCb: arXiv:1307.5024, PRL.111.101805 (2013)

CMS: arXiv:1307.5025, PRL. 111.101804 (2013)



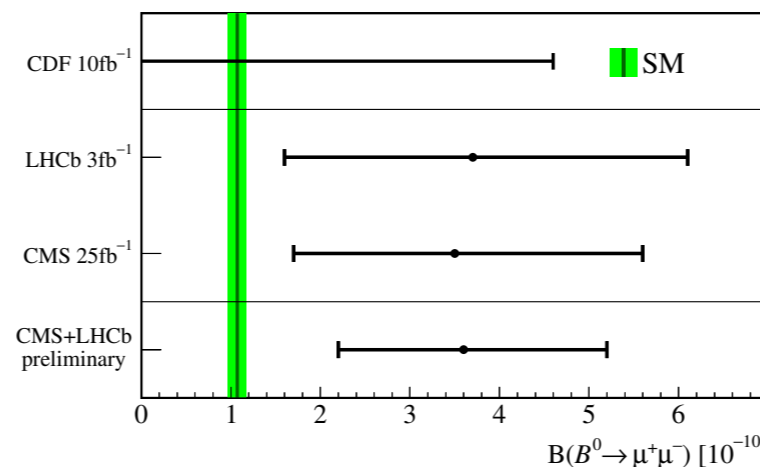
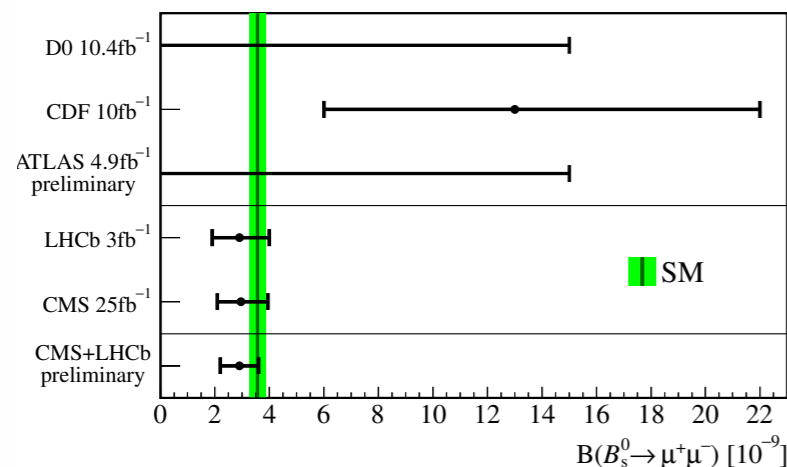
$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.9_{-1.0}^{+1.1}) \times 10^{-9}, \quad \text{--> } 4.0\sigma$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) = (3.7_{-2.1}^{+2.4}) \times 10^{-10}$$

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (3.0_{-0.9}^{+1.0}) \times 10^{-9}, \quad \text{--> } 4.3\sigma$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) = (3.5_{-1.8}^{+2.1}) \times 10^{-10}$$

combination:  $\text{BR}(B_s \rightarrow \mu^+ \mu^-) = (2.9 \pm 0.7) \times 10^{-9}$   
 $\text{BR}(B^0 \rightarrow \mu^+ \mu^-) = 3.6_{-1.4}^{+1.6} \times 10^{-10}$



CMS-PAS-BPH-13-007 ;  
LHCb-CONF-2013-012

# Other weak decays to $\mu\mu$ ...

$D^0 \rightarrow \mu\mu$

arXiv:1305.5059, PLB 725 (2013) 15-24

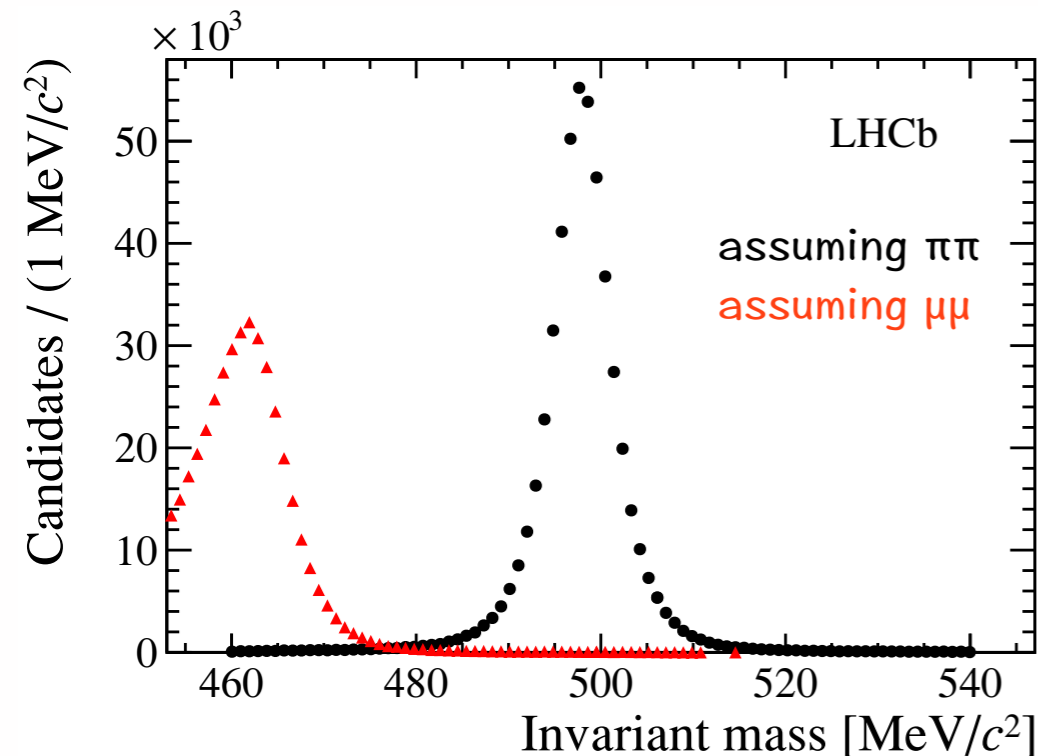
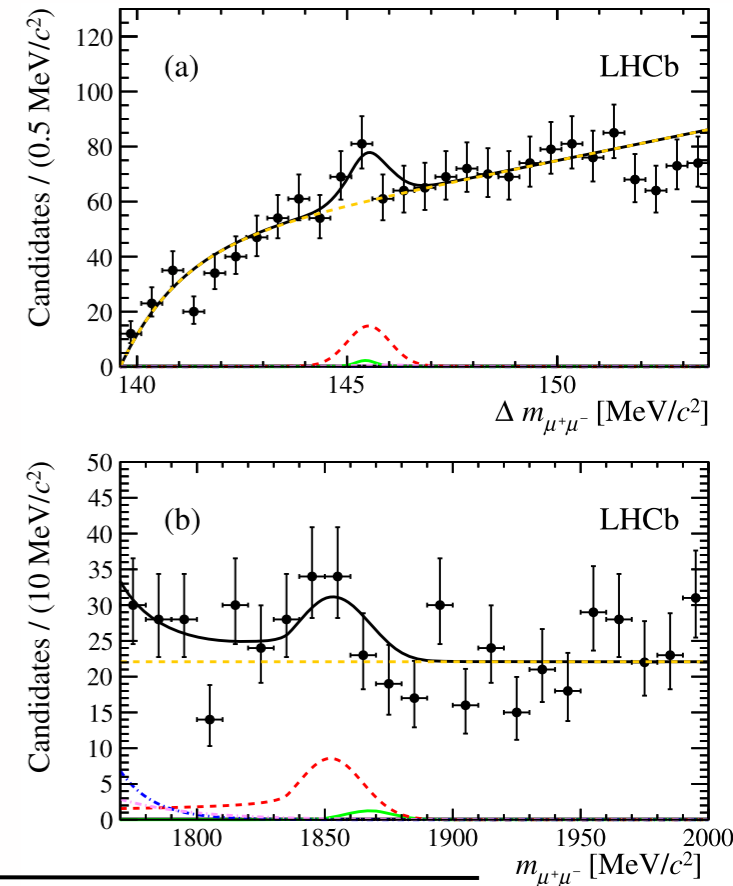
SM:

short-distance:  $\sim < 10^{-18}$

long-distance:  $< 6 \times 10^{-11}$

LHCb, 0.9/fb:  $\mathcal{B}(D^0 \rightarrow \mu^+ \mu^-) < 6.2 (7.6) \times 10^{-9}$   
at 90 (95)% CL

combin.  
 $D^0 \rightarrow K\pi$   
 $D^0 \rightarrow \pi\pi$   
 $D^0 \rightarrow \mu\mu$



$K_S \rightarrow \mu\mu$

1/fb: arXiv:1209.4029, JHEP 01 (2013) 090

SM:  $\mathcal{B}(K_S^0 \rightarrow \mu^+ \mu^-) = (5.0 \pm 1.5) \times 10^{-12}$

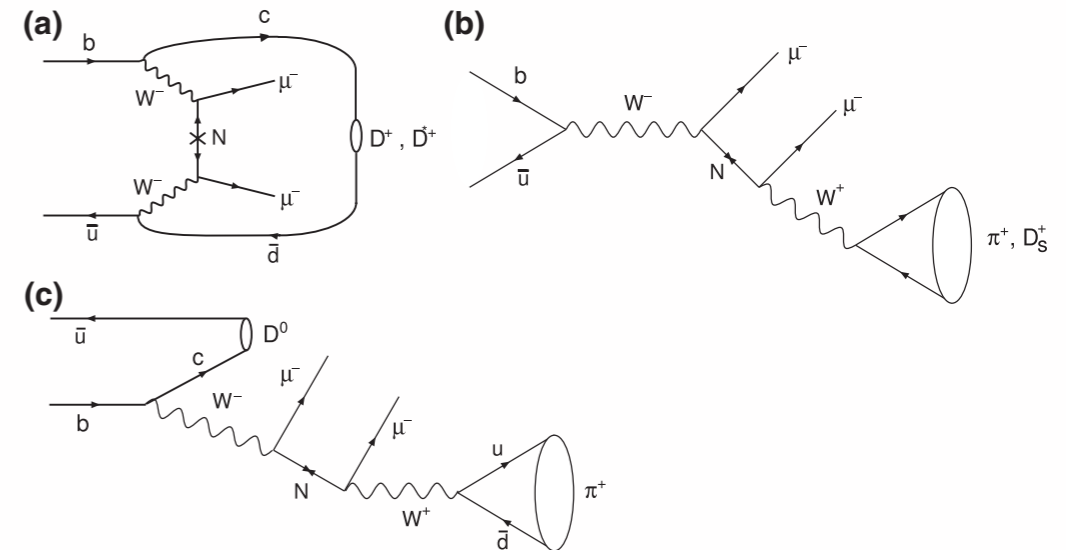
1/fb, number  $K_S$  in acceptance:  $\sim 10^{13}$

LHCb, 1/fb:  $\mathcal{B}(K_S^0 \rightarrow \mu^+ \mu^-) < 11(9) \times 10^{-9}$   
at 95 (90)% CL

# And more exotic signals with two muons ...

- B final states with like-sign muons are sensitive to 'medium-heavy' majorana neutrino N
- 5 decays analyzes in LHCb

$$\left. \begin{aligned}
 B^- &\rightarrow D^+ \mu^- \mu^- \\
 B^- &\rightarrow D^{*+} \mu^- \mu^- \\
 B^- &\rightarrow \pi^+ \mu^- \mu^- \\
 B^- &\rightarrow D_s^{*+} \mu^- \mu^- \\
 B^- &\rightarrow D^0 \pi^+ \mu^- \mu^-
 \end{aligned} \right\} \begin{array}{l} \text{via virtual N} \\ \text{also via on-shell N} \end{array}$$



- results, using 0.4/fb of LHCb data

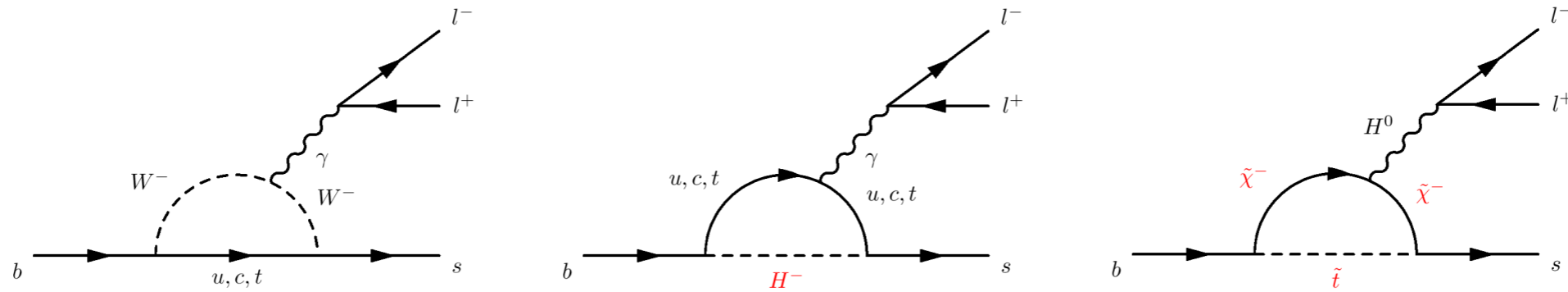
Mode	$\mathcal{B}$ upper limit	Approximate limits as function of $M_N$
$D^+ \mu^- \mu^-$	$6.9 \times 10^{-7}$	
$D^{*+} \mu^- \mu^-$	$2.4 \times 10^{-6}$	
$\pi^+ \mu^- \mu^-$	$1.3 \times 10^{-8}$	$(0.4 - 1.0) \times 10^{-8}$
$D_s^+ \mu^- \mu^-$	$5.8 \times 10^{-7}$	$(1.5 - 8.0) \times 10^{-7}$
$D^0 \pi^+ \mu^- \mu^-$	$1.5 \times 10^{-6}$	$(0.3 - 1.5) \times 10^{-6}$

nothing seen ...

... yet!

# B $\rightarrow$ K(\*) $\mu^+ \mu^-$

- another example if a b  $\rightarrow$  s transition (FCNC): very sensitive to NP



- many observables rates, angular distributions, CP/isospin asymmetries
  - clean theoretical predictions
  - clean experimental signature

- theoretical framework: Operator Product Expansion

$$\left( H_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \left[ \underbrace{C_i(\mu) O_i(\mu)}_{\text{left-handed part}} + \underbrace{C'_i(\mu) O'_i(\mu)}_{\text{right-handed part suppressed in SM}} \right] \right)$$

$i = 1, 2$	Tree
$i = 3 - 6, 8$	Gluon penguin
$i = 7$	Photon penguin
$i = 9, 10$	Electroweak penguin
$i = S$	Higgs (scalar) penguin
$i = P$	Pseudoscalar penguin

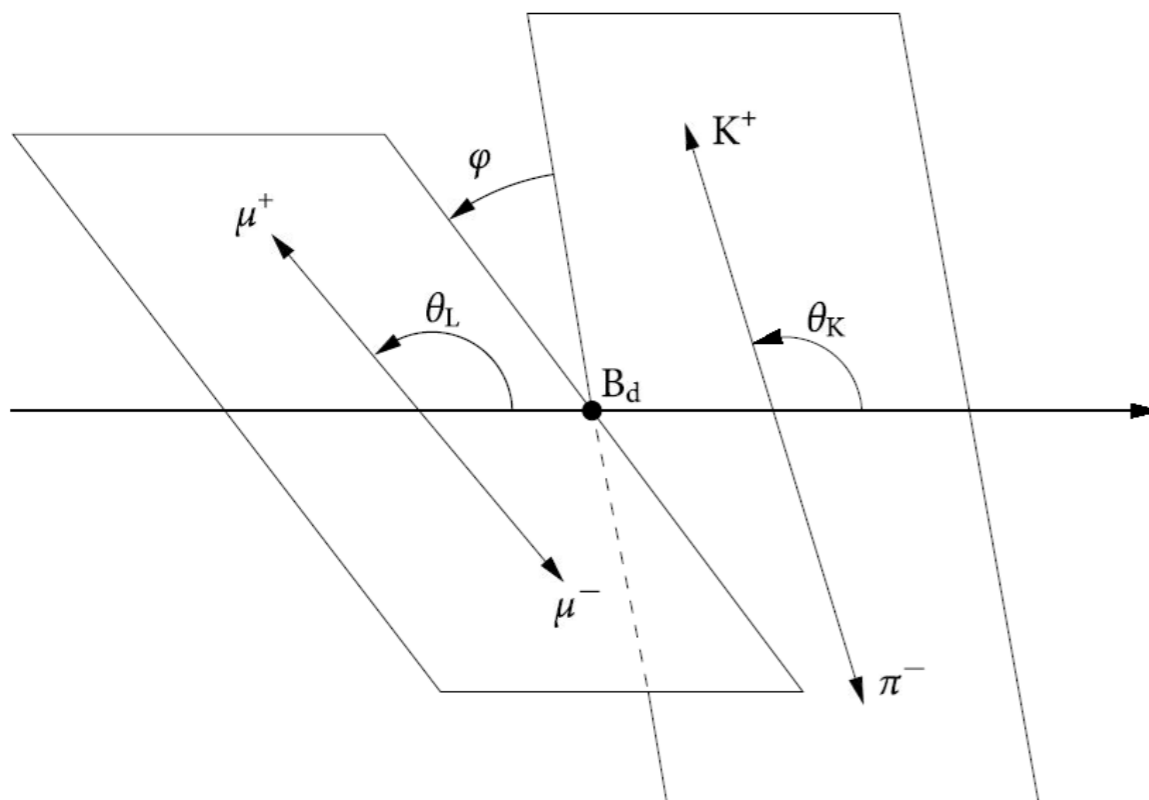


# Angular analysis of $B_d \rightarrow K^{*0} \mu^+ \mu^-$

generic parameterization of CP-averaged angular distribution:

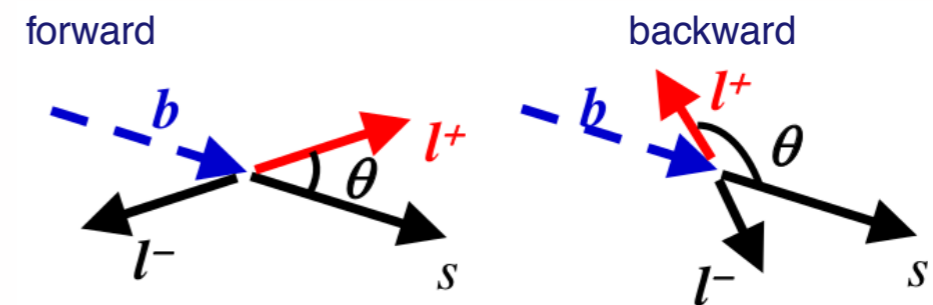
$$\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\theta_l d\theta_K d\varphi dq^2} = \frac{9}{32\pi} \left[ \frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_l \right. \\ \left. - F_L \cos^2 \theta_K \cos 2\theta_l + S_3 \sin^2 \theta_K \sin^2 \theta_l \cos 2\varphi \right. \\ \left. + S_4 \sin 2\theta_K \sin 2\theta_l \cos \varphi + S_5 \sin 2\theta_K \sin \theta_l \cos \varphi \right. \\ \left. + S_6^s \sin^2 \theta_K \cos \theta_l + S_7 \sin 2\theta_K \sin \theta_l \sin \varphi \right. \\ \left. + S_8 \sin 2\theta_K \sin 2\theta_l \sin \varphi + S_9 \sin^2 \theta_K \sin^2 \theta_l \sin 2\varphi \right]$$

$q^2 = \text{dimuon invariant mass}^2$



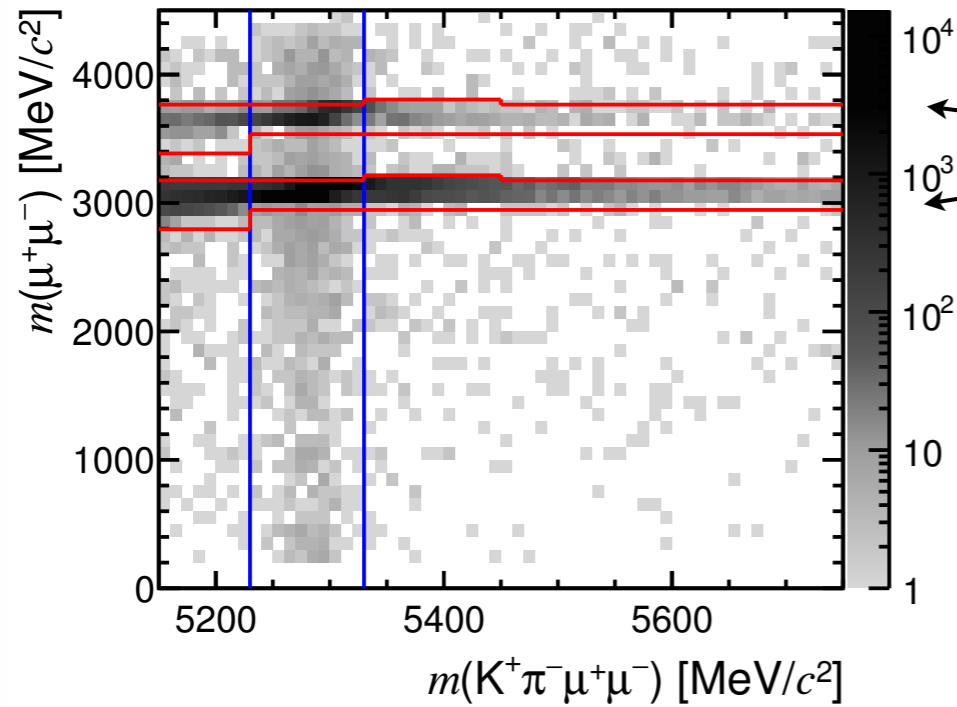
Most well known observables:

- ▶ muon forward backward asymmetry  $A_{FB} = 4/3 S_6^s$



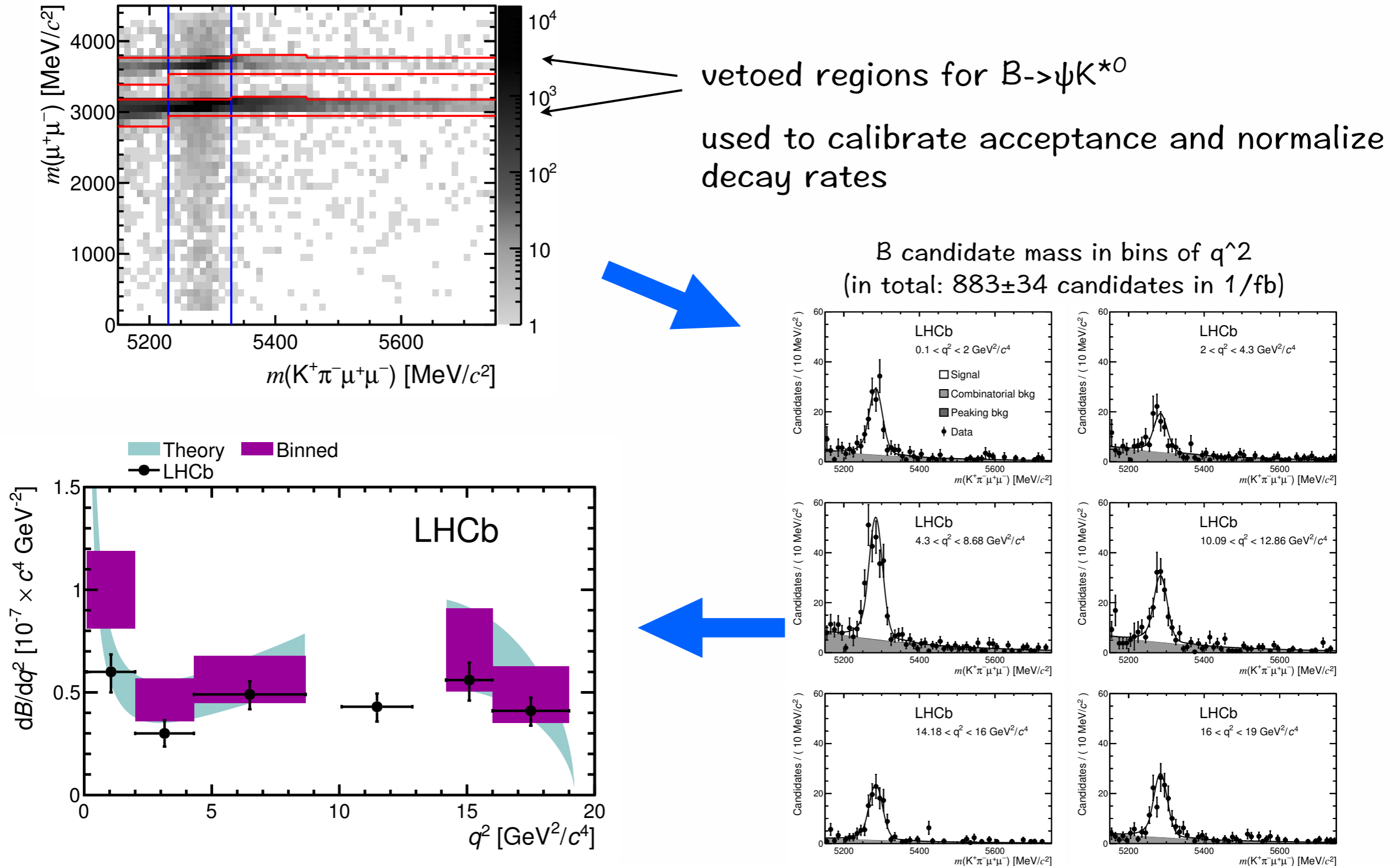
- ▶  $K^*$  longitudinal polarization fraction,  $F_L$

# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ at LHCb



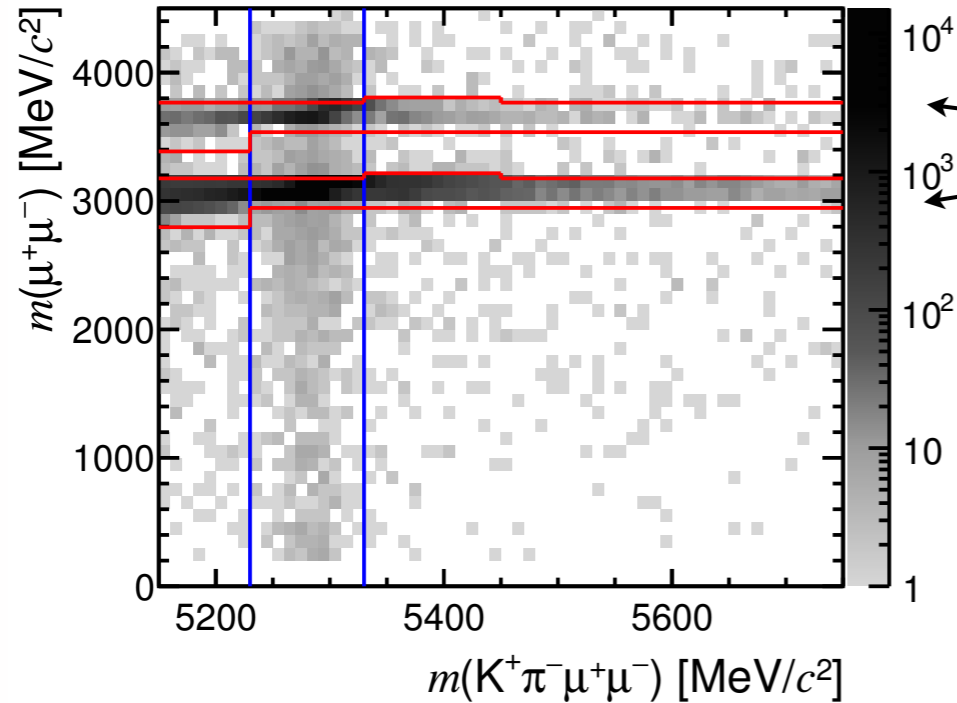
vetoed regions for  $B \rightarrow \psi K^{*0}$   
used to calibrate acceptance and normalize decay rates

# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ at LHCb



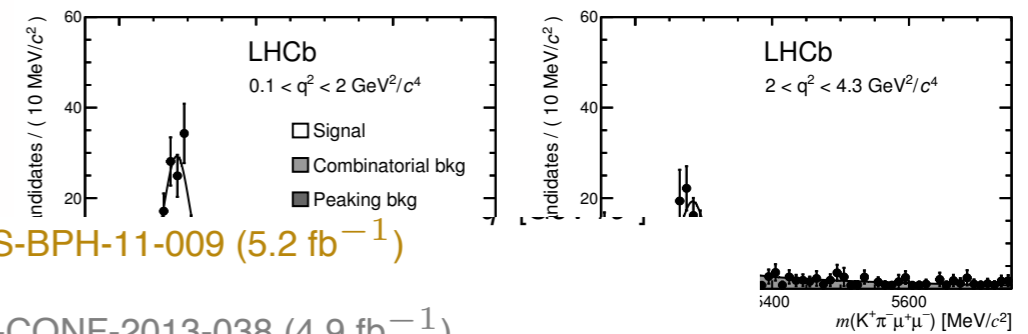
good agreement between data and predictions?

# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ at LHCb



vetoed regions for  $B \rightarrow \psi K^{*0}$   
used to calibrate acceptance and normalize decay rates

B candidate mass in bins of  $q^2$   
(in total:  $883 \pm 34$  candidates in  $1/\text{fb}$ )



CMS: CMS-PAS-BPH-11-009 ( $5.2 \text{ fb}^{-1}$ )

ATLAS: ATLAS-CONF-2013-038 ( $4.9 \text{ fb}^{-1}$ )

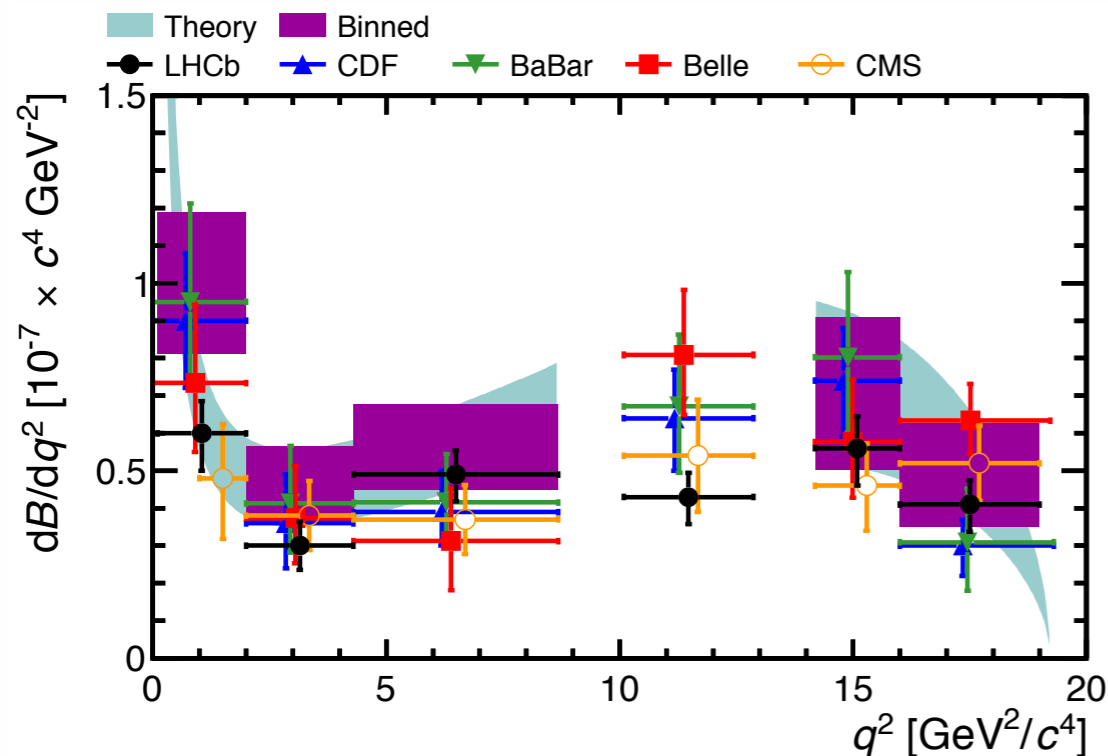
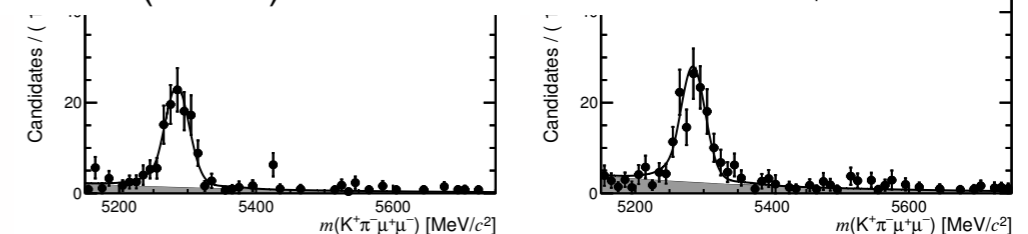
BELLE: Phys. Rev. Lett. 103 (2009) 171801 ( $605 \text{ fb}^{-1}$ )

BABAR: Phys. Rev. D73 (2006) 092001 ( $208 \text{ fb}^{-1}$ )

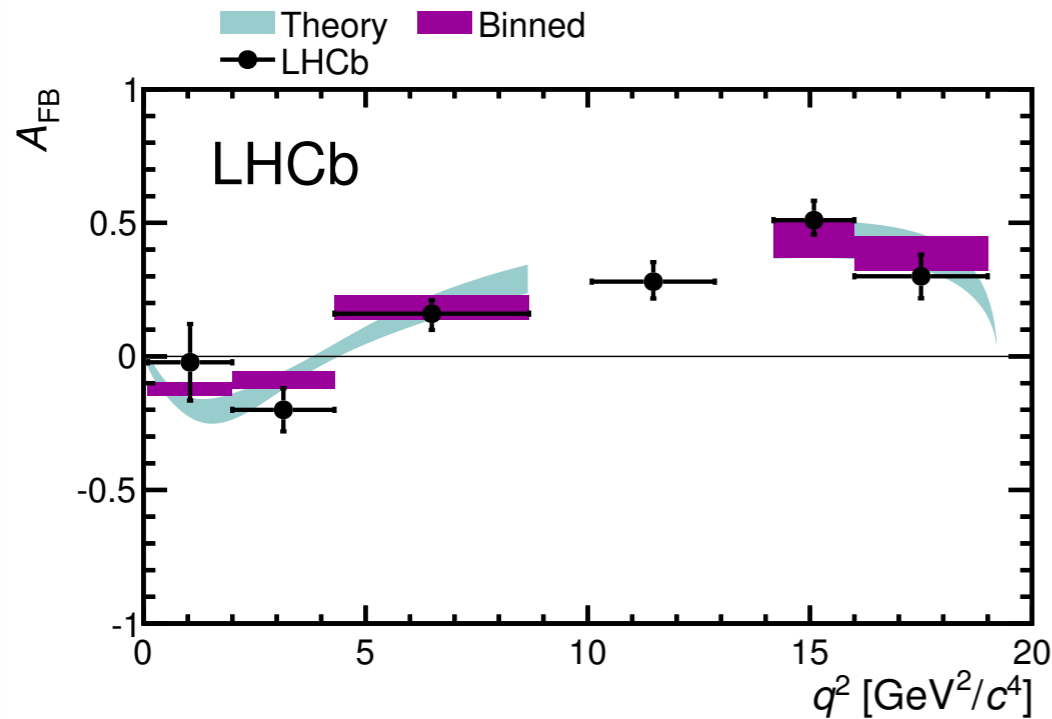
CDF: Phys. Rev. Lett 108 (2012) 081807 ( $6.8 \text{ fb}^{-1}$ )

(results from CDF Public Note 10894 ( $9.6 \text{ fb}^{-1}$ ) not included)

LHCb: arXiv:1304.6325 ( $1 \text{ fb}^{-1}$ )



good agreement between data and predictions?



‘zero crossing point’

$$q_0^2 = 4.9 \pm 0.9 \text{ GeV}^2/c^4$$

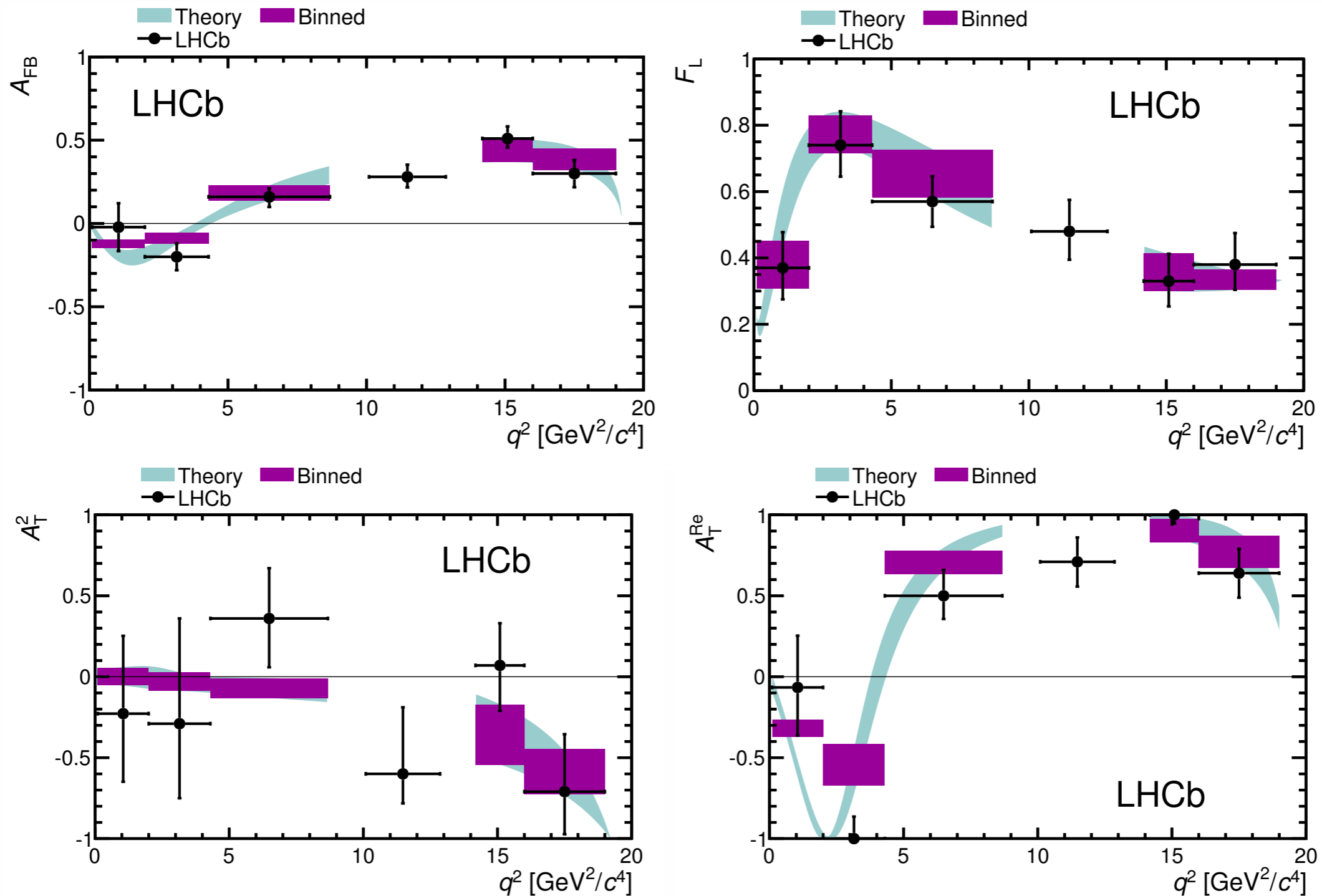
agrees well with SM prediction:

$$q_{0,\text{SM}}^2 \in [3.9, 4.4] \text{ GeV}^2/c^2$$

C. Bobeth, G. Hiller, D. van Dyk and C. Wacker, *The decay  $B \rightarrow K \ell^+ \ell^-$  at low hadronic recoil and model-independent  $\Delta B = 1$  constraints*, *JHEP* **01** (2012) 107 [[arXiv:1111.2558](#)] [[INSPIRE](#)].

M. Beneke, T. Feldmann and D. Seidel, *Exclusive radiative and electroweak  $b \rightarrow d$  and  $b \rightarrow s$  penguin decays at NLO*, *Eur. Phys. J. C* **41** (2005) 173 [[hep-ph/0412400](#)] [[INSPIRE](#)].

A. Ali, G. Kramer and G.-h. Zhu,  *$B \rightarrow K^+ \ell^+ \ell^-$  decay in soft-collinear effective theory*, *Eur. Phys. J. C* **47** (2006) 625 [[hep-ph/0601034](#)] [[INSPIRE](#)].

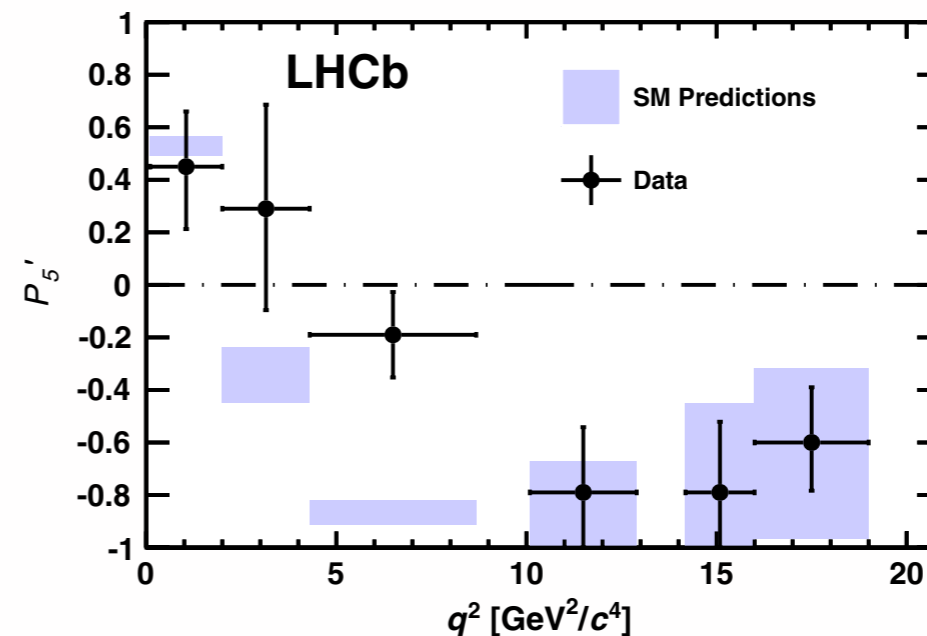
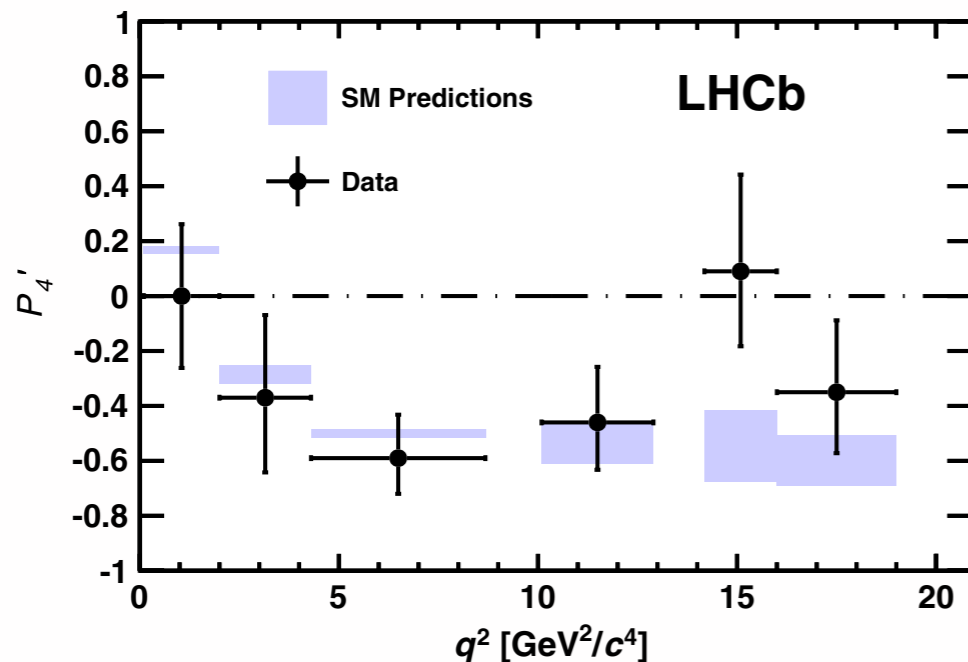


- good agreement with predictions ....
- ... but theory errors large due to 'hadronic form factors'
- > look at ratios of observables to reduce form factors uncertainties



- 
- at low  $q^2$ , ratios  $P'_{i=4,5,6,8} = \frac{S_{i=4,5,7,8}}{\sqrt{F_L(1-F_L)}}$  largely free of FF uncertainties, while sensitivity to NP remains (Descotes-Genon, Hurth, Matias, Virto, [arXiv:1303.5794](#))
  - $P6'$  and  $P8'$  close to SM predictions, which are close to zero for all  $q^2$
  - $P4'$  and  $P5'$  have less trivial structure:

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- note the tension in  $P5'$ :  $3.7\sigma$  (1 out of 24 bins: global significance  $2.8\sigma$ )
  - ▶ change in  $C9$ ? (Descotes, Matias, Virto, [arXiv:1307.5683](#)) by a  $Z'$ ? (Gauld et al, [arXiv:1308.1959](#); Buras and Girschbach, [arXiv:1309.2466](#))
  - ▶ or more mundane effects of QCD, 'non-factorizable form factors'?

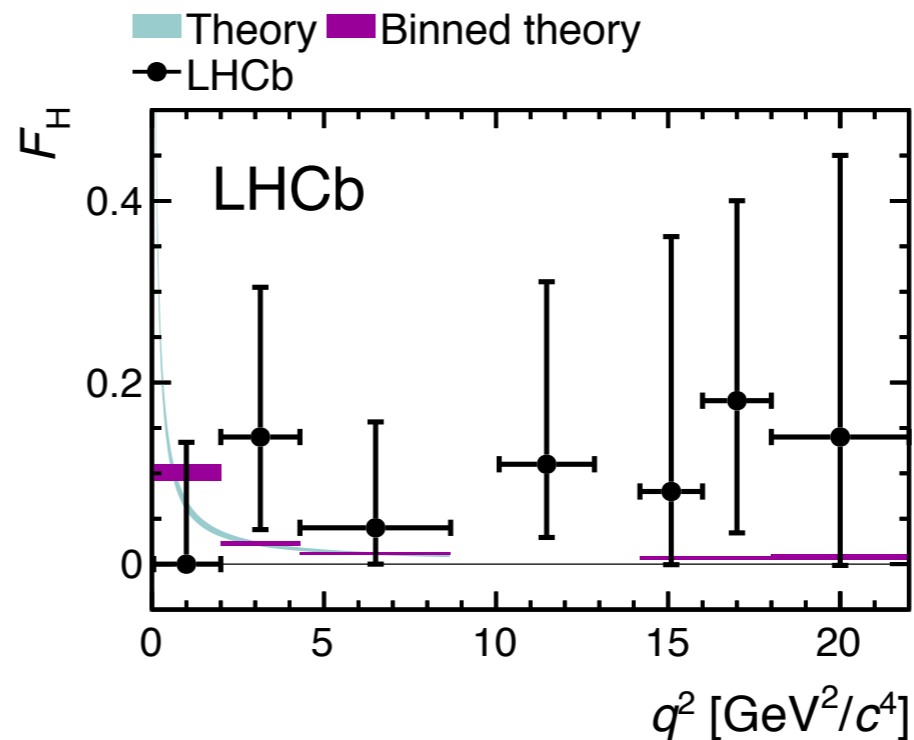
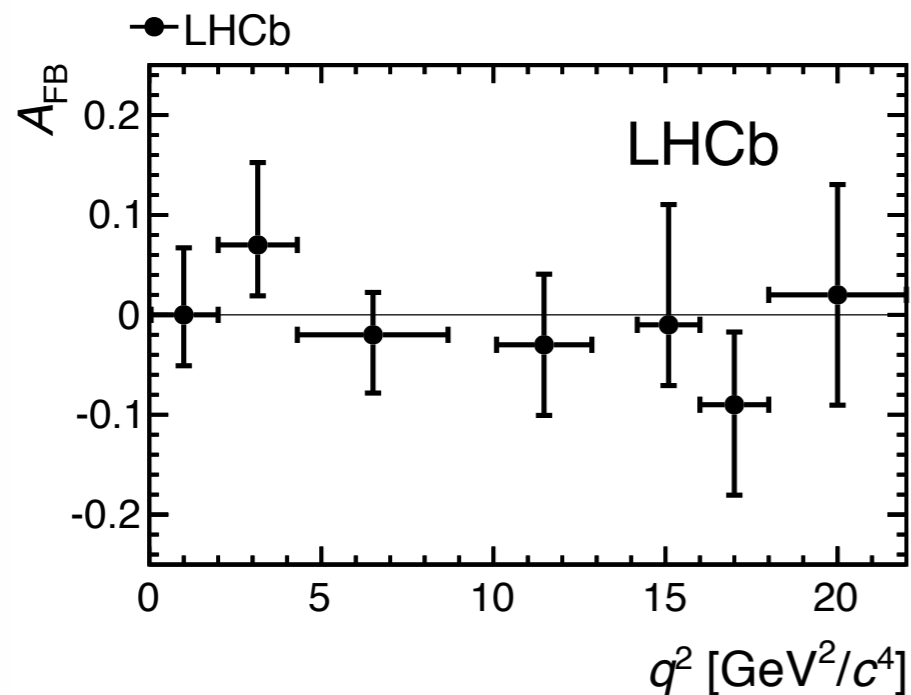
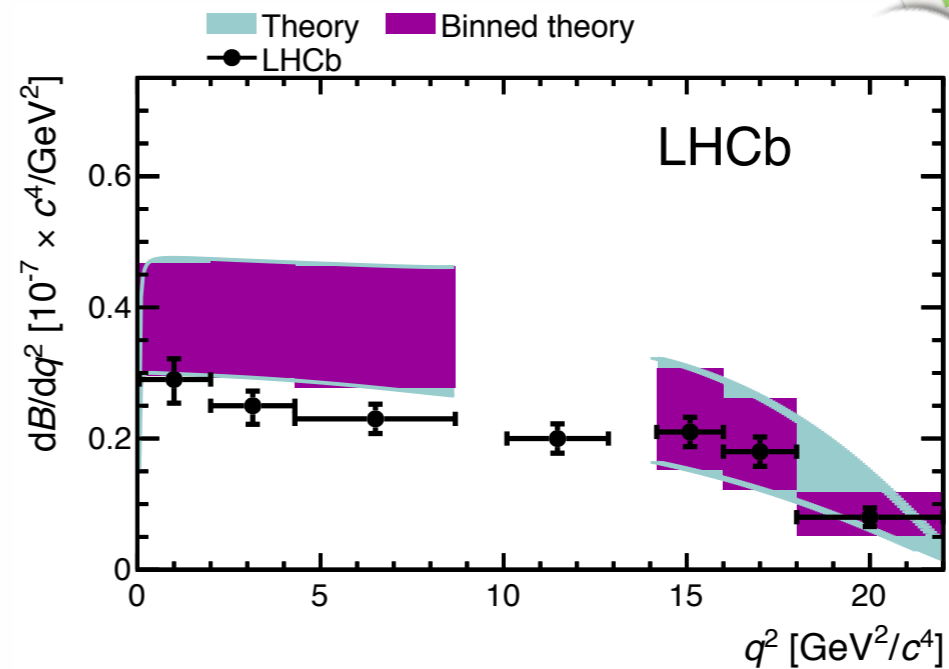
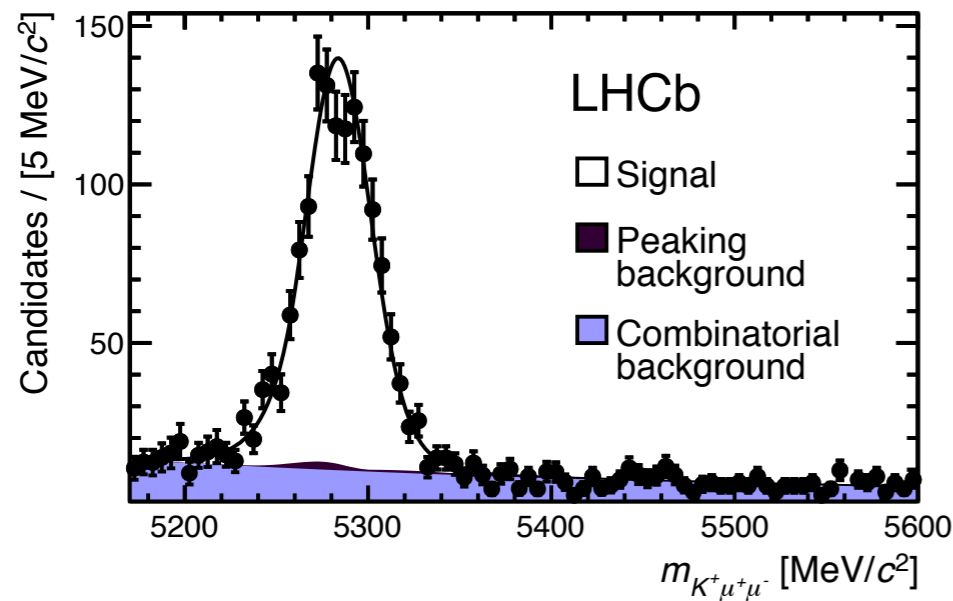
# $B^+ \rightarrow K^+ \mu^+ \mu^-$ at LHCb

- $B^+ \rightarrow K^+ \mu^+ \mu^-$  also extensively studied at LHCb

- differential distributions (1/fb) [arXiv:1209.4284, JHEP02\(2013\)105](https://arxiv.org/abs/1209.4284)

no surprises

1240 candidates in 1/fb



# $B^+ \rightarrow K^+ \mu^+ \mu^-$ at LHCb

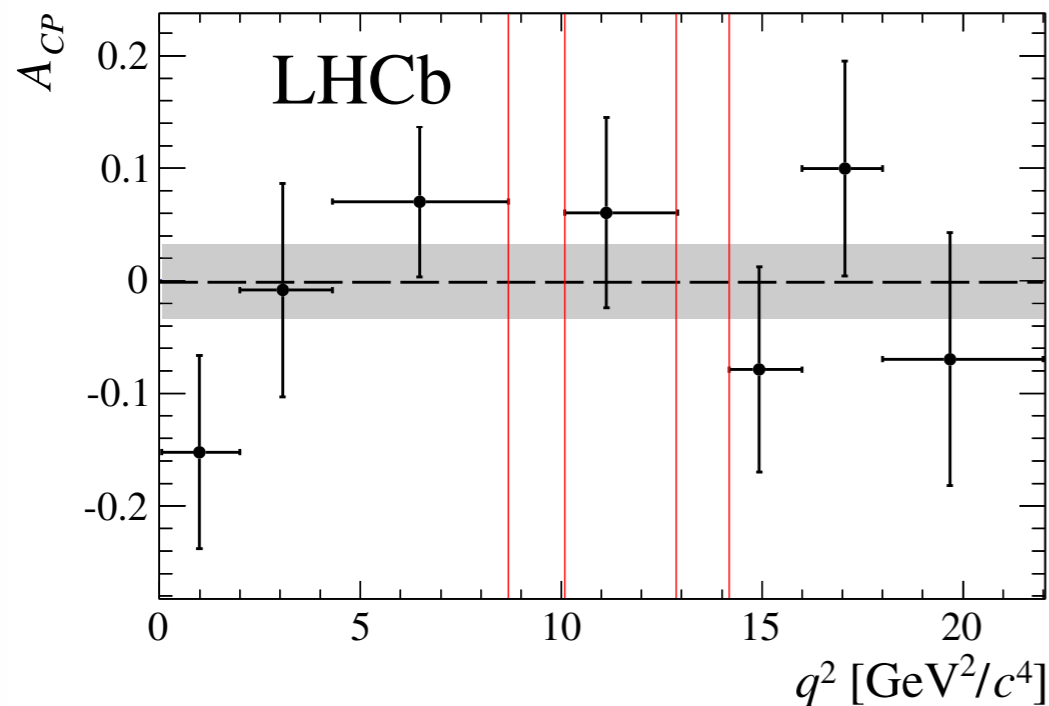
●  $B^+ \rightarrow K^+ \mu^+ \mu^-$  also extensively studied at LHCb

● differential distributions (1/fb)

arXiv:1209.4284, JHEP02(2013)105

● CP asymmetry (1/fb)

arXiv:1308.1340, PRL 111, 151801 (2013)



efficiency-corrected  $q^2$  average

$$A_{CP} = 0.000 \pm 0.033(\text{stat.}) \pm 0.005(\text{syst.}) \pm 0.007(J/\psi K).$$

consistent with SM ...

# $B^+ \rightarrow K^+ \mu^+ \mu^-$ at LHCb

- $B^+ \rightarrow K^+ \mu^+ \mu^-$  also extensively studied at LHCb

- differential distributions (1/fb)

arXiv:1209.4284, JHEP02(2013)105

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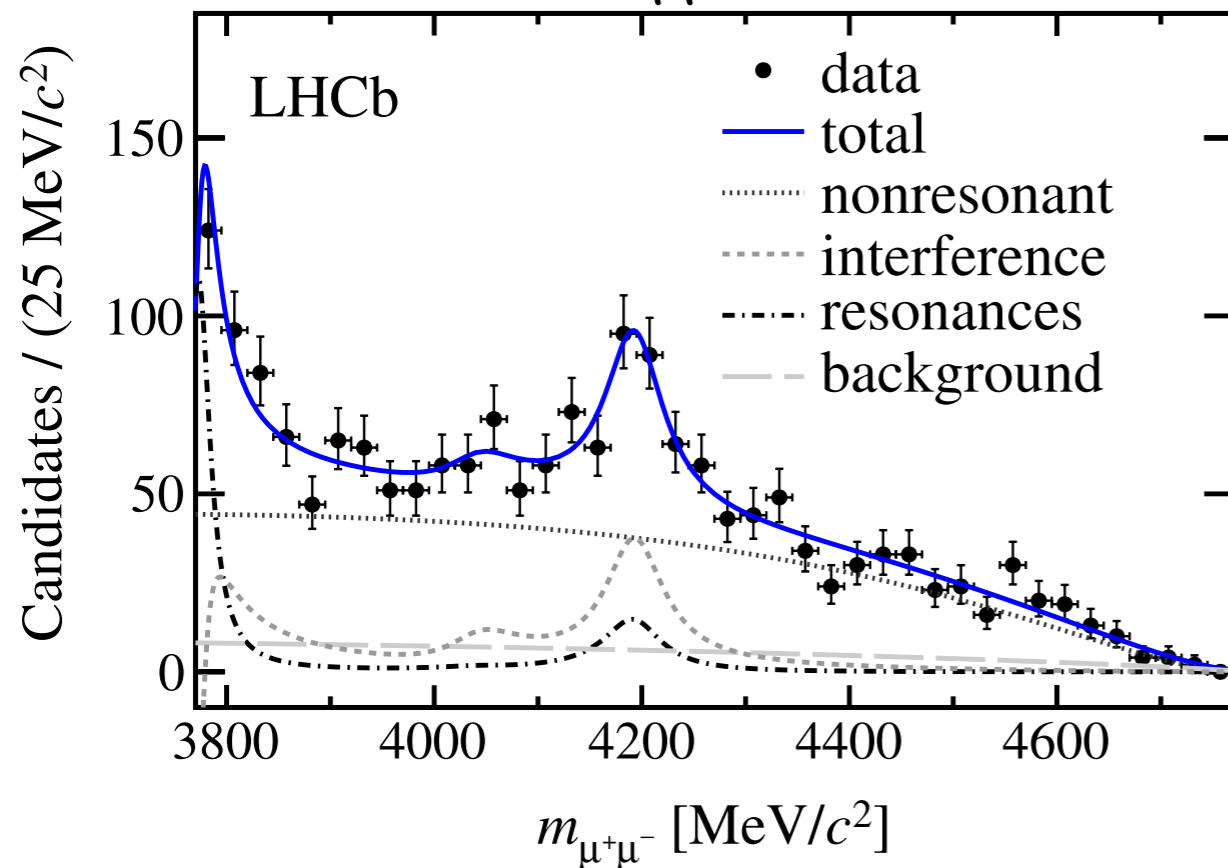
arXiv:1308.1340, PRL 111, 151801 (2013)

- resonance structure (3/fb)

arXiv:1307.7595, PRL 111, 112003 (2013)



830  $B^+ \rightarrow K^+ \mu^+ \mu^-$  candidates



- ▶  $\mu^+ \mu^-$  resonance 'at low recoil' (high- $q^2$ )
- ▶ significance  $> 6\sigma$
- ▶ consistent with  $\psi(4160)$  found by BES

	Unconstrained	$\psi(4160)$
$\mathcal{B}[\times 10^{-9}]$	$3.9^{+0.7}_{-0.6}$	$3.5^{+0.9}_{-0.8}$
Mass [MeV/ $c^2$ ]	$4191^{+9}_{-8}$	$4190 \pm 5$
Width [MeV/ $c^2$ ]	$65^{+22}_{-16}$	$66 \pm 12$
Phase [rad]	$-1.7 \pm 0.3$	$-1.8 \pm 0.3$

Q: what does this mean for reliability of predictions in  $B^{*0} \rightarrow K^{*0} \mu \mu$  at high  $q^2$ ?

isospin asymmetry:

$$A_I = \frac{\Gamma(B^0 \rightarrow K^{(*)0} \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^{(*)+} \mu^+ \mu^-)}{\Gamma(B^0 \rightarrow K^{(*)0} \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^{(*)+} \mu^+ \mu^-)}$$

predicted to be close  
to zero in SM

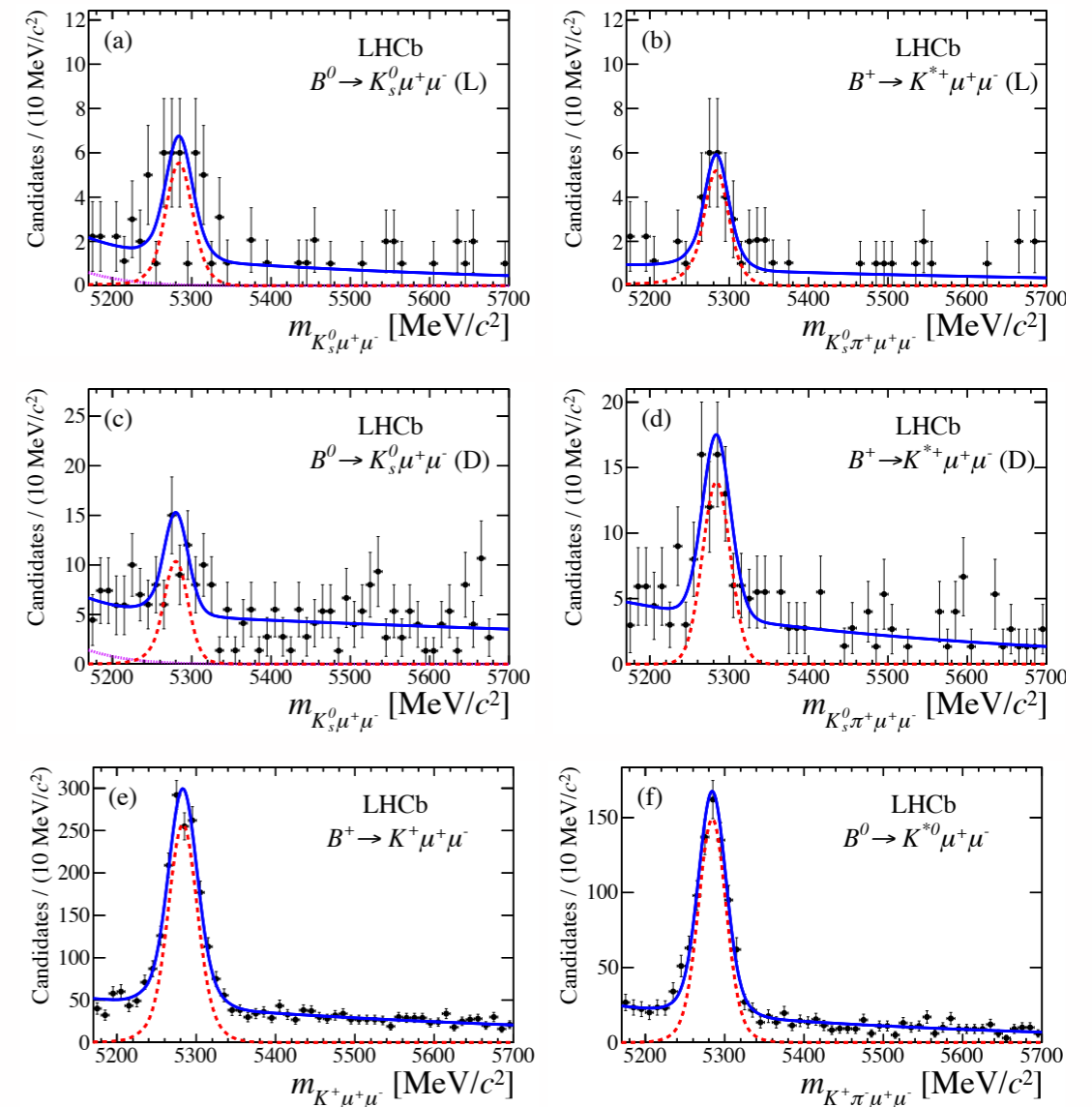
LHCb measurement on 1/fb using final  
states with charged particles only:

$$B^0 \rightarrow K^0 \mu^+ \mu^- \rightarrow K_s \mu^+ \mu^-$$

$$B^+ \rightarrow K^+ \mu^+ \mu^-$$

$$B^0 \rightarrow K^{*0} \mu^+ \mu^- \rightarrow K^+ \pi^- \mu^+ \mu^-$$

$$B^+ \rightarrow K^{*+} \mu^+ \mu^- \rightarrow K_s^0 \pi^+ \mu^+ \mu^-$$



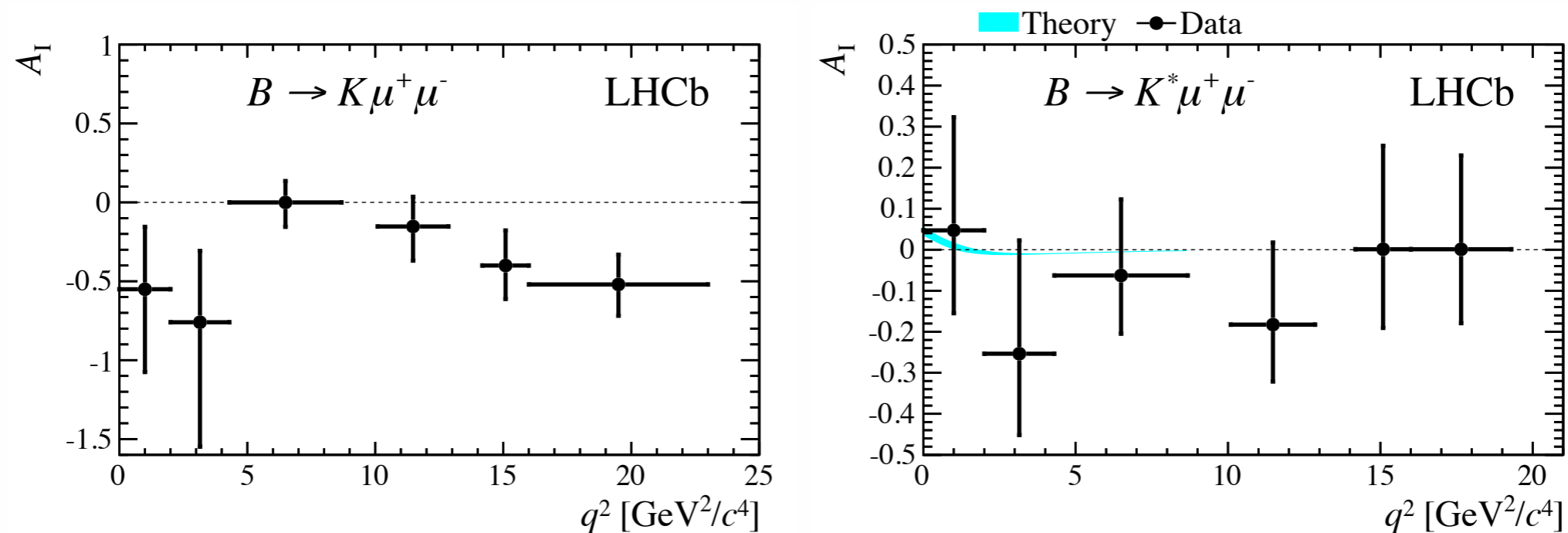


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LHCb measurement on 1/fb using final  
states with charged particles only:



Both consistent with previous measurements (Babar, Belle,  
CDF) ... but  $K \mu \mu$  has a  $4.4\sigma$  deviation from zero?!

# More $B \rightarrow X \mu \mu$ and $B \rightarrow X \gamma$

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Similar physics, omitted due to lack of time ...

- $B_s \rightarrow \phi \mu \mu$

arXiv:1305.2168, JHEP. 07 (2013) 084

- $\Lambda_b \rightarrow \Lambda \mu \mu$

arXiv:1306.2577, PLB 725 (2013) 25

- $B^+ \rightarrow \pi^+ \mu^+ \mu^-$

arXiv:1210.2645, JHEP. 12 (2012) 125

- $B^+ \rightarrow K \pi \pi \gamma$  (photon polarization)

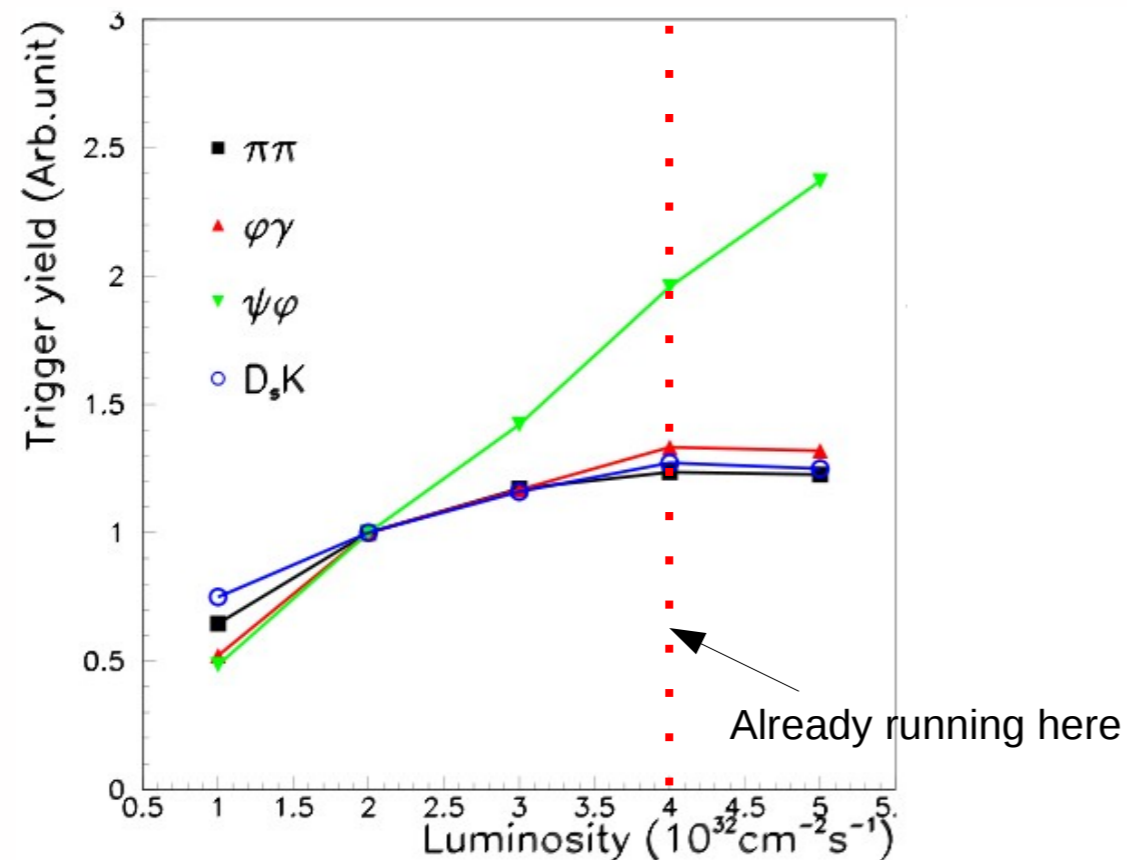
LHCb-CONF-2013-009

- $B^0 \rightarrow K^* \gamma$  and  $B_s \rightarrow \phi \gamma$

arXiv:1209.0313, Nucl. Phys. B 867 (2012) 1

# LHCb future

- 2010-2012: collected about 3/fb at effective  $L_{\text{inst}}$  up to  $4 \times 10^{32} / \text{cm}^2/\text{s}$
- 2015-2017: collect another  $\sim 5/\text{fb}$  at similar lumi
- 2017-2018: LHC LS2 ... then what?



- ▶ current LO output rate cannot exceed  $\sim 1$  MHz
- ▶ to keep rate limited, cut hard on hadron  $p_T$
- ▶ with current detector, no gain in hadronic event yields from increased lumi!

- LHCb upgrade: prepare detector for  $L_{\text{inst}} \sim 20 \times 10^{32} / \text{cm}^2/\text{s}$ 
  - read out full detector at 40 MHz
  - full software trigger with momentum/vertex reconstruction
  - velo pixel detector, Si and scintillating fibre tracker, new RICH, ....

# Summary

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- flavour physics at the LHC works!
  - large yields, clean signals (cleaner than expected by many of us!)
  - many important new results
  
- the Standard Model has survived ... so far
  
  
  
  
  
  
  
  
  
  
- exciting prospects for the future
  - 2015-2017: double (triple?) existing dataset
  - 2017-2019: LS2, upgrade installation
  - 2019- ...: collect  $\sim 50/\text{fb}$ , with up to 2x more efficient trigger

backup

# Mixing in the $D^0$ system

1/fb: arXiv:1211.1230, PRL 110.101802 (2013)

3/fb: arXiv:1309.6534, subm. to PRL

- in  $D^0$  system flavour can be tagged via  $D^*$ :

right sign:  $D^{*+} \rightarrow \pi^+ D^0 \rightarrow \pi^+ (K^- \pi^+)_{D^0}$

wrong sign:  $D^{*+} \rightarrow \pi^+ D^0 \rightarrow \pi^+ (K^+ \pi^-)_{D^0}$

- mixing probed with time-dependence of  $R = \frac{N(\text{wrong})}{N(\text{right})}$



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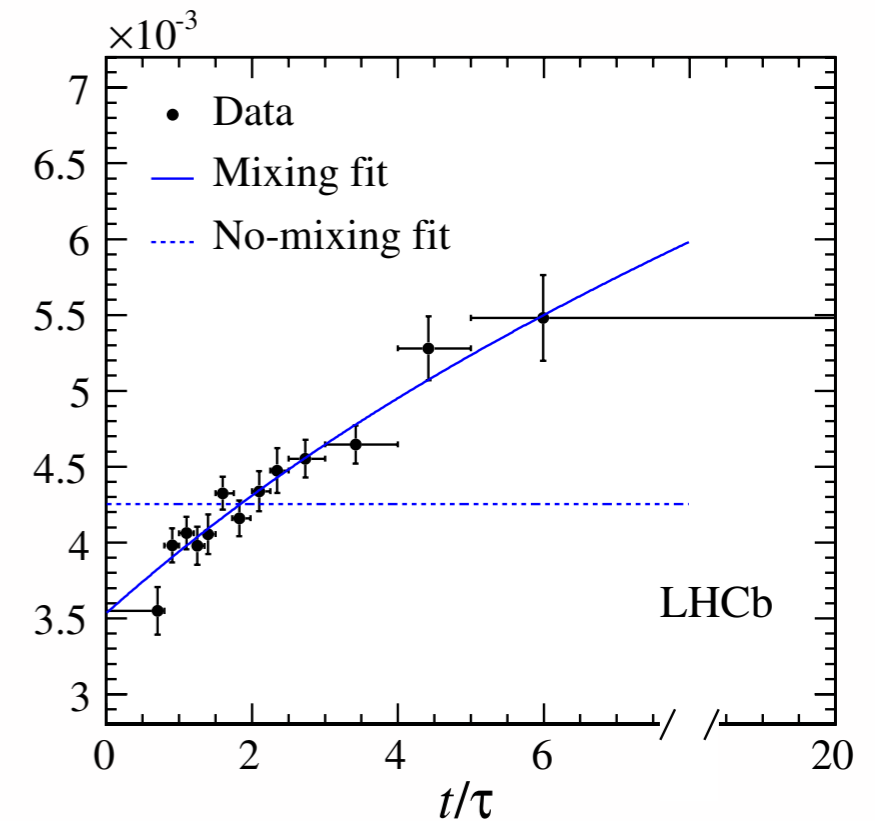
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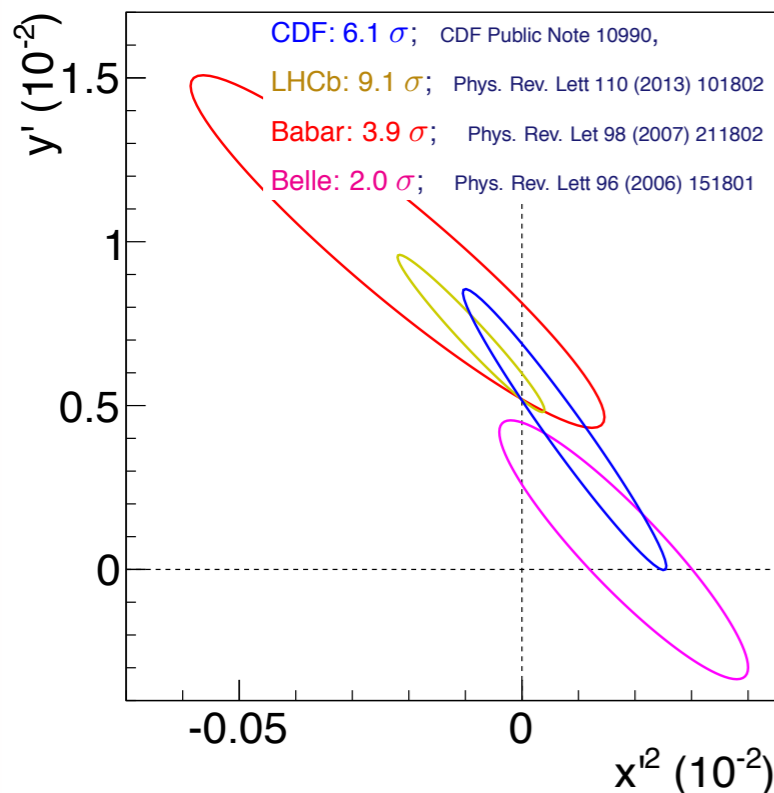
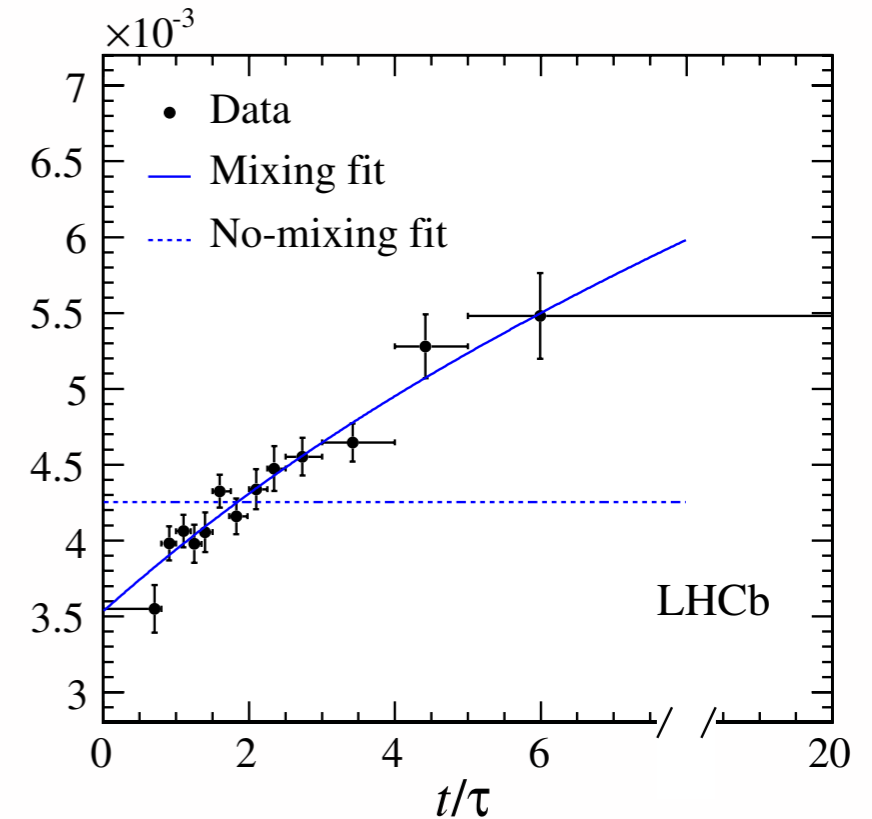
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- mixing probed with time-dependence of  $R = \frac{N(\text{wrong})}{N(\text{right})}$

$$R(t) \equiv \frac{N_{\text{ws}}(t)}{N_{\text{rs}}(t)} \approx R_D + \sqrt{R_D} y' t + \frac{1}{4} (x' + y') t^2$$



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1/fb: arXiv:1211.1230, PRL 110.101802 (2013)

3/fb: arXiv:1309.6534, subm. to PRL

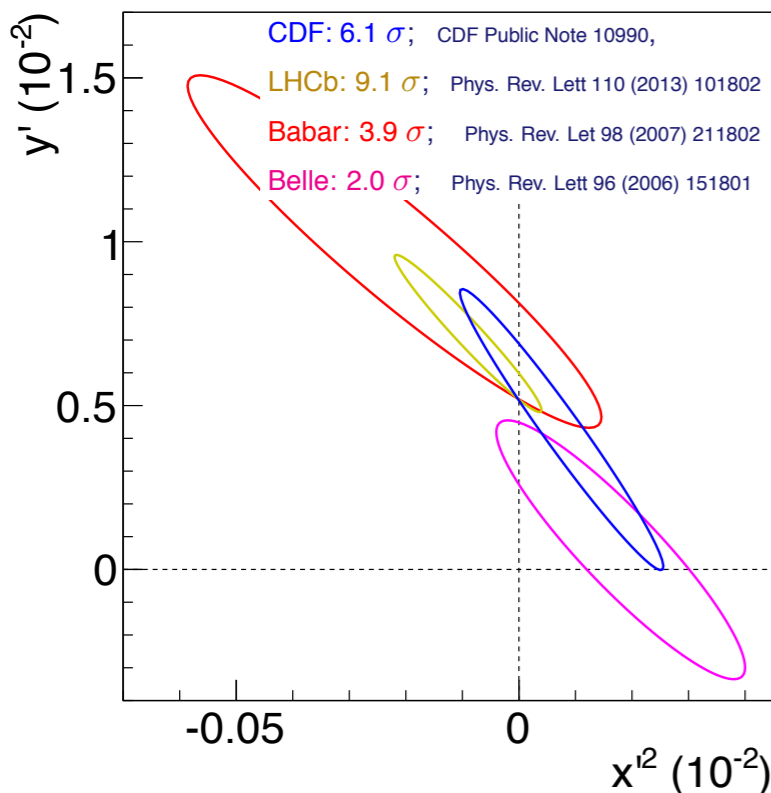
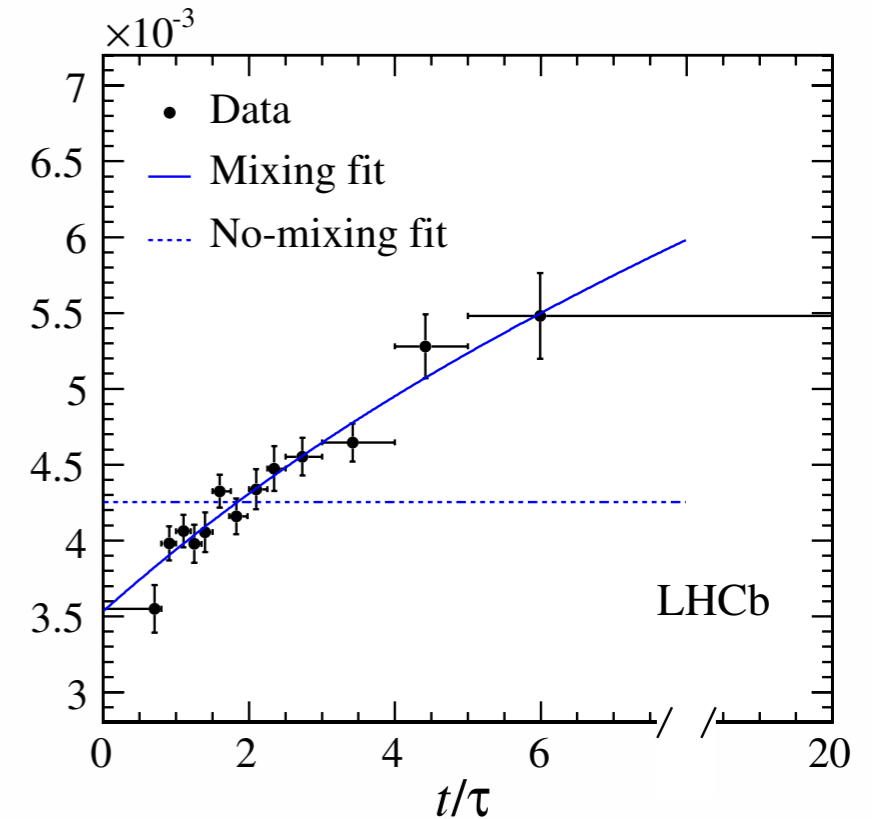
- in  $D^0$  system flavour can be tagged via  $D^*$ :

right sign:  $D^{*+} \rightarrow \pi^+ D^0 \rightarrow \pi^+ (K^- \pi^+)_{D^0}$

wrong sign:  $D^{*+} \rightarrow \pi^+ D^0 \rightarrow \pi^+ (K^+ \pi^-)_{D^0}$

- mixing probed with time-dependence of  $R = \frac{N(\text{wrong})}{N(\text{right})}$

$$R(t) \equiv \frac{N_{\text{ws}}(t)}{N_{\text{rs}}(t)} \approx R_D + \sqrt{R_D} y' t + \frac{1}{4} (x' + y') t^2$$



- $D^0$  mixing now firmly established:

$$x'^2 = (5.5 \pm 4.9) \times 10^{-5}, \quad y' = (4.8 \pm 1.0) \times 10^{-3}$$

- latest LHCb analysis (3/fb) also measures CPV in decay and mixing:

$$A_D \equiv (R_D^+ - R_D^-)/(R_D^+ + R_D^-) = (-0.7 \pm 1.9)\%$$

$$|q/p| = 1.00 \pm 0.25$$

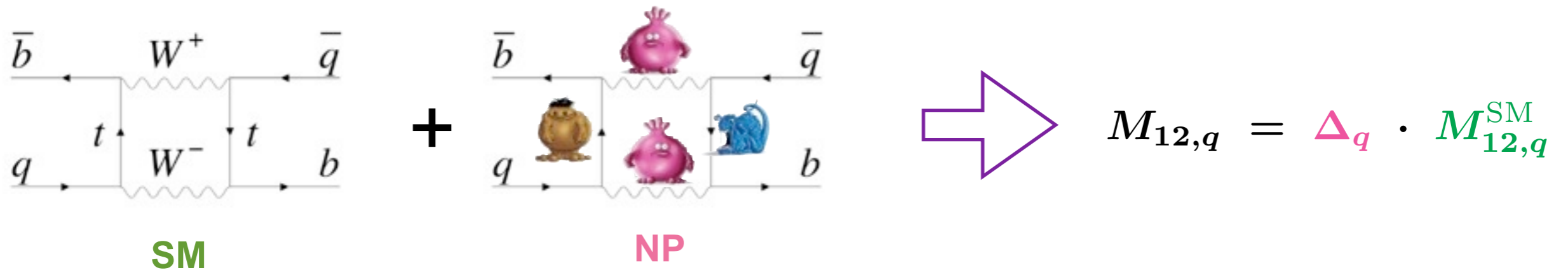
no sign of CPV

- related: measurement of indirect CPV using effective lifetimes in  $KK$  and  $\pi\pi$ :

arXiv:1310.7201, acc. PRL

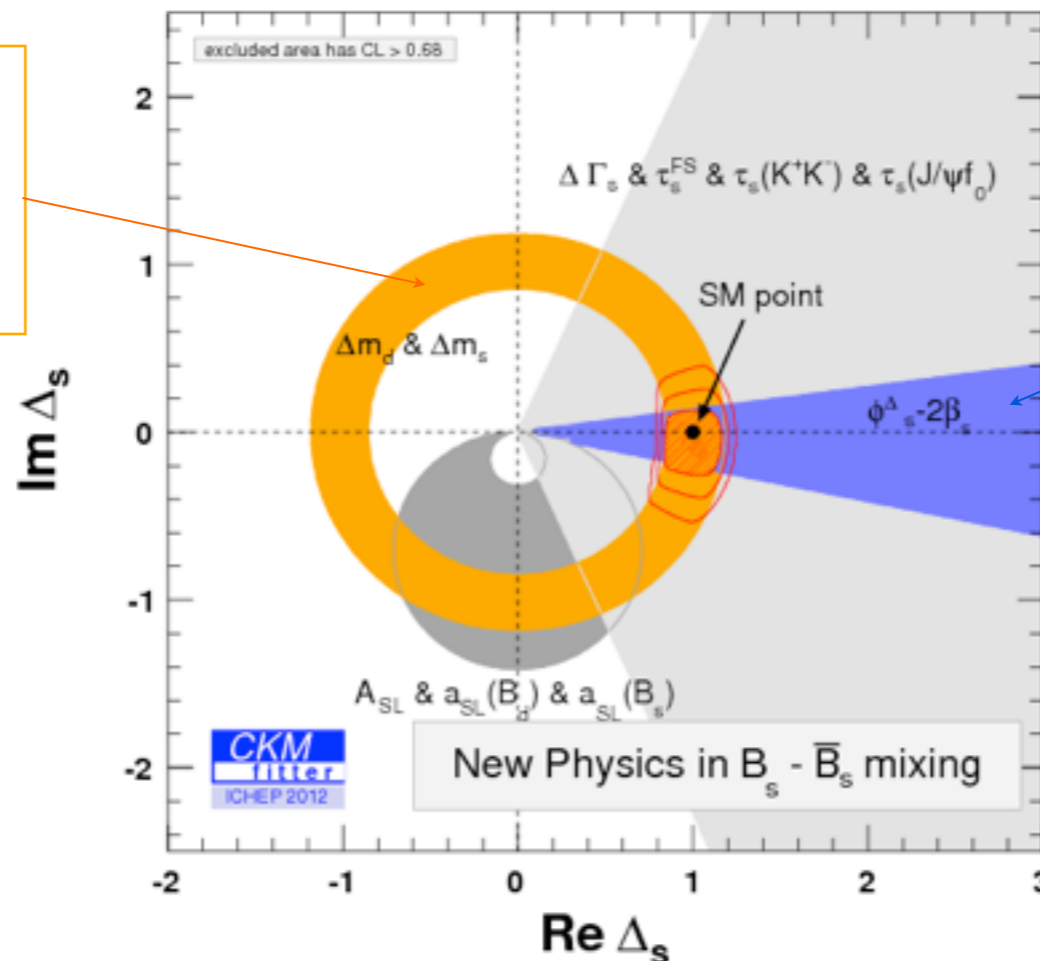
# Constraints on $M_{12}$

- suppose new physics affects  $M_{12}$



- constraints from mixing measurements

constraint from oscillation frequencies, theory limited



constraint from indirect CP violation, statistically limited

From Lenz, Nierste and CKMFitter, arXiv:1203.0238

# The future of gamma: $B_s \rightarrow D_s K$

- extract 'sin(2beta\_s - gamma)' from amplitudes of oscillations in  $B_s^0 \rightarrow D_s^\mp K^\pm$

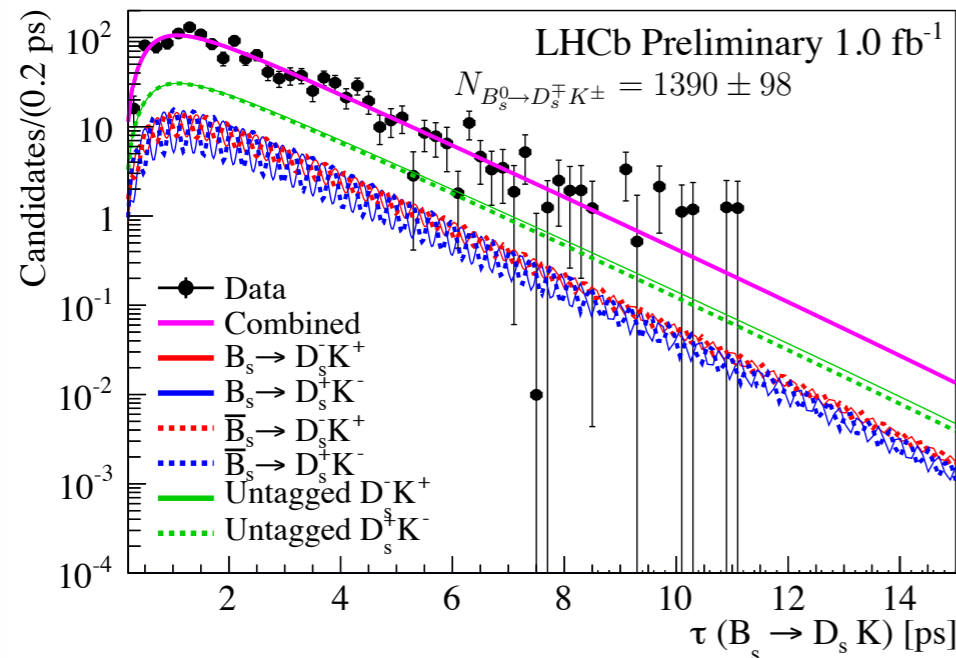
$$\frac{d\Gamma_{B_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2}|A_f|^2(1 + |\lambda_f|^2)e^{-\Gamma_s t} \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) - D_f \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t) \right],$$

$$C = \frac{1 - r_{D_s K}^2}{1 + r_{D_s K}^2},$$

$$D_f = \frac{2r_{D_s K} \cos(\Delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad D_{\bar{f}} = \frac{2r_{D_s K} \cos(\Delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2},$$

$$S_f = \frac{2r_{D_s K} \sin(\Delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad S_{\bar{f}} = \frac{2r_{D_s K} \sin(\Delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}.$$

- small theory uncertainties, but requires flavour tagging (see later)
- first results on 2011 data (LHCb-CONF-2012-029)



$$C = 1.01 \pm 0.50 \pm 0.23,$$

$$S_f = -1.25 \pm 0.56 \pm 0.24,$$

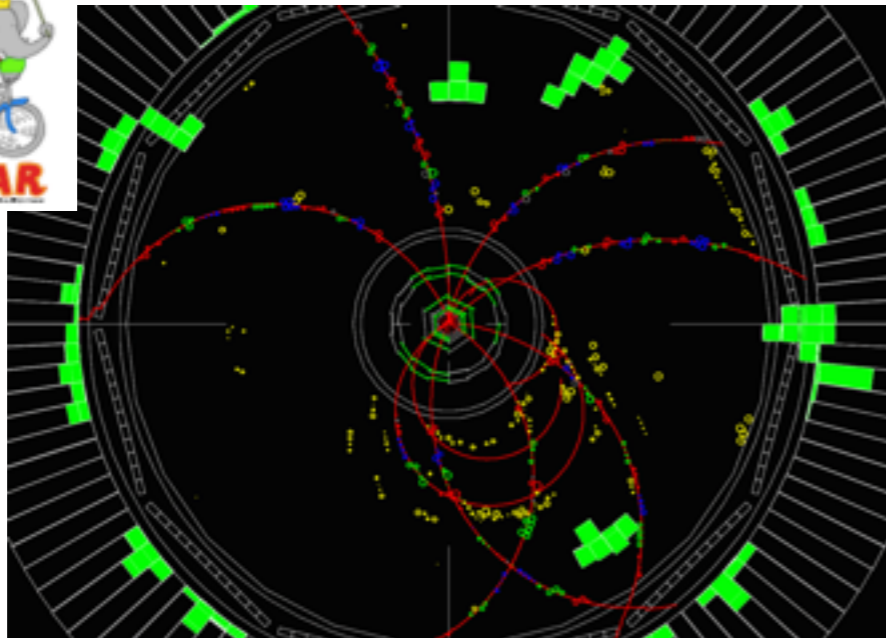
$$S_{\bar{f}} = 0.08 \pm 0.68 \pm 0.28,$$

$$D_f = -1.33 \pm 0.60 \pm 0.26,$$

$$D_{\bar{f}} = -0.81 \pm 0.56 \pm 0.26,$$

- first constraints of gamma expected from 3/fb analysis (in progress)
- systematics limited by lifetimes, mixing and tagging calibration: will all improve with more LHCb data

# The price to pay ...



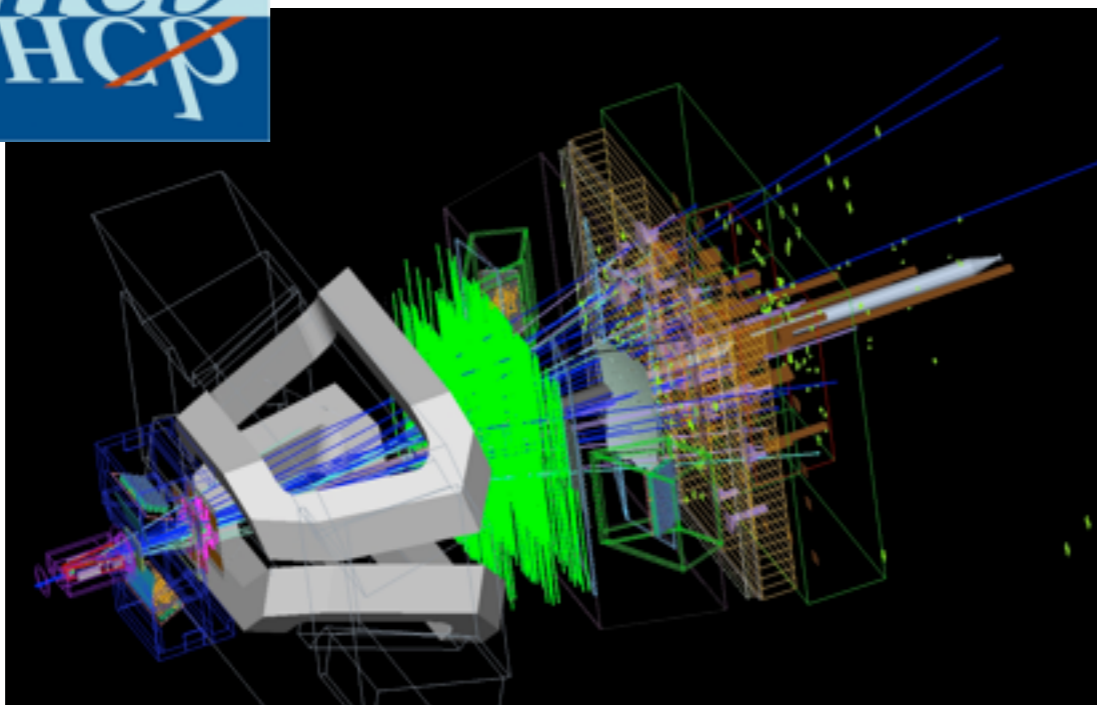
$\sim 10$  b-bbar events / s  
 $O(20)$  particles/event

full event reconstruction

typical B $\rightarrow$ X efficiency  $\sim 30\%$

$Q_{\text{tag}} \sim 30\%$

$\sigma(t) \sim 0.6$  ps



50000 b-bbar events / s  
 $O(200)$  particles/event

typical B $\rightarrow$ X efficiency 1%

$Q_{\text{tag}} \sim 2-4\%$

$\sigma(t) \sim 0.05$  ps

good with muons, reasonably good with hadrons, but not so good with photons and flavour tagging

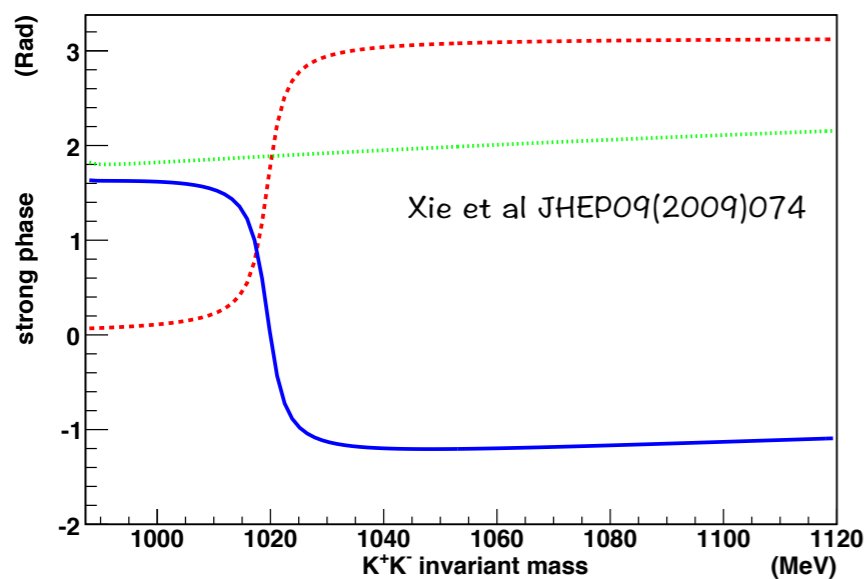


■ ambiguity:

$$(\phi_s, \Delta\Gamma_s, \delta_0, \delta_{\parallel}, \delta_{\perp}, \delta_S) \mapsto (\pi - \phi_s, -\Delta\Gamma_s, -\delta_0, -\delta_{\parallel}, \pi - \delta_{\perp}, -\delta_S)$$

■ solution: measure phase difference between S and P-wave in bins of mass (Xie, e.a. 2009)

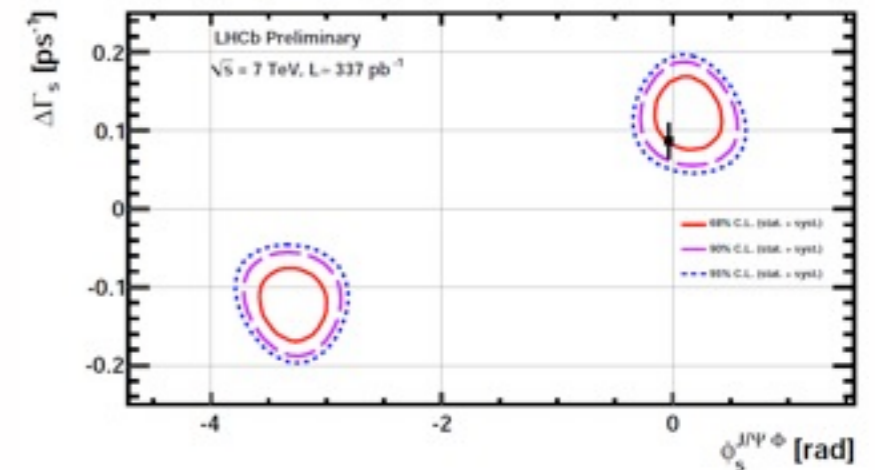
expected:



P-wave: relativistic Breit wigner

S-wave: not well known, but mild variation with mKK difference

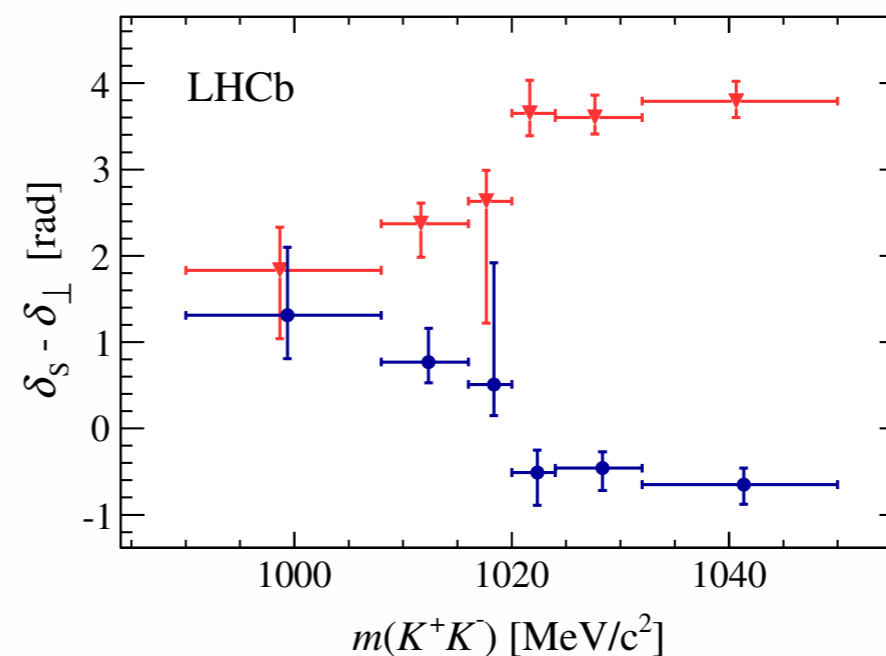
(note: plot for first 0.37/fb only)



■ unambiguous!

only  $\Delta\Gamma > 0$  solution fits expected pattern

observed:



$\Delta\Gamma < 0$

$\Delta\Gamma > 0$

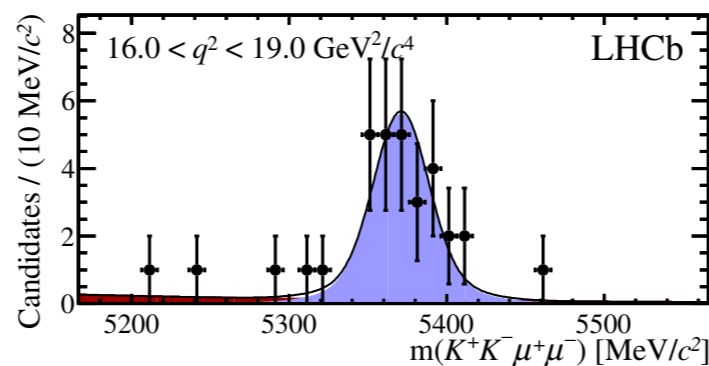
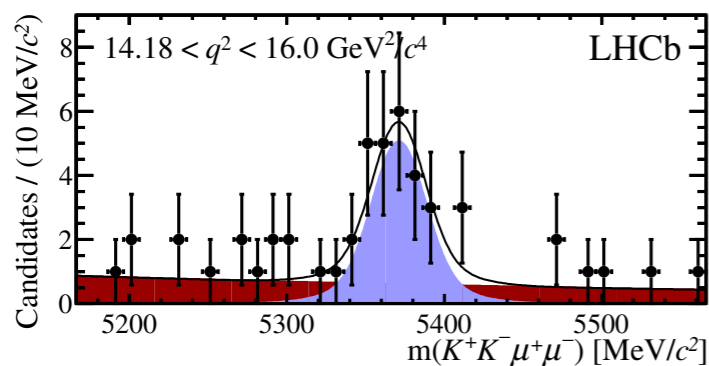
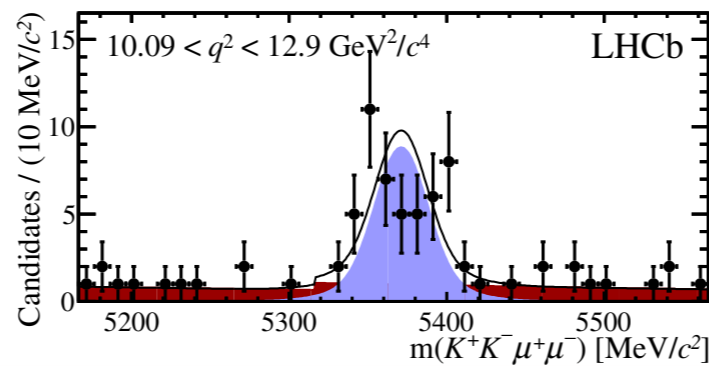
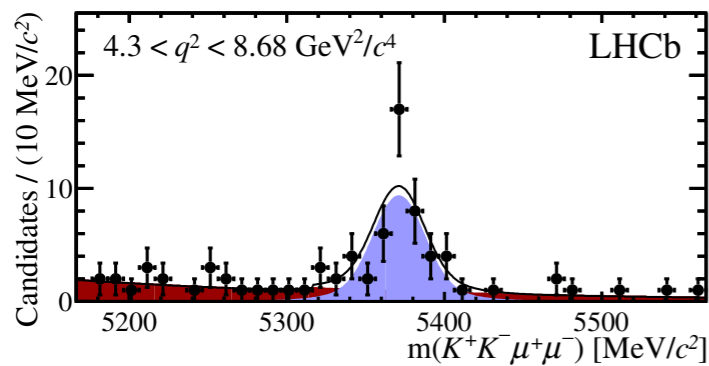
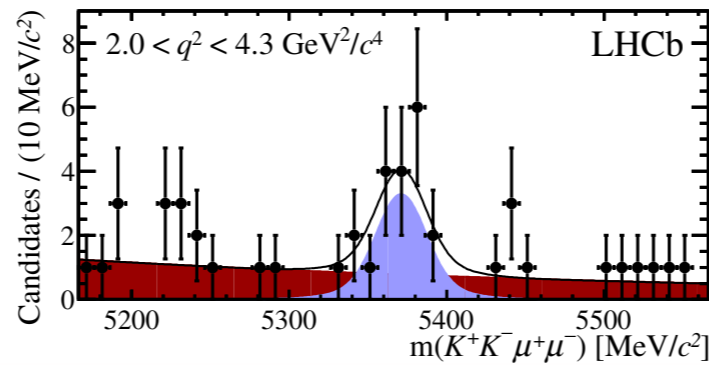
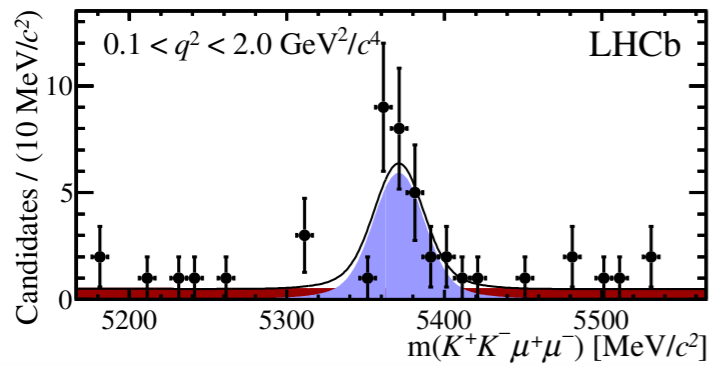


# $B_s \rightarrow \phi \mu \mu$

174 +/- 15 candidates in 1/fb

integrated BF:  $\mathcal{B}(B_s^0 \rightarrow \phi \mu^+ \mu^-) = (7.07^{+0.64}_{-0.59} \pm 0.17 \pm 0.71) \times 10^{-7}$

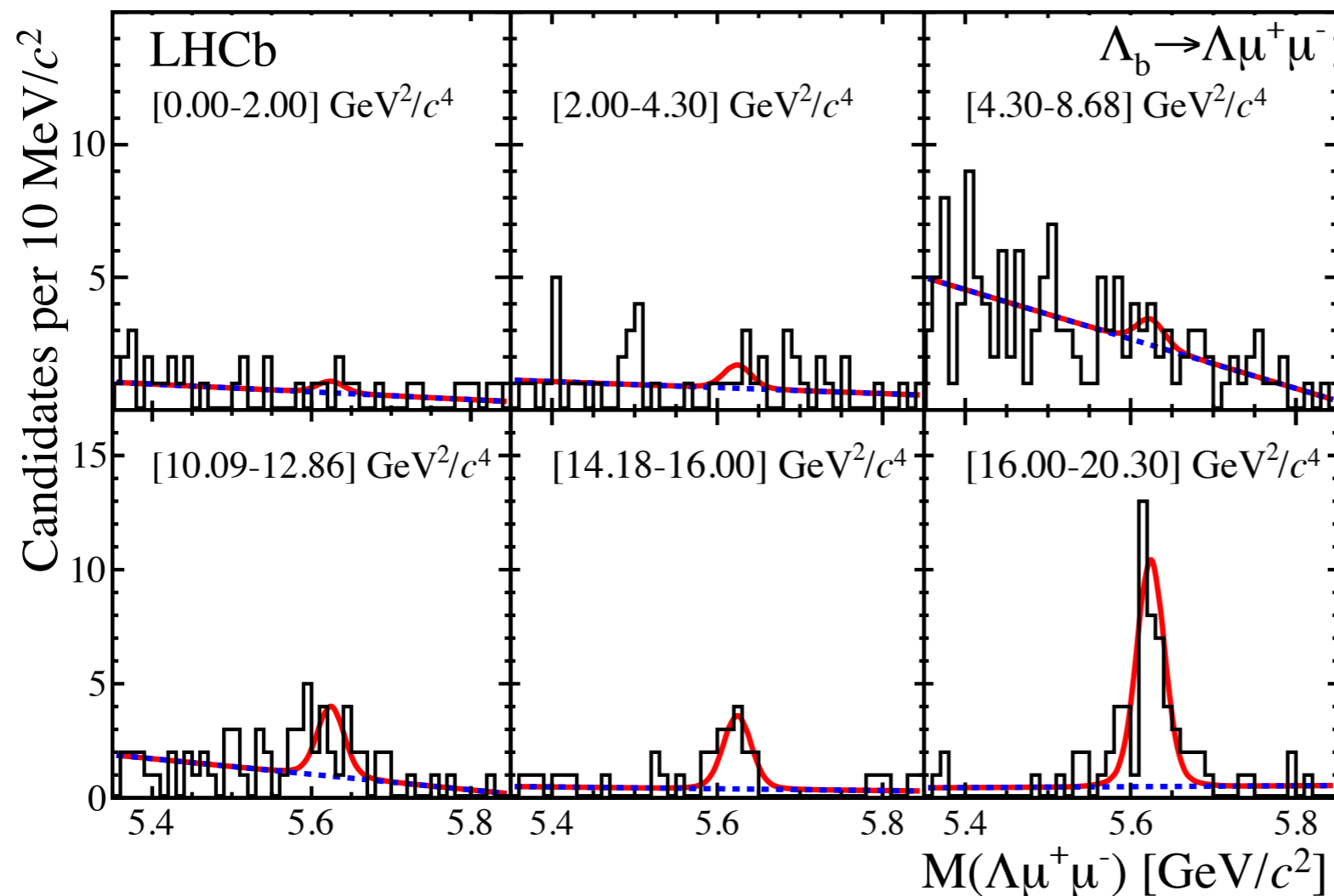
error on BF( $J/\psi\phi$ )



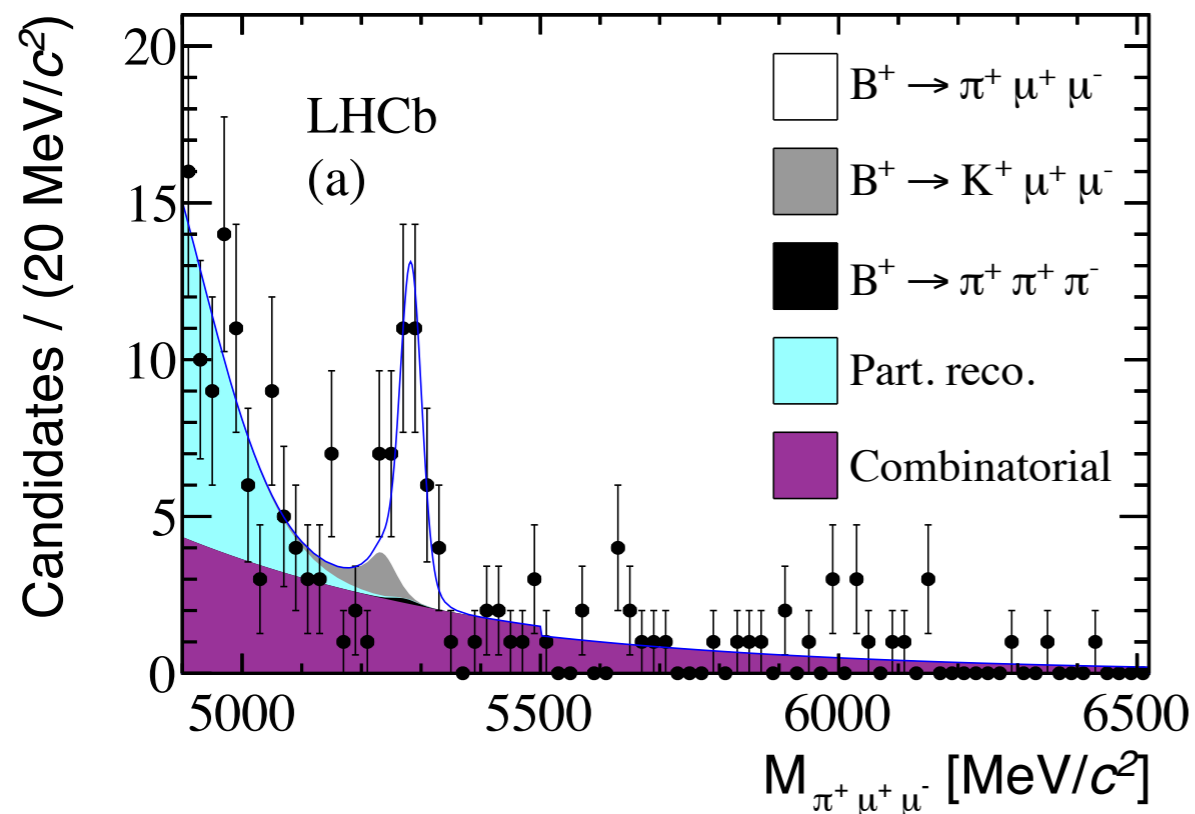
# $\Lambda_b \rightarrow \Lambda \mu \mu$

78 +/- 12 events in 1/fb

integrated BF:  $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-) = (0.96 \pm 0.16 \text{ (stat)} \pm 0.13 \text{ (syst)} \pm 0.21 \text{ (norm)}) \times 10^{-6}$



about 25 +/- 7 events in 1/fb



$$\frac{\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)} = 0.053 \pm 0.014 \text{ (stat.)} \pm 0.001 \text{ (syst.)}.$$

$$\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-) = (2.3 \pm 0.6 \text{ (stat.)} \pm 0.1 \text{ (syst.)}) \times 10^{-8}.$$

my very naive estimate of the  $B \rightarrow \rho \mu \mu$  yield:

$$25 \times 883 / 1240 = 17 \text{ events / fb ?}$$

(that's why nobody looked yet)

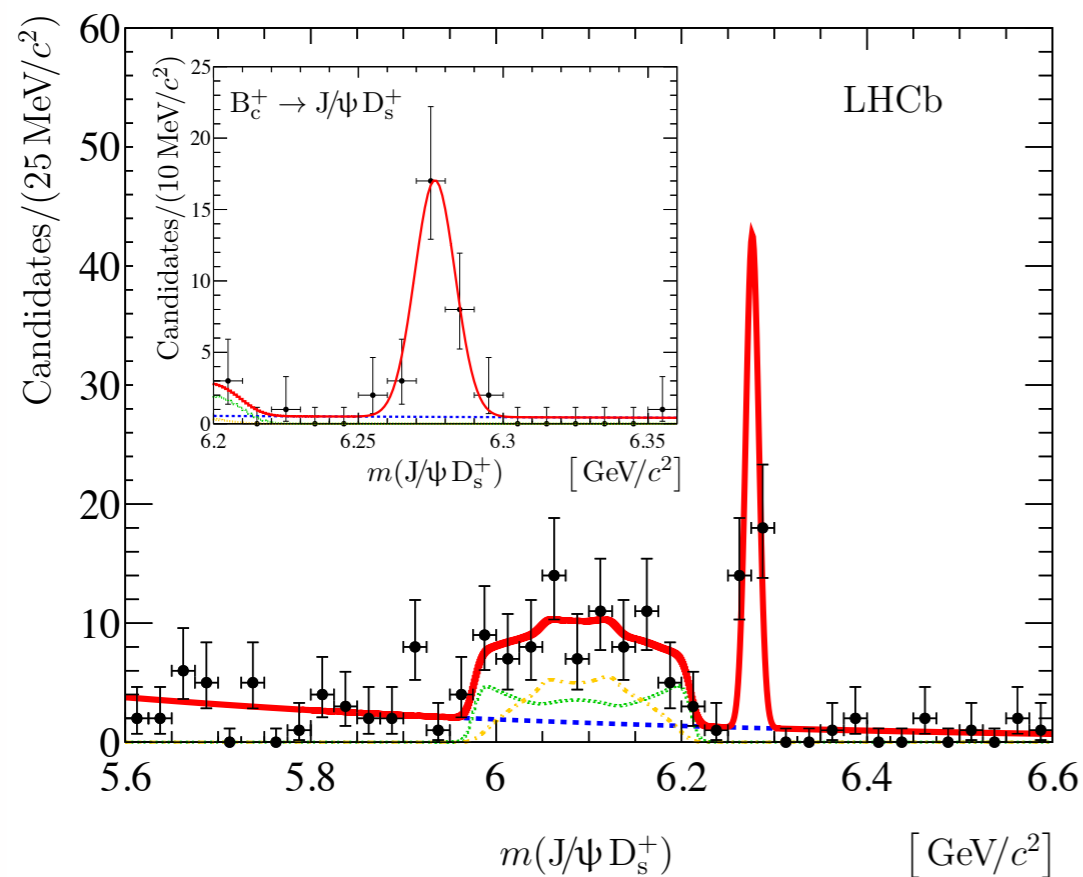
# can LHCb do $B_c \rightarrow D_s \mu\mu$ ?

well ... we have just seen  $B_c \rightarrow J/\psi D_s^{(*)}$

arXiv:1304.4530,  
PRD 87 (2013) 112012

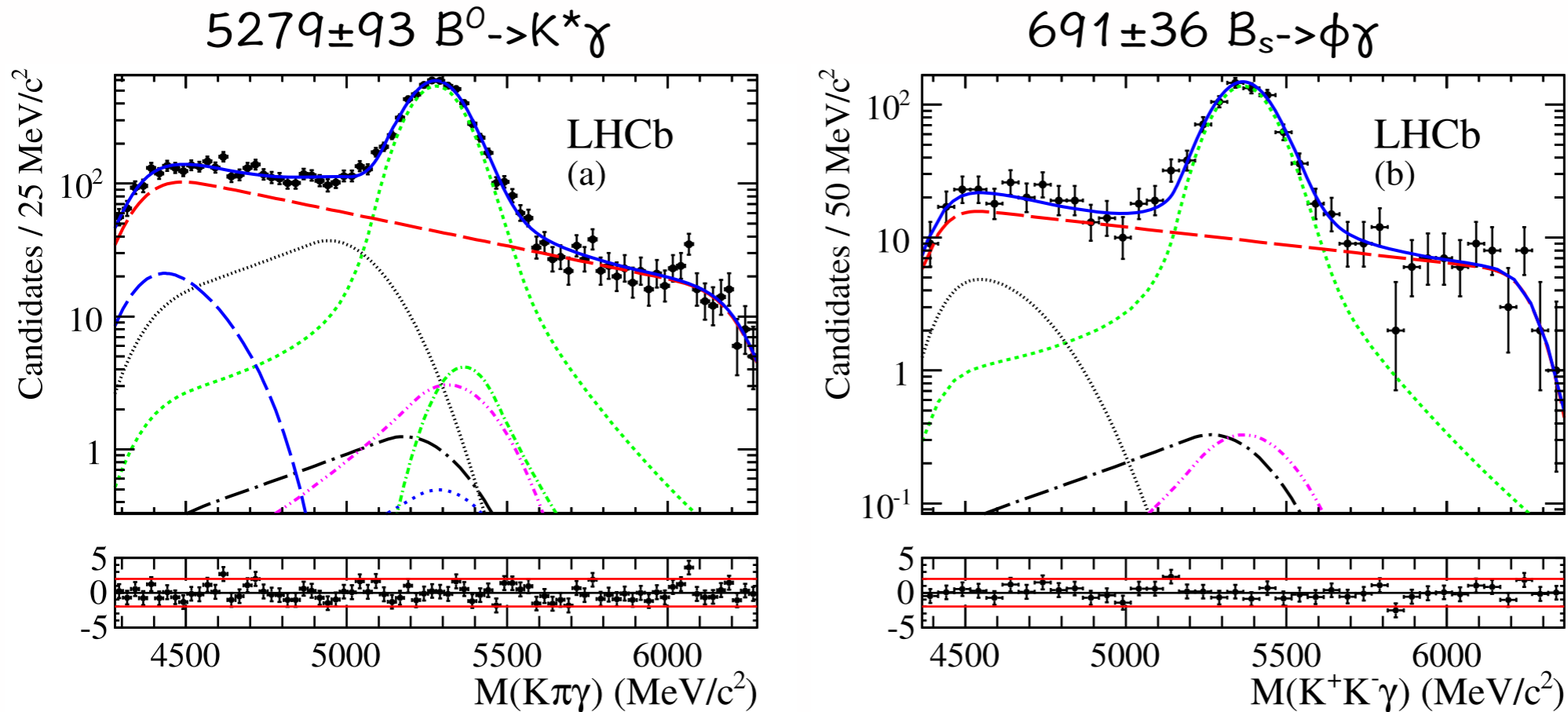
$$N(B_c \rightarrow J/\psi D_s) = 28.9 \pm 5.6$$

$$N(B_c \rightarrow J/\psi D_s^*) = 68.5 \pm ?$$



draw you own conclusion ...

measurements of BF ratios and direct CP asymmetry in 1/fb:



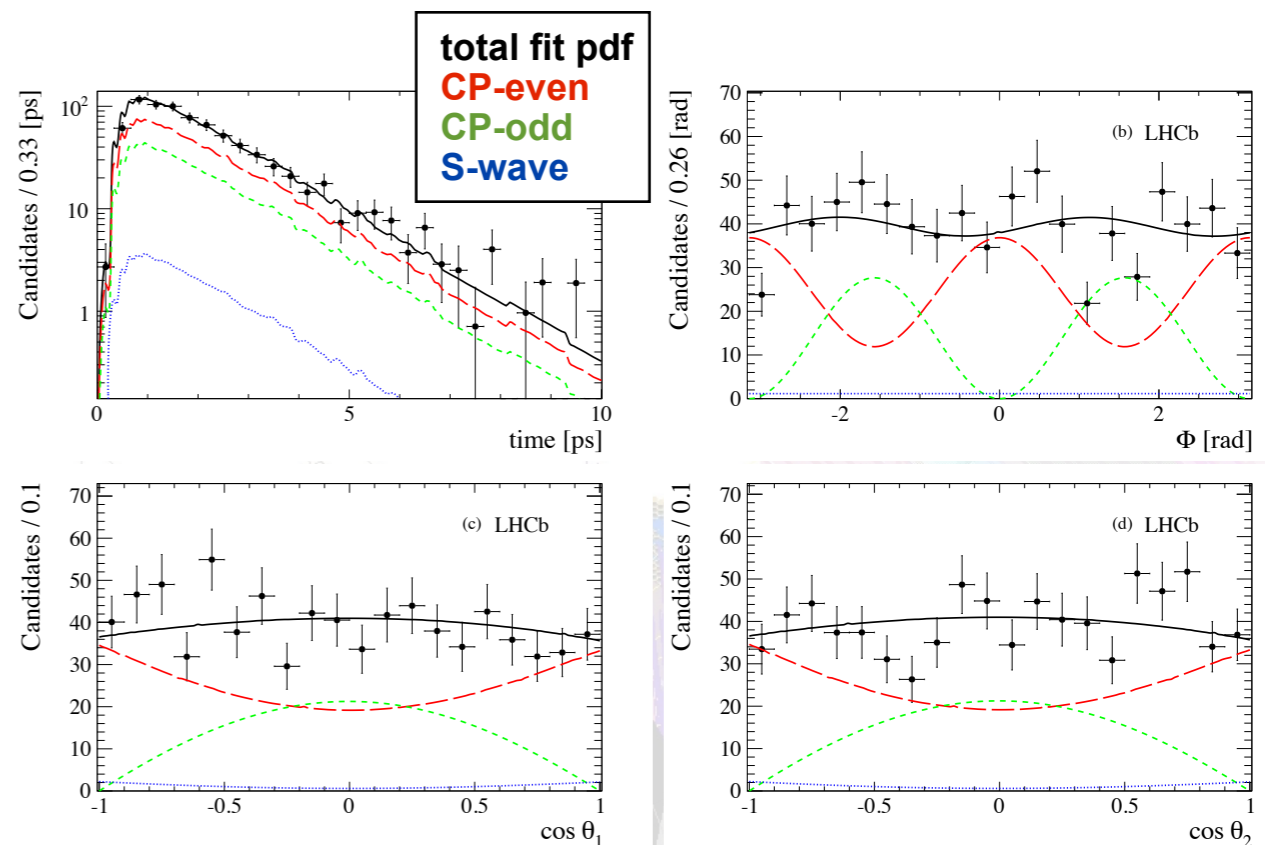
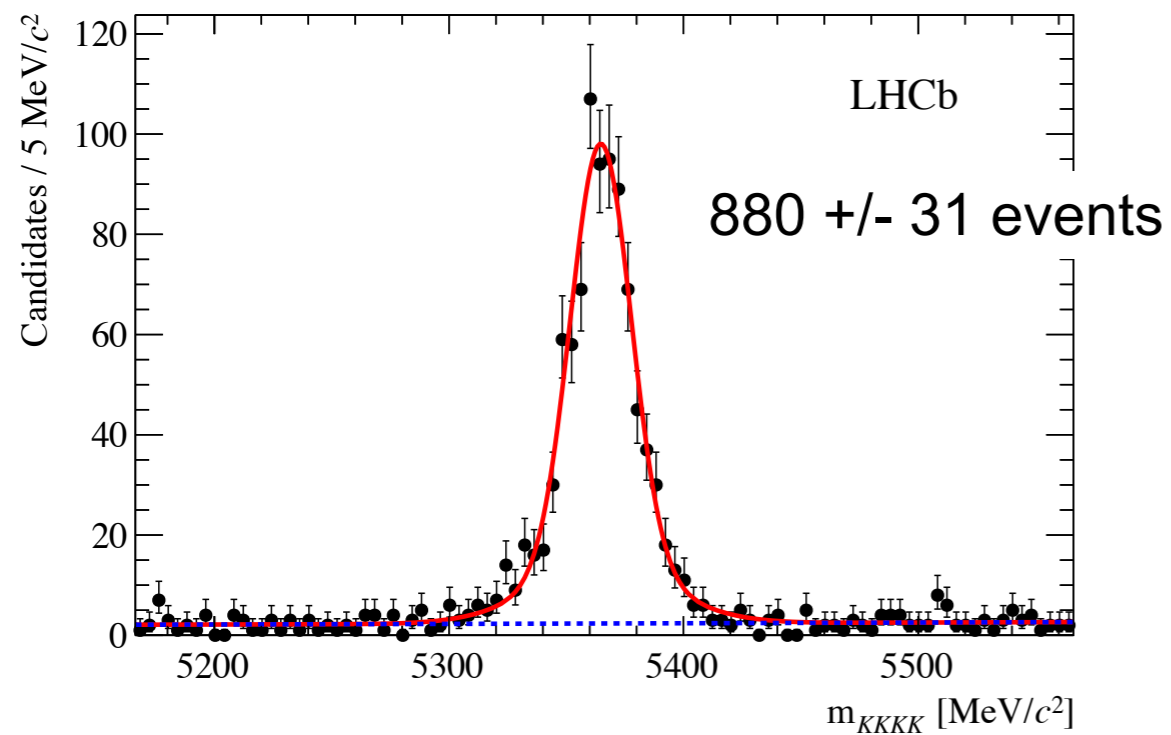
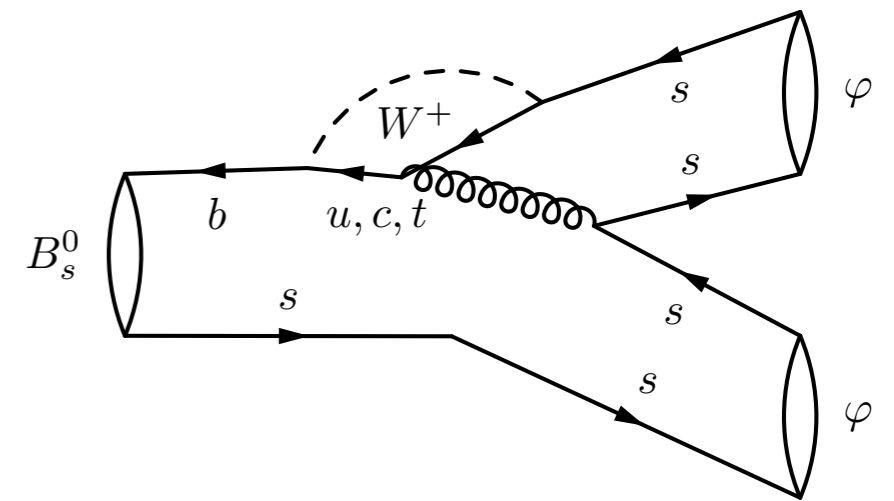
$$\frac{\mathcal{B}(B^0 \rightarrow K^{*0} \gamma)}{\mathcal{B}(B_s^0 \rightarrow \phi \gamma)} = 1.23 \pm 0.06 \text{ (stat.)} \pm 0.04 \text{ (syst.)} \pm 0.10 \text{ (} f_s/f_d \text{)},$$

$$\mathcal{A}_{CP}(B^0 \rightarrow K^{*0} \gamma) = (0.8 \pm 1.7 \text{ (stat.)} \pm 0.9 \text{ (syst.)})\%.$$

note: with 3/fb should have enough events for tagged time-dependent CPV analysis in  $B_s \rightarrow \phi \gamma$

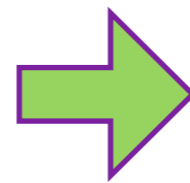
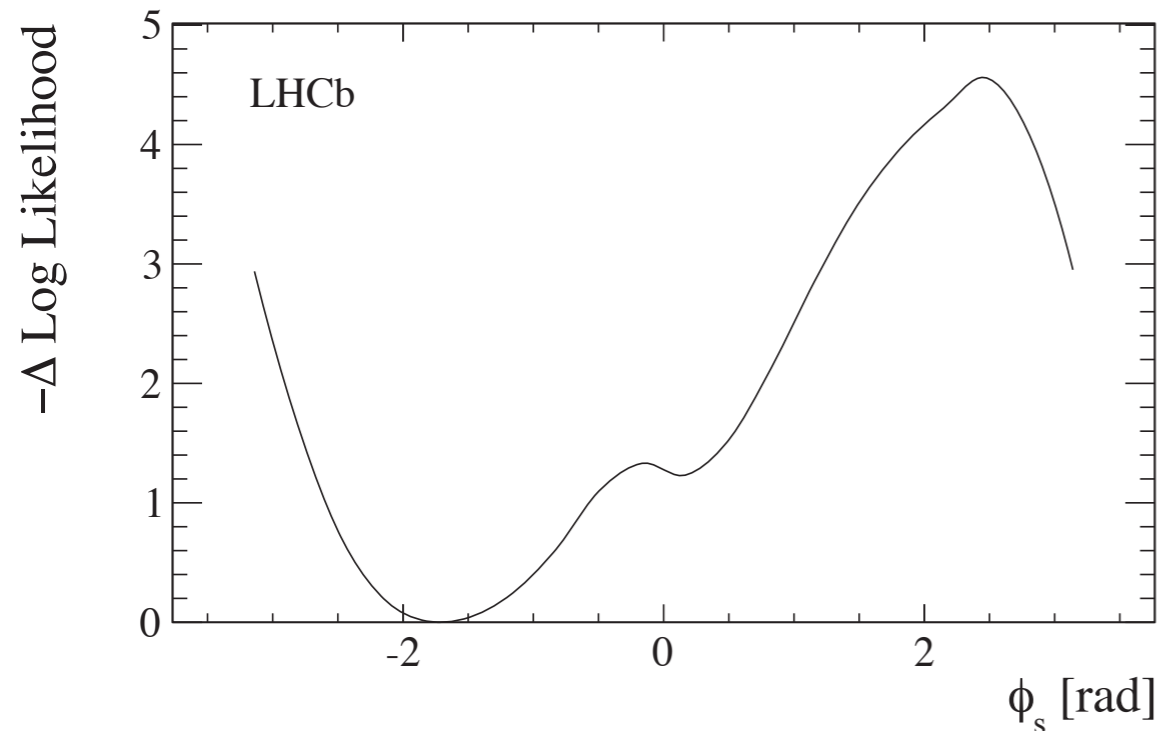
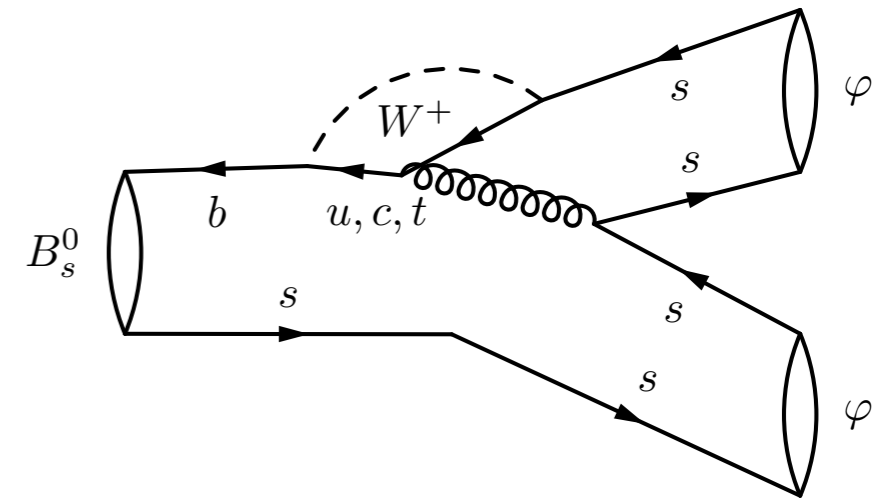
# $B_s \rightarrow \varphi\varphi$

- 'golden penguin' mode: similar CKM phases as  $J/\psi\varphi$ , but only penguin amplitude
- expect very small CPV phase in SM:  $\phi_s^{s\bar{s}s} \approx 0.01$   
[see e.g. Bartsch, Buchalla, and Kraus, 2008]
- requires very similar tagged time-dependent angular analysis



# $B_s \rightarrow \phi\phi$

- ‘golden penguin’ mode: similar CKM phases as  $J/\psi\phi$ , but only penguin amplitude
- expect very small CPV phase in SM:  $\phi_s^{s\bar{s}s} \approx 0.01$   
[see e.g. Bartsch, Buchalla, and Kraus, 2008]
- requires very similar tagged time-dependent angular analysis
- result for the CPV phase



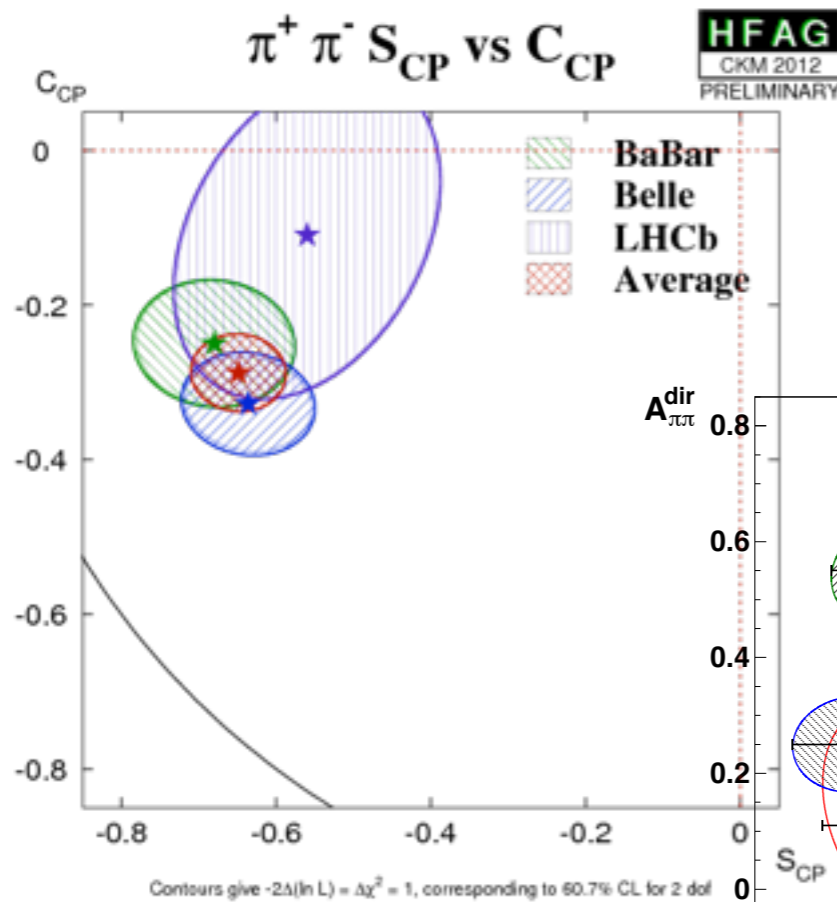
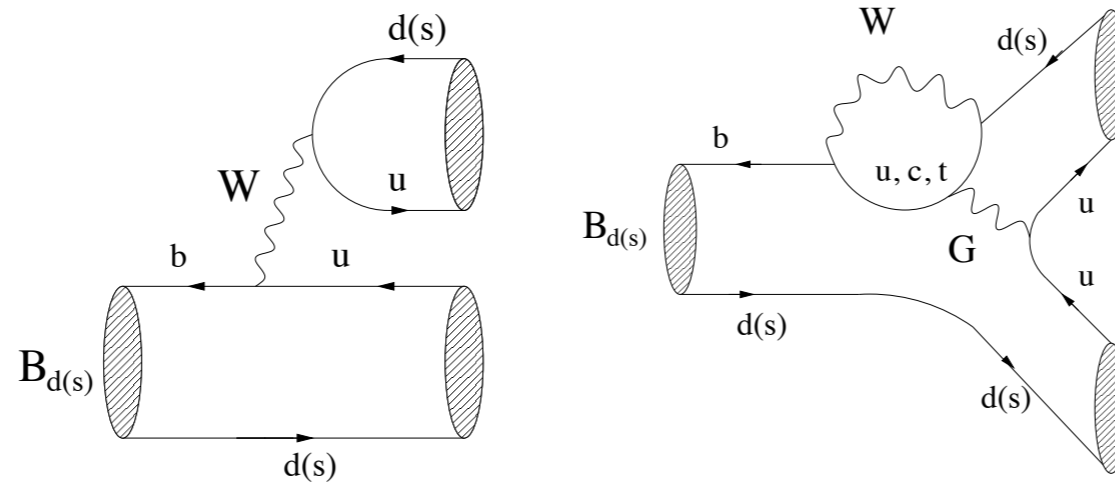
LHCb-PAPER-2013-007, *preliminary*

$$\phi_s^{s\bar{s}s} \in [-2.46, -0.76] \quad \text{at 68\% CL}$$

First measurement of  $B_s$  mixing phase in a pure penguin decay



two amplitudes, with different weak phases: also expect large direct CPV ( $C \neq 0$ )



DLL = 2 contours (“39%CL”)

