the flipped classroom: a fresh look at teaching







how do we teach science?

The Hungry Stones and Other Stories

Rabindranath Tagore

we repeat the book

why do we teach the way we do?





inertia

is this teaching efficient?



information loss





information loss &

we don't know their problems

they all know it

Dear professor,

Due to an appointment that was fixed long ago, I will not be able to attend the first lecture of your course next week.

Do I miss something important if I skip this first lecture?

Yours sincerely,

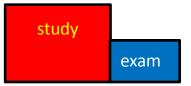
can we 1mprove

ideal world



lecture study

lecture study

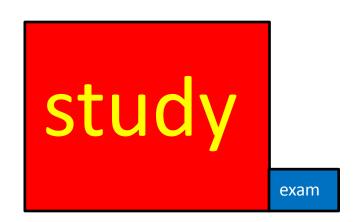


real world



lecture

lecture

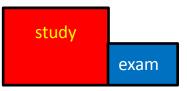


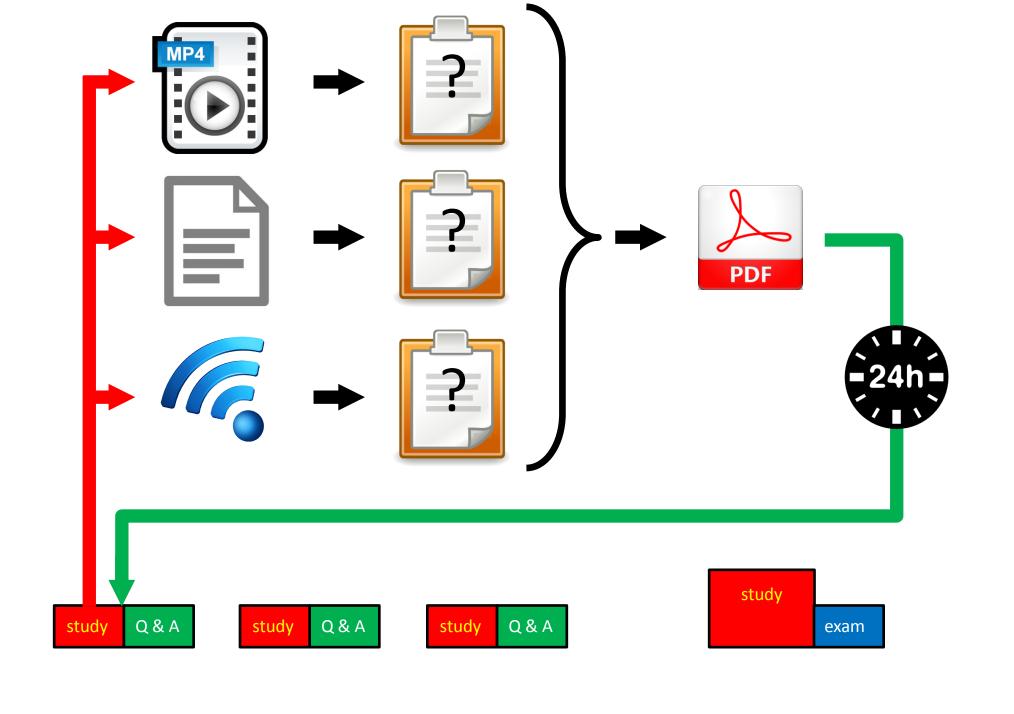
flipped classroom



study Q&A

study Q&A



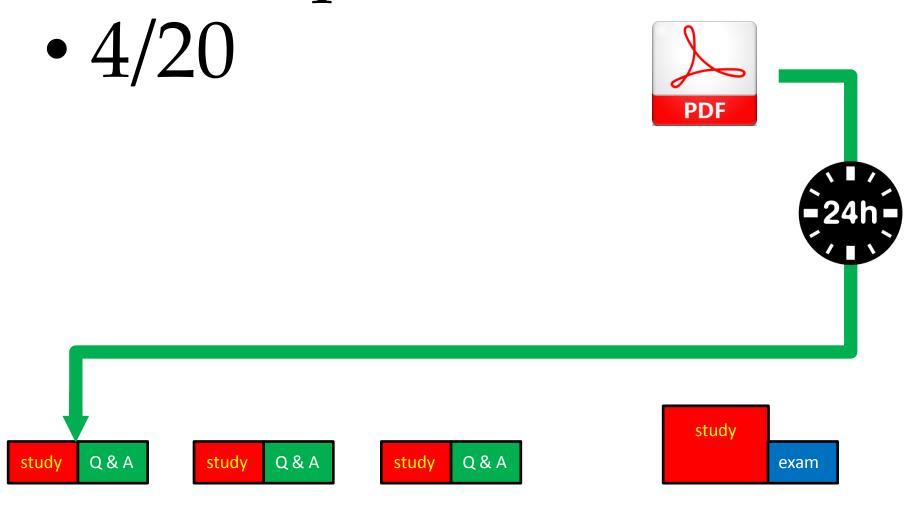




avoid covering pre-class material once again in class



• social pressure



pre-class material







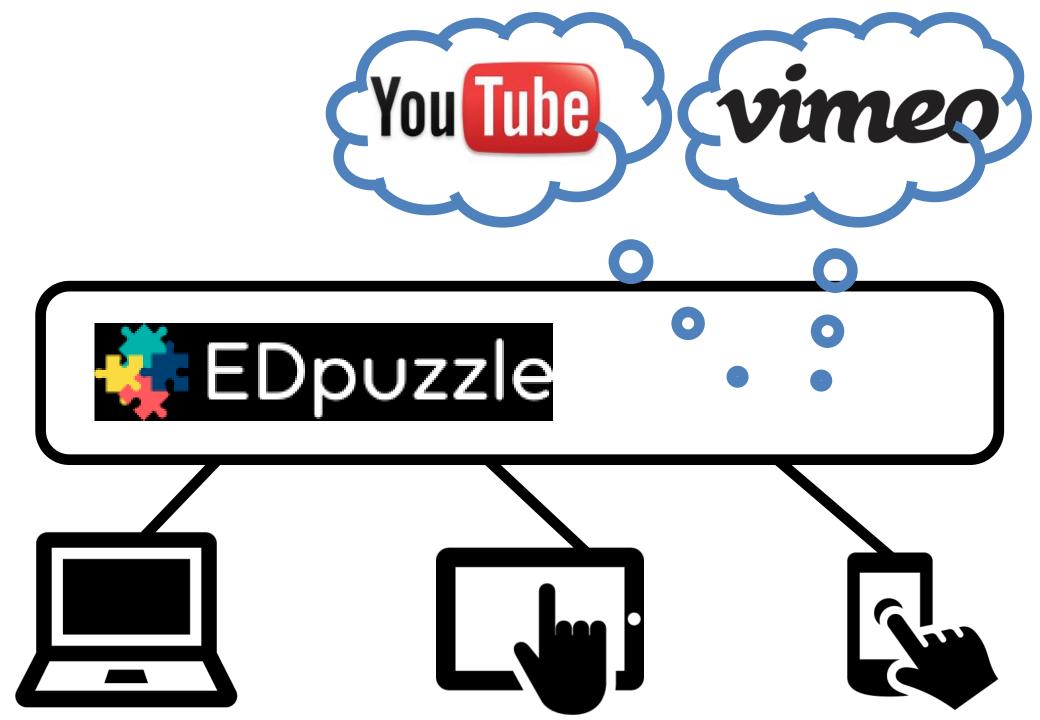


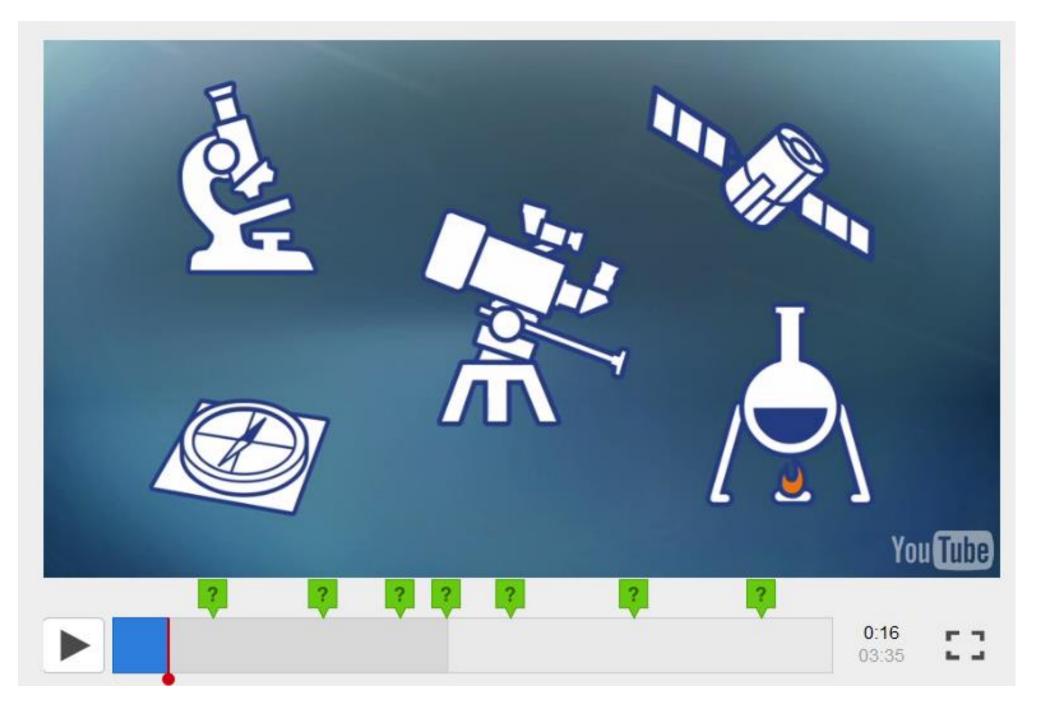
videofragment 1

keep videos short: 6-9 minutes

- attention curve
- easier to change







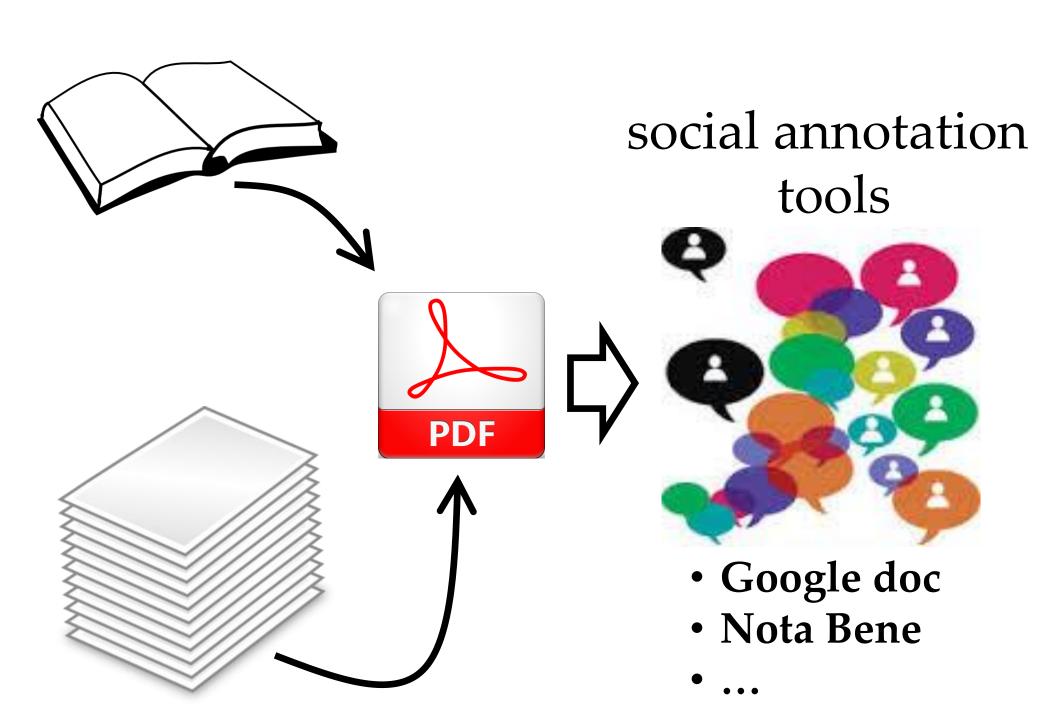


pre-class material









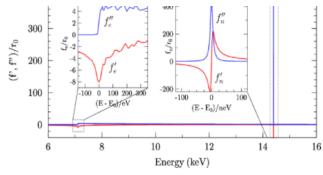


Fig. 2. Anomalous scattering lengths (real part f and imaginary part f') for 57 Fe including the K-edge around 7 keV and the nuclear resonance around 14.4 keV, measured in units of the free electron radius r₀. While around the K-edge the anomalous contributions amount to only a few electrons, around the nuclear resonance they correspond to more than 200 electrons!

 $f_{LM} \approx 0.8$ so that $f_0 \approx 1.2 \times 10^{-12}$ m. In the case of magnetic materials, a finite magnetic hyperfine field acts on the resonant nuclei. In ferromagnetic α -Fe, for example, the magnitude of this field is 33.3T at room temperature. As a consequence, the degeneracy of nuclear levels with spin $I_N \geq 1/2$ is lifted, i.e., these states split into $2I_N + 1$ magnetic sublevels between which electromagnetic transitions can occur according to the multipole selection rules. For an electric dipole transition the nuclear scattering length reads:

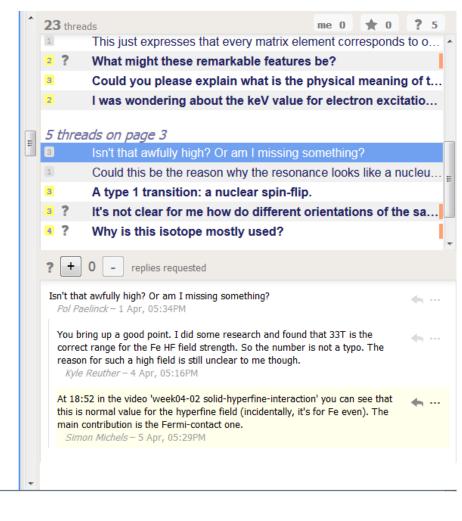
$$[N(\omega)]_{\mu\nu} = \frac{3}{16\pi} \left\{ \begin{aligned} &(\boldsymbol{\varepsilon}_{\mu} \cdot \boldsymbol{\varepsilon}_{\nu})[F_{+1} + F_{-1}] \\ &-i(\boldsymbol{\varepsilon}_{\mu} \times \boldsymbol{\varepsilon}_{\nu}) \cdot \mathbf{m}[F_{+1} - F_{-1}] \\ &+(\boldsymbol{\varepsilon}_{\mu} \cdot \mathbf{m})(\boldsymbol{\varepsilon}_{\nu} \cdot \mathbf{m})[2F_{0} - F_{+1} - F_{-1}] \end{aligned} \right\}$$

The quantities $F_0=F_0(\omega)$ and $F_\pm=F_\pm(\omega)$ are functions that describe the energy dependence of the scattering length for the transitions with change of magnetic quantum number $\Delta m=0,\pm 1$ [13]. In case of a magnetic dipole transition (as for the 14.4 keV transition of 57 Fe) the role of electric and magnetic fields of the radiation are interchanged. Then one has to transform the polarization vectors in the expression above via $\mathbf{E} \to \mathbf{E} \times \mathbf{k}_0$.

The scalar and vector products involving the polarization basis vectors $\mathbf{\epsilon}_{\mu}$ and $\mathbf{\epsilon}_{\nu}$ as well as the unit vector \mathbf{m} in the direction of the local magnetization (corresponding to the direction of the local hyperfine field at the position of the nucleus) indicate a strong

| | Geometry | Scattering matrix | Time spectrum |
|--------------------|--------------------------|---|---------------|
| A | $m \stackrel{k_0}{\sim}$ | $\begin{pmatrix} F_{+1} + F_{-1} & -i(F_{+1} - F_{-1}) \\ i(F_{+1} - F_{-1}) & F_{+1} + F_{-1} \end{pmatrix}$ | MMMMM |
| В | m k ₀ | $\begin{pmatrix} F_{+1} + F_{-1} & 0 \\ 0 & F_{+1} + F_{-1} \end{pmatrix}$ | MMMMMM |
| С | k_0 m | $\begin{pmatrix}F_{+1}+F_{-1}&0\\0&2F_0\end{pmatrix}$ | MMMMMM |
| D | m k_0 σ | $\begin{pmatrix} 2F_0 & 0 \\ 0 & F_{+1} + F_{-1} \end{pmatrix}$ | MMMM |
| $0 	 t/\tau_0 	 1$ | | | |

Fig. 3. Time spectra of nuclear resonant scattering for selected orientations of the magnetic hyperfine field m relative to the incident wave vector \mathbf{k}_0 . The 2×2 matrix of the nuclear scattering length is given in a linear polarization basis (σ, π) . The time spectra are calculated for a 2 nm thick, ferromagnetic ω^{-5} Fe film on a tungsten substrate, assuming purely σ polarized incident radiation and unpolarized detection. This is the most frequently used scattering geometry in experiments with synchrotron radiation. A. C. D display the results for a undirectional magnetization of the sample. B results from the superposition of two magnetic sublattices in



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ideal for existing courses with good text-based material

pre-class material











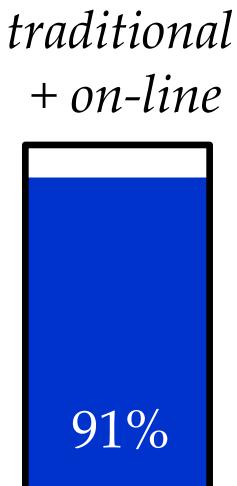
show me some results!

learning gain

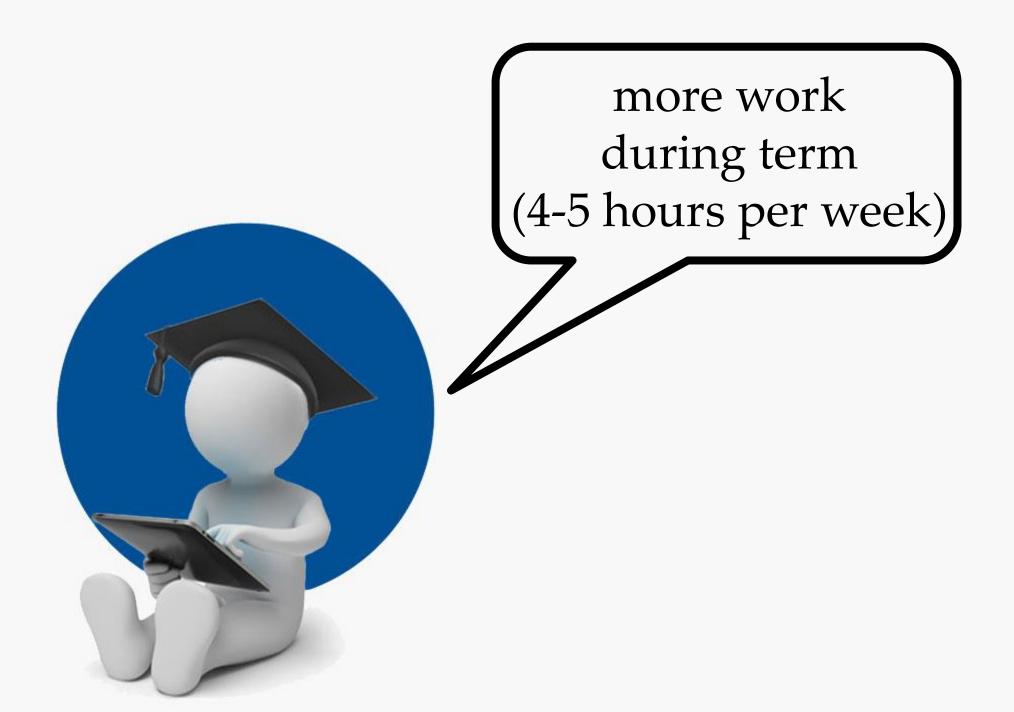


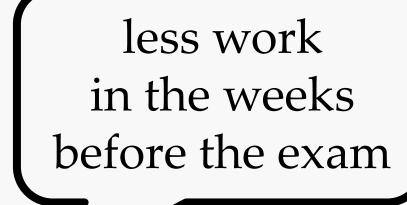
San Jose State and edX (fall 2012)

traditional

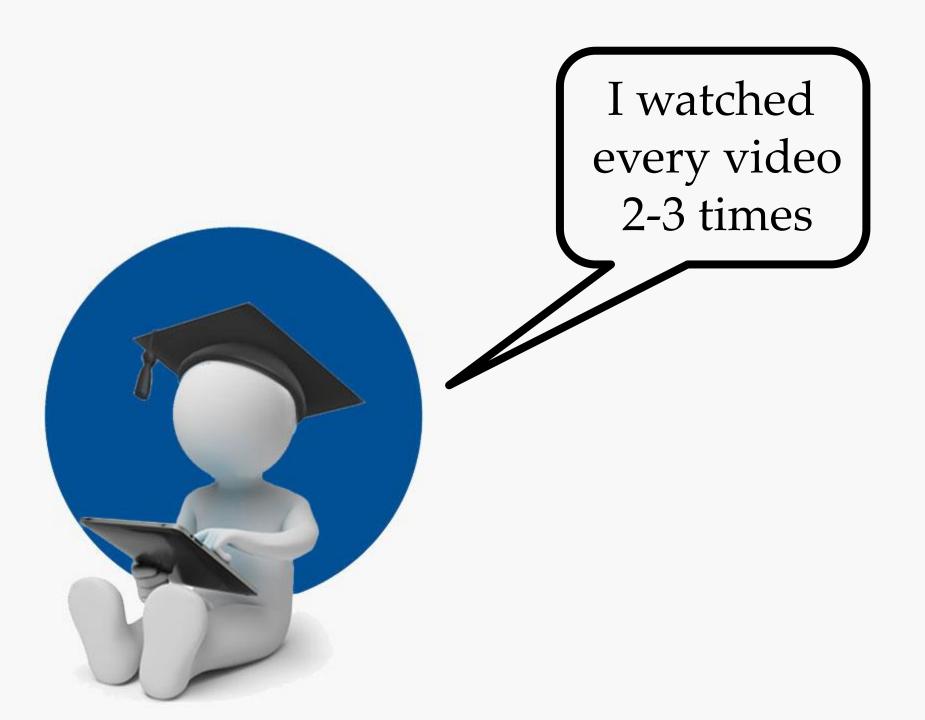


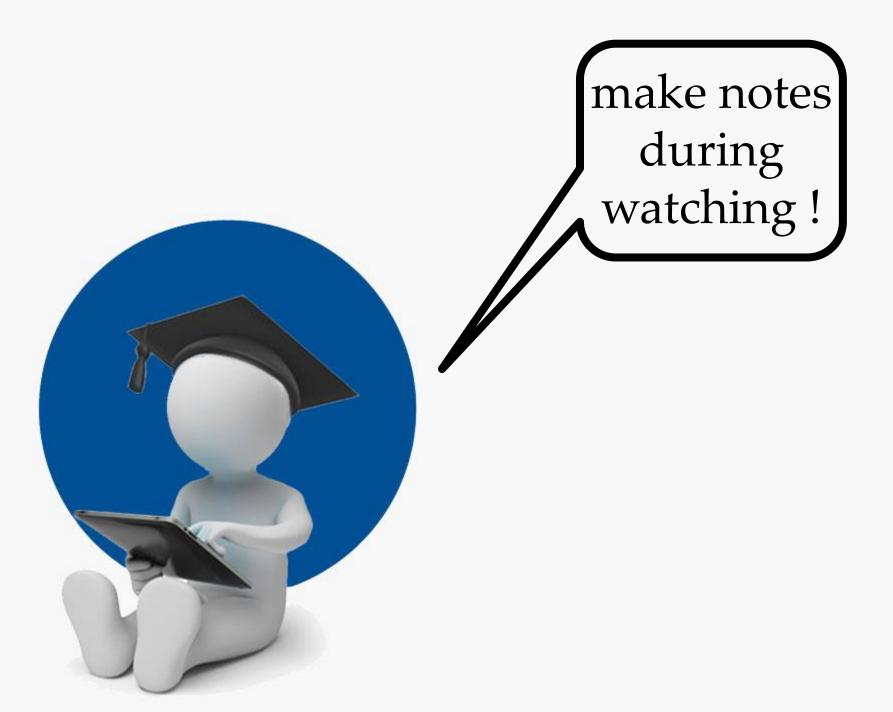
study behaviour













Computational Materials Physics @ TIFR (August 2015)

appreciation

It's my 5th year as a student, and never I was as well-prepared for an exam



It's a lot of work, but you benefit from it



At home, your thoughts more easily wander. But the same holds true in a lecture, and there you can't rewind.



time pressure if all teachers would do this?





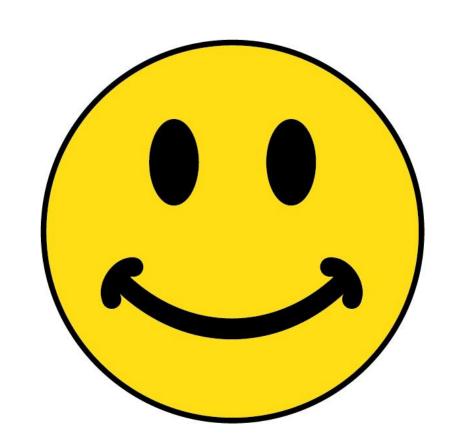
economize on class time!

teacher's time

You don't gain time



You don't gain time, but you spend it better.



something for me?



each year, I tell my same story



my students are not responsive at all



for some students my course goes too fast, for others it goes too slowly





my talks on YouTube http://goo.gl/P2b1Hs

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