

# D-meson semileptonic decays with Lattice QCD

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# 1. Introduction

Goals in the study of semileptonic  $D$  decays

- \* Precise determination of CKM matrix elements ( $|V_{cd,cs}|$ )

$$\text{Experiment} = (\text{known factors}) \times (V_{CKM}) \times \underbrace{(\text{hadronic matrix elements})}_{\text{lattice QCD}}$$

- \* Check Standard Model
  - \*\* Consistency of different determinations of CKM matrix elements
  - \*\* Test unitarity of CKM matrix in the second row.
  - \*\* Comparison of shape of form factors with experimental data.
- \* Validate lattice QCD techniques to use in  $B$  physics
- \* Good candidate for New Physics searches (constraining NP models)
  - \*\* Correlated signals of NP in leptonic and semileptonic decays.

See **S. Fajfer talk**

# 1. Introduction: Lattice QCD

**Lattice QCD:** Numerical evaluation of QCD path integral (rely only on first principles).

\* Control over systematic errors:

\*\* Unquenched calculations:  $N_f = 2$ ,  $\mathbf{N}_f = \mathbf{2} + \mathbf{1}$  or  $\mathbf{N}_f = \mathbf{2} + \mathbf{1} + \mathbf{1}$ .

\*\* Discretization: improved actions + simulations at several lattice spacings  $a$ 's  
→ continuum limit.

\*\* Chiral extrapolation: simulate at several  $m_\pi$  and extrapolate to  $m_\pi^{\text{phys}}$

Next step: configurations with **physical light quark masses**.

Five collaborations have generated sets of configurations with physical light quark masses (**PACS-CS, BMW, MILC, RBC/UKQCD, ETM**)

\*\* Renormalization: non-perturbative, perturbative.

\*\* Tuning lattice scale and quark masses  $M_{H,lat} = M_{H,exp} \rightarrow m_f = m_{f,phys}$

\*\* Finite volume, **isospin effects, electromagnetic effects, ...**

# 1. Introduction: Lattice QCD

Several extrapolations/interpolations: lattice spacing, lattice volume, quark masses.

**Effective Field Theory:** good tool for

- \* Provides functional form for extrapolation (or interpolation).
- \* Used to build improved lattice actions/methods
- \* Used to anticipate size of systematic effects (discretization, FV, chiral extrap. ...)

**Charm: heavy or light quark?**

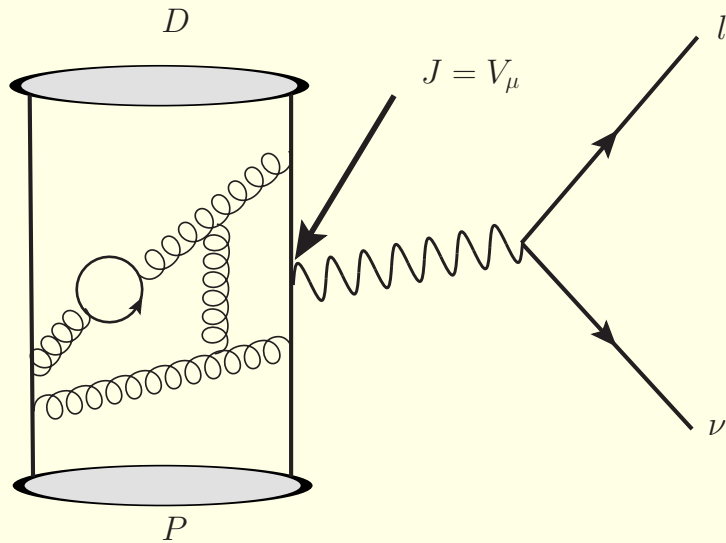
Dominant discretization errors for light quarks are  $\mathcal{O}(\alpha_s^k (a\Lambda_{QCD})^n)$

Dominant discretization errors for heavy quarks are  $\mathcal{O}(\alpha_s^k (am_h)^n)$

( $am_c \sim 0.15 - 0.6$  in current lattice spacings)

For charm use light quark methods  
with improved actions (HISQ, tmWilson, NP imp. Wilson ...)

## 2. Semileptonic $D$ decays



$$P = \pi, K$$

$x = d, s$  daughter light quark

$q = (p_D - p_P)$  (momentum of lepton pair)

$$\underbrace{\frac{d\Gamma(D \rightarrow Pl\nu)}{dq^2}}_{\text{experimental}} = \frac{G_F^2}{24\pi^3} \frac{(q^2 - m_l^2)^2 \sqrt{E_P^2 - m_P^2}}{q^4 m_D^2} |V_{cx}|^2$$

$$\left[ \left(1 + \frac{m_l^2}{2q^2}\right) m_D^2 (E_P^2 - m_P^2) \underbrace{|f_+(q^2)|^2}_{\text{lattice QCD}} + \frac{3m_l^2}{8q^2} (m_D^2 - m_P^2)^2 \underbrace{|f_0(q^2)|^2}_{\text{lattice QCD}} \right]$$

With vector and scalar form factors  $f_+(q^2)$  and  $f_0(q^2)$  defined by

$$\langle P(p_P) | V_\mu | D(p_D) \rangle = \left( p_{P\mu} + p_{D\mu} - \frac{m_D^2 - m_P^2}{q^2} q_\mu \right) f_+(q^2) + \frac{m_D^2 - m_P^2}{q^2} q_\mu f_0(q^2)$$

## 2. Semileptonic $D$ decays

For  $l = e, \mu$  the contribution from  $f_0(q^2)$  can be neglected and

$$\underbrace{\frac{d\Gamma(D \rightarrow Pl\nu)}{dq^2}}_{\text{experimental}} = \frac{G_F^2}{24\pi^3} |\vec{p}_P|^3 |V_{cx}|^2 \underbrace{|f_+^{DP}(q^2)|^2}_{\text{lattice QCD}}$$

The errors on those studies are still dominated by errors in the calculation of the relevant form factors.

$$\frac{d}{dq^2} \Gamma(D \rightarrow K(\pi)l\nu) \propto |V_{cs(cd)}|^2 |f_+^{D \rightarrow K(\pi)}(q^2)|^2$$

HFAG2016 1.0(2.6)% error 5(8.7)% error

(in leptonic determinations of  $|V_{cd(cs)}|$  the main error is experimental)

See [A. Soffer's talk](#)

## 2. Semileptonic $D$ decays: Form factors

For  $l = e, \mu$  the contribution from  $f_0(q^2)$  can be neglected and

$$\underbrace{\frac{d\Gamma(D \rightarrow Pl\nu)}{dq^2}}_{\text{experimental}} = \frac{G_F^2}{24\pi^3} |\vec{p}_P|^3 |V_{cx}|^2 \underbrace{|f_+^{DP}(q^2)|^2}_{\text{lattice QCD}}$$

With lattice QCD we can:

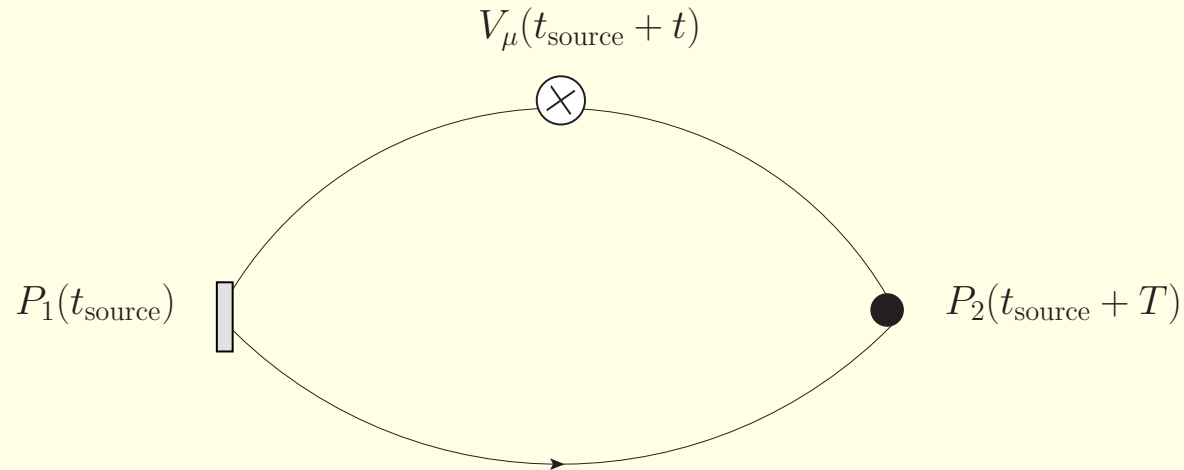
- \* Calculate both  $f_+^{DP}(q^2)$  and  $f_0^{DP}(q^2)$  for the entire  $q^2$  range.
- \* Extend to rare semileptonic decay form factors ( $f_T^{DP}(q^2)$ ).

Use **z-expansion** for model-independent parametrization of  $q^2$  dependence

→ compare shape of lattice form factor with experimental data

### 3. Semileptonic $D$ decays with Lattice QCD

Two main strategies to eliminate the need of renormalize the lattice currents



# Double ratios of 3-point correlators **Becirevic, Haas, Mescia 0710.1741**  
 (get the form factors from linear combinations of double ratios, with  $P = \pi, K$ )

$$R_\mu = \frac{C_{3pt,\mu}^{DP}(t, T, \vec{p}_D, \vec{p}_P) C_{3pt,\mu}^{PD}(t, T, \vec{p}_P, \vec{p}_D)}{C_{3pt,\mu}^{PP}(t, T, \vec{p}_P, \vec{p}_P) C_{3pt,\mu}^{DD}(t, T, \vec{p}_D, \vec{p}_D)}$$

$$\xrightarrow{t+t_{\text{source}} \gg a, (T-t-t_{\text{source}}) \gg a} \frac{1}{4} |\langle P(p_P) | V_\mu | D(p_D) \rangle|^2 (p_D)_\mu (p_P)_\mu$$



### 3. Semileptonic $D$ decays with Lattice QCD

Two main strategies to eliminate the need of renormalize the lattice currents

# Use the Ward identity ( $S = \bar{x}c$ ) **HPQCD**, Phys.Rev.D82:114506(2010)

$$q^\mu \langle P | V_\mu^{cont.} | D \rangle = (m_c - m_x) \langle P | S^{cont.} | D \rangle$$

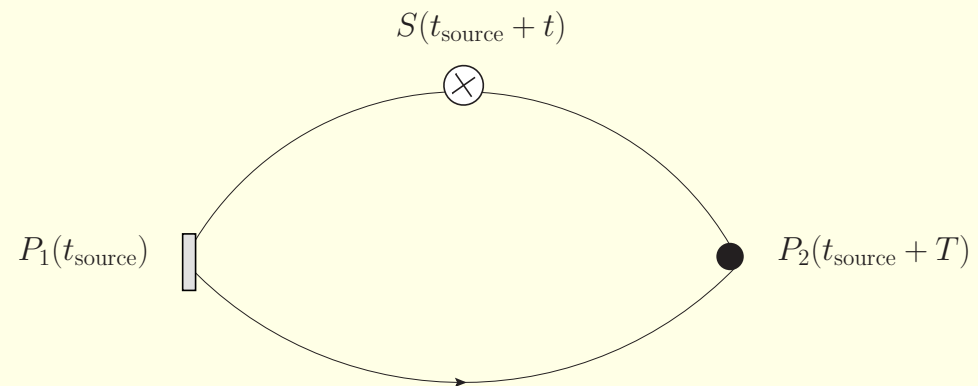
that relates matrix elements of vector and scalar currents. In the lattice

$$q^\mu \langle P | V_\mu^{lat.} | D \rangle Z = (m_c - m_x) \langle P | S^{lat.} | D \rangle$$

→ replace the  $V_\mu$  with a scalar current in the 3-point function

$$f_0^{DP}(q^2) = \frac{m_c - m_x}{m_D^2 - m_P^2} \langle P | S | D \rangle_{q^2} \implies$$

$$f_+^{DP}(0) = f_0^{DP}(0) = \frac{m_c - m_x}{m_D^2 - m_P^2} \langle P | S | D \rangle_{q^2=0}$$



### 3. Semileptonic $D$ decays with Lattice QCD

**2010/2011:** Important reduction of errors in the lattice determination of the form factors  $f_+^{D\pi(K)}(0)$  by  $N_f = 2 + 1$  **HPQCD Collaboration** calculations, due mainly to

\* Use a relativistic action, **HISQ**, to describe light and charm quarks.

\* Used WI to relate scalar matrix elements to vector matrix element

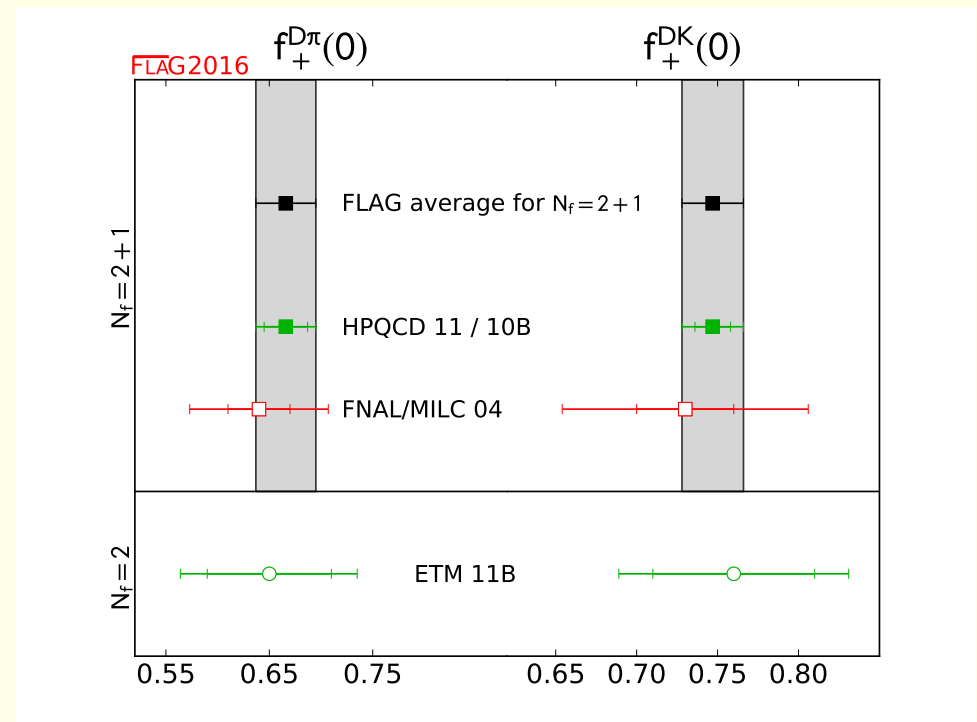
**HPQCD**, 1008.4562, 1109.1501

$$f_+^{D\pi}(0) = 0.666(29)$$

$$f_+^{DK}(0) = 0.747(19)$$

**FLAG 2016 averages** 1607.00299,

<http://itpwiki.unibe.ch/flag/>



Using **HFGA2014** experimental averages and **HPQCD** form factors above:

$$|V_{cd}| = 0.2140(29)_{\text{exp}}(93)_{\text{lat}} \quad |V_{cs}| = 0.975(7)_{\text{exp}}(25)_{\text{lat}}$$

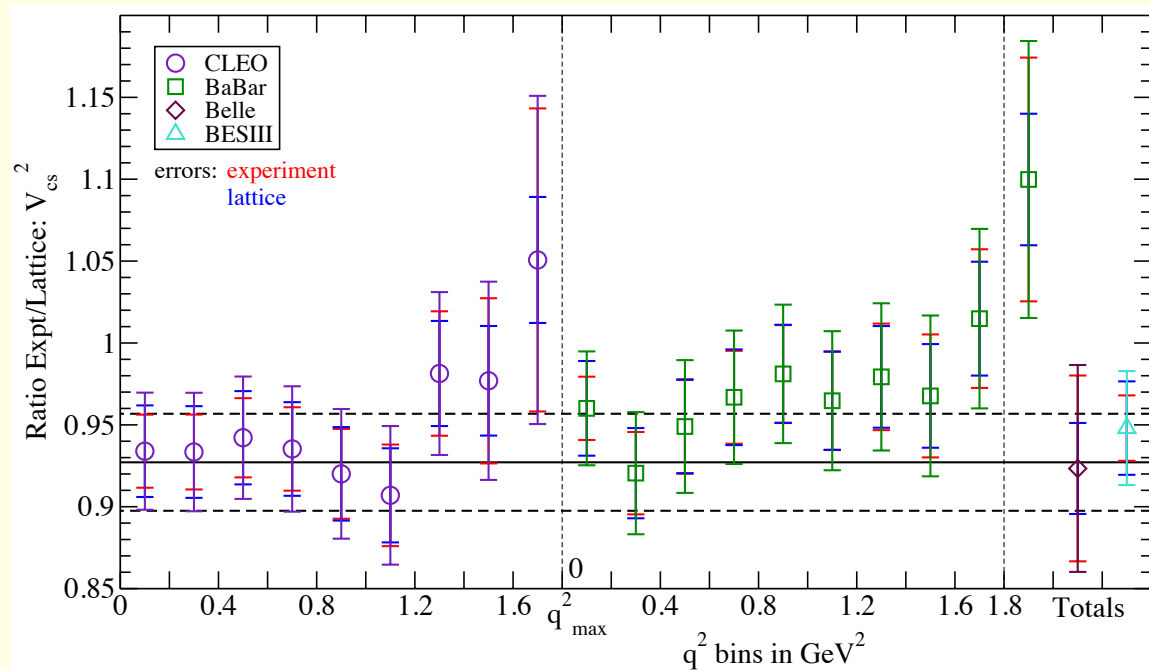
Error dominated by form factor (lattice) uncertainty. Main sources: statistics and  $am_c$  disc.

### 3. Done but not published: $q^2 \neq 0$

$N_f = 2 + 1$  determination of  $|V_{cs}|$  from  $D \rightarrow Kl\nu$  at non-zero momentum transfer **HPQCD, 1305.1462**

Calculation of  $f_0^{DK}(q^2)$  (using Ward identity method) and  $f_+^{DK}(q^2)$  (using its definition)

- \* Global fit to available experimental data (using  $z$ -expansion)  $\rightarrow$  extraction of  $|V_{cs}|$  using all experimental  $q^2$  bins  $\rightarrow$  large reduction of error

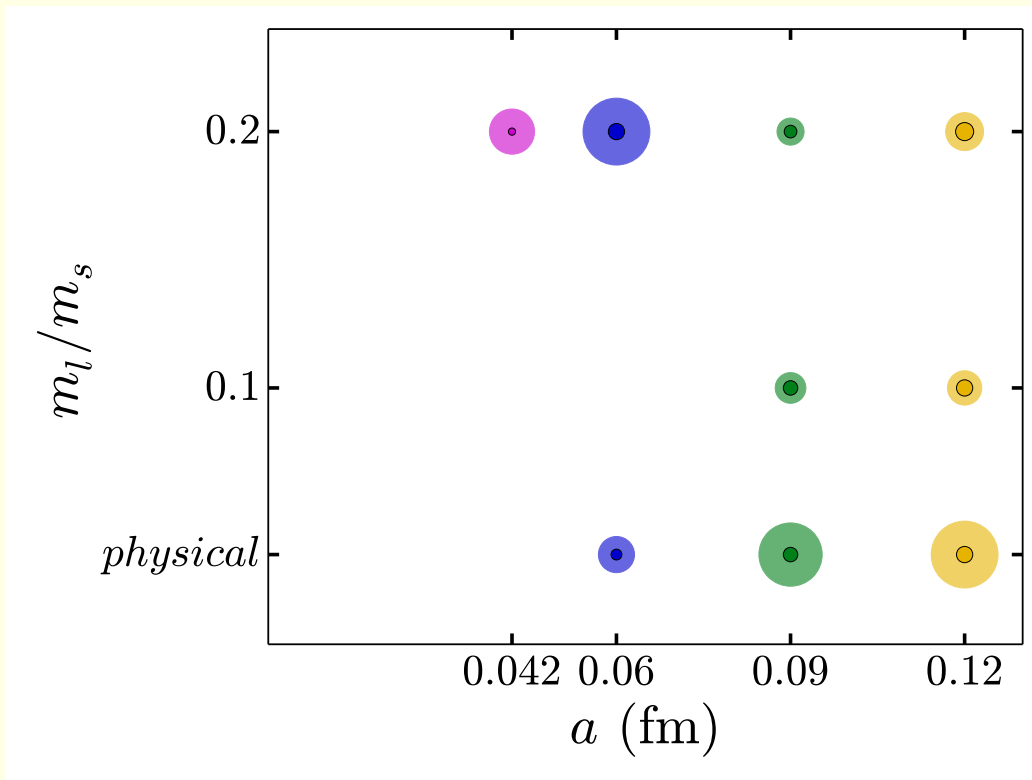


$$|V_{cs}| = 0.963(5)_{exp}(14)_{lat}$$

Lattice form factors are more precise at smaller external momenta (near/at  $q^2_{\max}$ )

## 4. Work in progress: $q^2 = 0$

FNAL/MILC, talk by S. Gottlieb (T. Primer) at Lattice 2016  $N_f = 2 + 1 + 1$



- \* **HISQ** ensembles including
- \*\* Charm quarks on the sea
- \*\* Three set of data with physical light quark masses
- \* Simulate directly at  $q^2 = 0$

Use the relation  $f_+^{DP}(0) = f_0^{DP}(0) = \frac{m_c - m_x}{m_D^2 - m_P^2} \langle P|S|D \rangle_{q^2=0}$

$\rightarrow f_+^{D\pi}(0)$  and  $f_+^{DK}(0)$

## 4. Work in progress: $q^2 = 0$

FNAL/MILC, talk by S. Gottlieb at Lattice 2016  $N_f = 2 + 1 + 1$

\* Continuum-chiral extrapolation in Heavy Meson ChPT framework, including leading order discretization effects, in the hard pion/kaon limit.

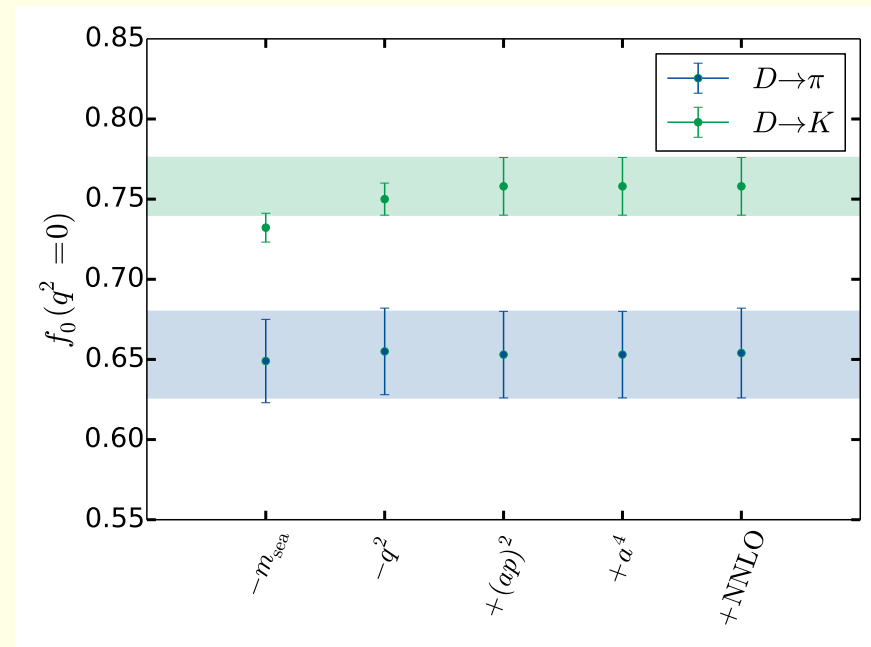
Aubin, Bernard 0704.0795, Becirevic, Prelovsek, Zupan hep-lat/0305001, Bijens, Jemos 1006.1197

$$f_0(q^2 = 0) = \frac{C_0}{f_\pi} \left[ 1 + \text{logs}(a, m_v, m_{\text{sea}}) + C_v m_v + C_s m_{\text{sea}} + C_a a^2 + C_q q^2 \right]$$

Very stable (continuum-chiral)  
fits under variations of fit function

\* Estimate of finite volume effects: negligible correction

Preliminary systematic error analysis:  $\sim 4.1\%$  and  $\sim 2.4\%$  for  $f_+^{D\pi}(0)$  and  $f_+^{DK}(0)$



## 4. Work in progress: $q^2 = 0$

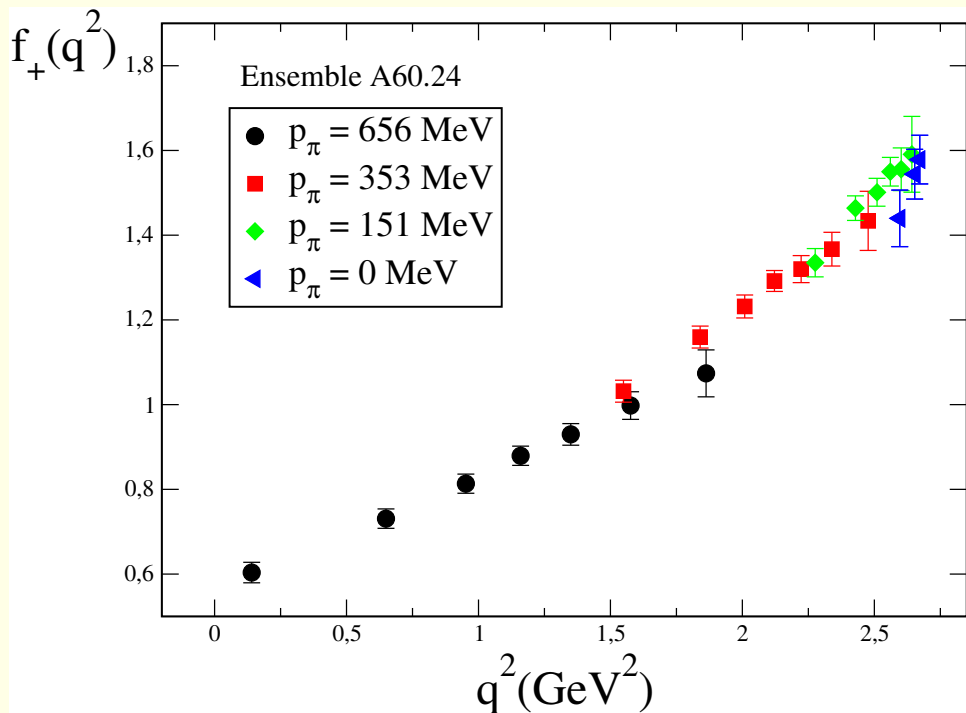
FNAL/MILC, talk by S. Gottlieb at Lattice 2016  $N_f = 2 + 1 + 1$

**Future work:** Scalar and vector form factors at multiple  $q^2$

→ combine with experiment to reduce errors. Compare shape.

## 4. Work in progress: $q^2 \neq 0$

ETM, talk by G. Salerno at Lattice 2016, 1611.00022  $N_f = 2 + 1 + 1$



- \* **tmWilson** ensembles including
- \*\* Charm quarks on the sea
- \*\* Pion masses in the range [210 – 450] MeV
- \* Calculate  $f_+^{D\pi}$  and  $f_0^{D\pi}$  over whole  $q^2$  range.
- \* Three different lattice spacings.
- \* Use double ratios to avoid renormalization.

Global fit to all  $f_+(q^2)$ ,  $f_0(q^2)$  data  $\rightarrow$  study simultaneously  $q^2$ ,  $m_l$  and  $a^2$  dependence.

\* Modified **z-expansion** **Bourenly, Caprini, Lellouch 0807.2722** imposing  $f_+(0) = f_0(0) = f(0)$  (parametrization of  $f(0)$  inspired by hard pion  $SU(2)$  ChPT)

\* Add finite volume term in the fit  $K_{FSE}^{+(0)} = 1 + C_{FSE}^{+(0)} \xi_l e^{-m_\pi L} / (m_\pi L)$

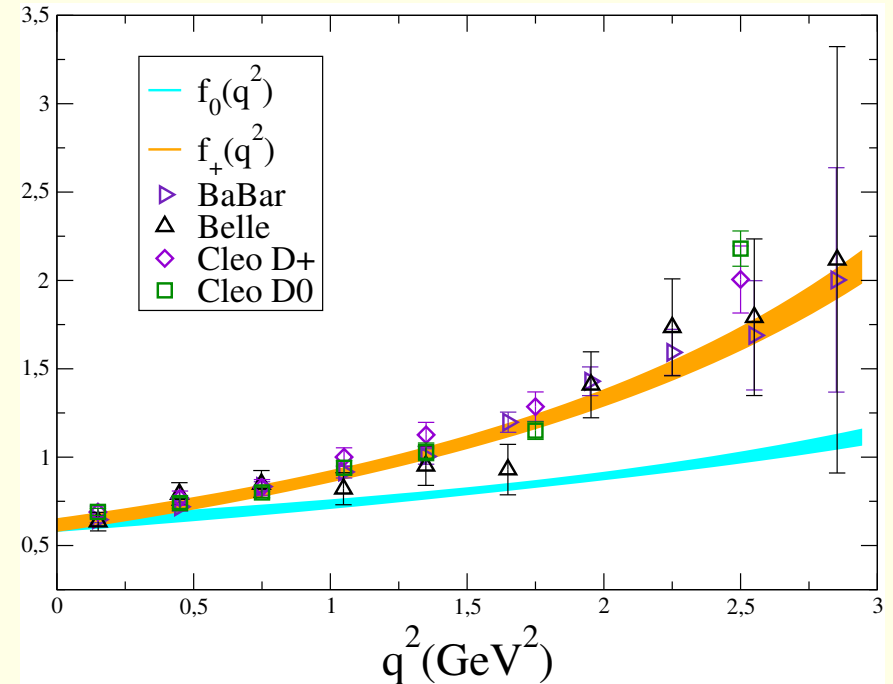
## 4. Work in progress: $q^2 \neq 0$

ETM, talk by G. Salerno at Lattice 2016, 1611.00022  $N_f = 2 + 1 + 1$

\*  $z_0 \equiv z(q^2 = 0)$

\* parametrization of  $f(0)$  inspired by hard pion  $SU(2)$  ChPT (with  $\xi_l \propto m_l$ )

$$f(0) = b_0 \left[ 1 - \frac{3}{4}(1 + 3g^2)\xi_l \log \xi_l + b_1 \xi_l + b_2 a^2 \right]$$



$$f_{+(0)}(q^2) = \frac{f(0) + c_{+(0)}(a^2)(z - z_0)[1 + (z - z_0)/2]}{1 - K_{FSE}^{+(0)} q^2 / M_{V(S)}^2}$$

**Preliminary systematic error analysis:**

$$f_+^{D\pi}(0) = 0.631(37)_{stat}(14)_{chiral}(08)_{disc} = 0.631(40)$$

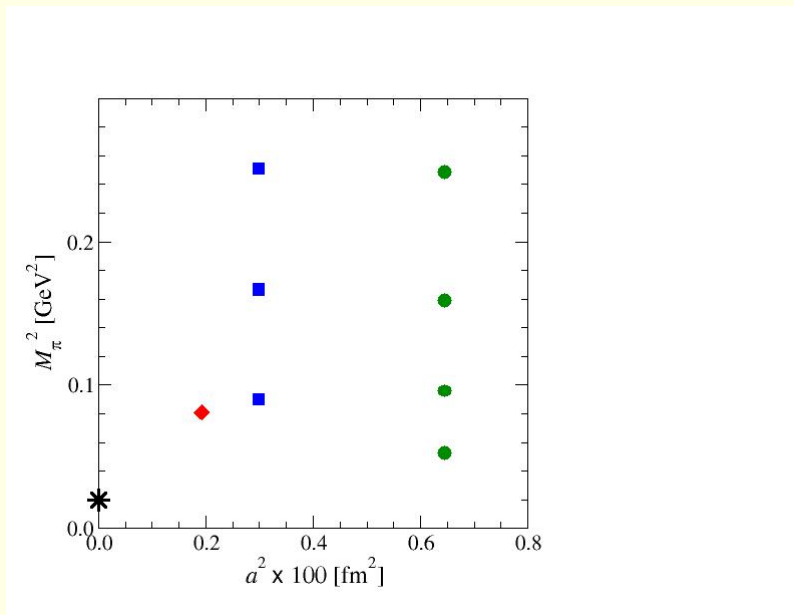
**Future work:** Improve statistics and finish systematic error analysis.

Extension to  $D \rightarrow K$ .



## 4. Work in progress: $q^2 \neq 0$

JLQCD, talk by T. Kaneko at Lattice 2016  $N_f = 2 + 1$



- \* **domain wall** ensembles including (on-going)
- \*\* Pion masses in the range [230 – 500] MeV
- \*\* Large volumes  $m_\pi L \geq 4$
- \* Calculate  $f_+^{D\pi}$  and  $f_0^{DK}$  for several  $q^2$  values
- \* Three different lattice spacings.
- \* Calculate renormalization non-perturbatively.
- \* Still adding ensembles to the analysis.

Use **z-expansion** to extrapolate results to  $q^2 = 0$  **preliminary**

$$f_{+,0}^{D\pi}(t) = \frac{1}{1 - t/M_{\text{pole}}^2} \sum_k a_k z^k$$

- \*  $f_+(t)$ : Include measured  $M_{D(s)^*}$ ,  $k = 1$  fit ( $k = 2$  for systematic error)
- \*  $f_0(t)$ :  $k = 1$  fit with no pole ( $k = 1$  with pole for systematic error)

## 4. Work in progress: $q^2 \neq 0$

JLQCD, talk by T. Kaneko at Lattice 2016  $N_f = 2 + 1$

Test a simple linear extrapolation (similar for  $f_{+,0}^{DK}$ )

$$f_{+,0}^{D\pi}(0, m_\pi^2, m_{\eta_s}^2, a) = c_0 + c_\pi m_\pi^2 + c_{\eta_s} m_{\eta_s}^2 + c_a a^2$$

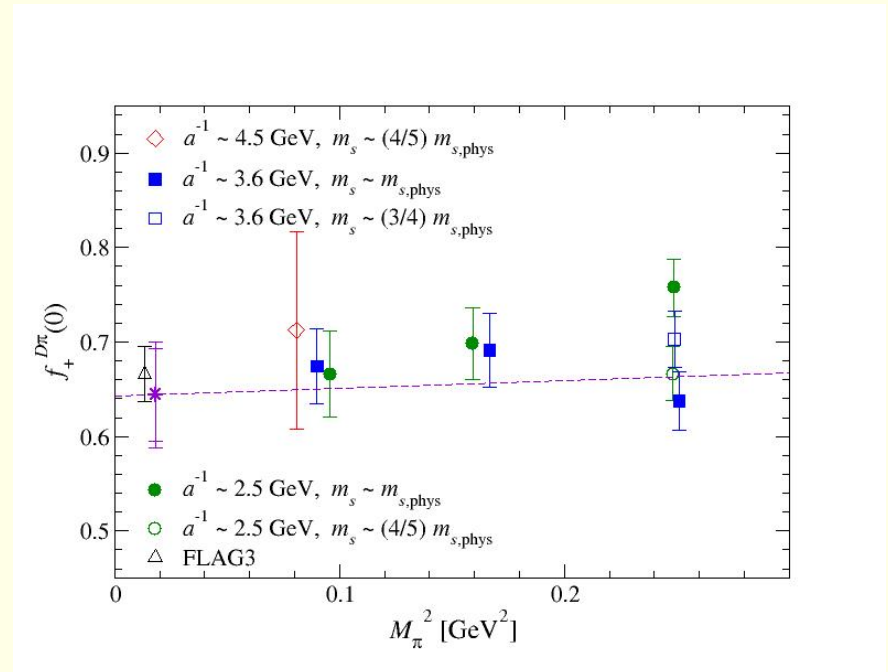
**Preliminary estimate**

$$f_+^{D\pi}(0) = 0.644(49)_{stat} (36)_{q^2 \rightarrow 0} (27)_{cont+chiral}$$

$$f_+^{DK}(0) = 0.701(46)_{stat} (12)_{q^2 \rightarrow 0} (33)_{cont+chiral}$$

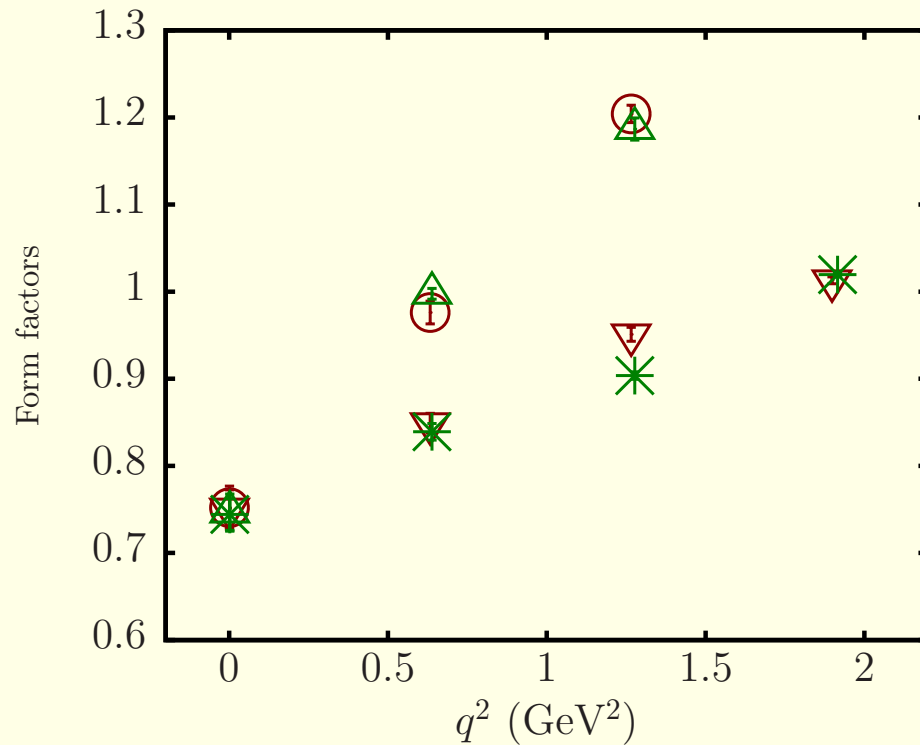
**Dominated by statistical error**

**Future work:** Add more ensembles to the analysis, more sophisticated cont+chiral extrapolation  $\rightarrow$  significant reduction of errors in the near future. Extension to  $B$  physics.



## 4. Work in progress: $q^2 \neq 0$

HPQCD, B. Chakraborty  $N_f = 2 + 1 + 1$

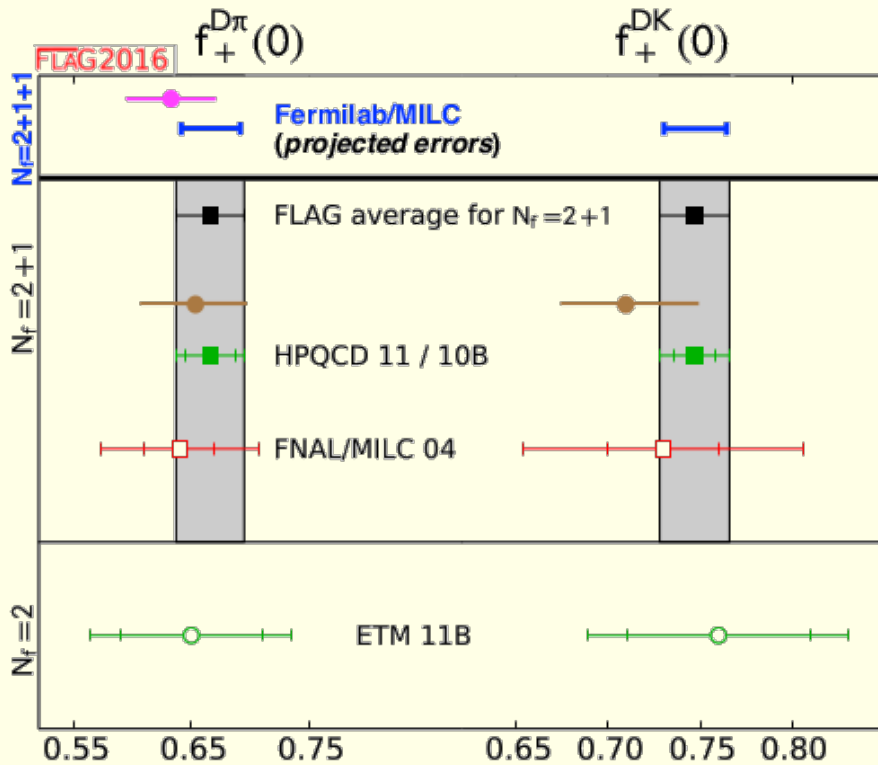


- \* red circles:  $f_+$  at  $a \sim 0.09$  fm
- \* red triangles:  $f_0$  at  $a \sim 0.09$  fm
- \* green triangles:  $f_+$  at  $a \sim 0.12$  fm
- \* green stars:  $f_0$  at  $a \sim 0.12$  fm

Physical light quark masses

HISQ lattice action: Very small discretization errors.

# 4. Work in progress: Summary



## Preliminary results

blue: FNAL/MILC (assuming central values do not change from average)

magenta: ETM

brown: JLQCD

Plot by A. El-Khadra

$f_+^{D\pi}(0)/f_+^{DK}(0)$ : **BESIII**, PRD92, 072012(2015)  $0.8649 \pm 0.0112 \pm 0.0073$

**LCSR**, P. Ball, PLB641,50(2006)  $0.84 \pm 0.04$

**FLAG** averages:  $0.89 \pm 0.04$

Further reduction of error if new calculations take correlations into account in the ratio

## 5. Correlations with leptonic decays

Cancel **CKM** matrix elements building ratios of semileptonic and leptonic decay widths

\* Experimental averages from **PDG16** (leptonic) and **F. Porter 1604.04940** (semil.)

$$\left[ \frac{f_+^{D\pi}(0)}{f_{D^+}} \right]_{exp} = (3.12 \pm 0.08) \text{ GeV}^{-1} \quad \left[ \frac{f_+^{DK}(0)}{f_{D_s}} \right]_{exp} = (2.87 \pm 0.05) \text{ GeV}^{-1}$$

\* Theoretical (lattice) averages from **PDG16** (leptonic including  $N_f = 2 + 1$  and  $N_f = 2 + 1 + 1$  calculations) and **FLAG 1607.00299** (semileptonic)

$$\left[ \frac{f_+^{D\pi}(0)}{f_{D^+}} \right]_{lat} = (3.14 \pm 0.14) \text{ GeV}^{-1} \quad \left[ \frac{f_+^{DK}(0)}{f_{D_s}} \right]_{lat} = (3.00 \pm 0.08) \text{ GeV}^{-1}$$

Good agreement experiment-theory for  $|V_{cd}|$ -ratios, not so good for  $|V_{cs}|$ -ratios  $\rightarrow$  slight tension between leptonic and semileptonic determinations of  $|V_{cs}|$

## 6. Semileptonic $D$ decays: beyond gold-plated quantities

# Alternative determination of  $|V_{cs}|$ :  $D_s \rightarrow \phi l \nu$  **HPQCD**, 1311.6669

More challenging: five form factors (vector meson), unstable meson ...

\* Treat  $\phi$  as stable and estimate the error.

\*  $q^2$  and angular distributions agree with **BaBar** data.

$$|V_{cs}| = 1.017(44)_{lat}(35)_{exp}(30)_{K\bar{K}}$$

\* Expected reduction of exper. errors at **BESIII**  $\rightarrow$  need improvement of theor. calculation (lattice error dominated by statistical error)

# Exploratory  $N_f = 2 + 1$  calculation of  $D \rightarrow \eta^{(\prime)} l \nu$  **G. Bali et al**, 1406.5449

\* Calculate  $\eta - \eta'$  mixing angles and disconnected contributions

## 7. Conclusions and outlook

Relativistic description of charm → important reduction of lattice QCD errors in decay constants and semileptonic form factors ...

$$\text{Error } f_+^{DK(\pi)} \sim 2.5 - 4.3\%$$

... still theory errors are dominant in  $|V_{cd(cs)}|$  extractions from semileptonic decays.

**Goal:**  $\sim 1\%$  error in the form factors.

# Not new results since 2013 but several on-going calculations (no statistically correlated with current calculations) will further reduce the error in the following 1-2 years

\*  $N_f = 2 + 1 + 1$  FNAL/MILC:  $f_+^{D\pi(K)}(q^2 = 0)$

\*  $N_f = 2 + 1 + 1$  ETM and HPQCD,

and  $N_f = 2 + 1$  JLQCD: shape of  $f_{+(0)}^{D\pi(K)}(q^2)$  over entire recoil range

Extensions to FCNC form factors ( $f_T$ ) straightforward

## 7. Conclusions and outlook

What do we need to achieve the targetted errors?

- \* Form factors over the entire recoil range.
- \* Physical quark masses, especially for  $D \rightarrow \pi$  quantities.
- \* Small lattice spacings and statistical errors (straightforward, but expensive)
- \* And will eventually need to include subdominant effects:
  - \*\* Include **charm** in the sea (already started).
  - \*\* **EM effects**  $\rightarrow$  Eventually will do QCD+QED simulations.
  - \*\* **Strong isospin breaking effects**: leading order corrections included via tuning light valence quarks (effects of degenerate sea are NNLO in **CHPT**).
- # Shed light over current **tensions** in the **unitarity of the second row**, and between leptonic and semileptonic  $|V_{cs}|$  determinations. See **A. Soffer** talk
- \* Interesting to improve theory error in  $D_s \rightarrow \phi l \nu$  (upcoming improvement of experimental error by **BESIII**)
- # Validate lattice techniques  $\rightarrow$  extend to  $B$  physics



