



# EXISTENCE AND FORMATION OF SMALL AMPLITUDE ELECTROSTATIC DOUBLE-LAYER STRUCTURE IN SUPERTHERMAL PLASMA

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# Introduction and Motivation

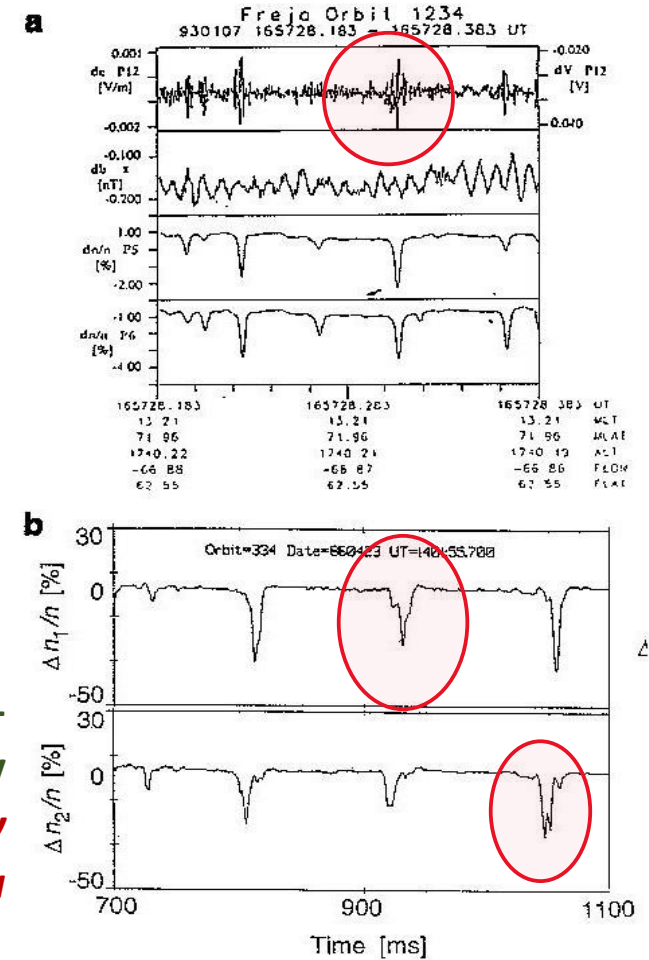
- *Ion acoustic solitary structures observed by the FREJA satellite. FREJA and VIKING satellite both observe the electric field in space plasma which suggest that they are mostly likely electrostatic in nature [Dovner et al., 1994].*

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## Electrostatic solitary structures in non-thermal plasmas

R.A. Cairns<sup>1</sup>, A.A. Mamun<sup>1</sup>, R. Bingham<sup>2</sup>, R. Boström<sup>3</sup>, R.O. Dendy<sup>4</sup>,  
C.M.C. Nairn<sup>5</sup>, P.K. Shukla<sup>6</sup>

- He show that in the presence of a distribution of electron which is non-thermal, with an excess of energetic particles the nature of ion sound solitary structures changes and *that is possible to obtain the solution with density depletions and dimensions roughly agreement with those observed Freja and Viking satellite.*

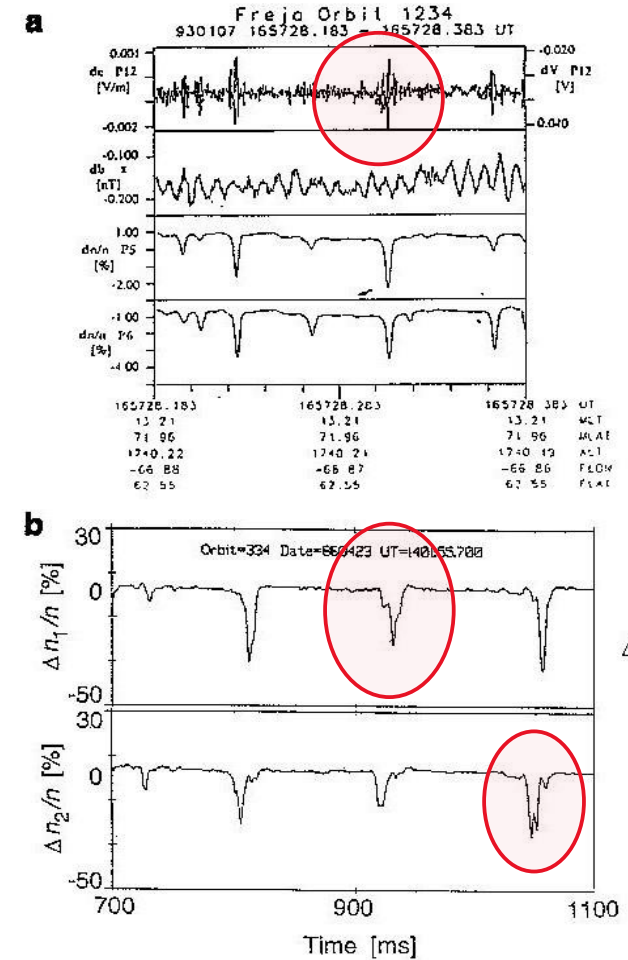


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# Introduction and Motivation

- Cairns et al. (1995) describe the distribution for non-thermal electrons
- FREJA and VIKING satellite both observe the electric field in space plasma which suggest that they are mostly likely electrostatic in nature
- FREJA satellite observe density known as lower-hybrid cavitons
- VIKING satellite also observed similar structure without associated lower hybrid waves.
- Maxwell distribution cannot give the solution to the upper and lower cavities occur in the electrostatic waves



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# Model Equations

$$f(v) = \frac{n_{e0}}{(1+3\alpha)(\pi v_e^2)} \left( 1 + \alpha \frac{v^4}{v_e^4} \right) \exp\left(-\frac{v^2}{v_e^2}\right)$$



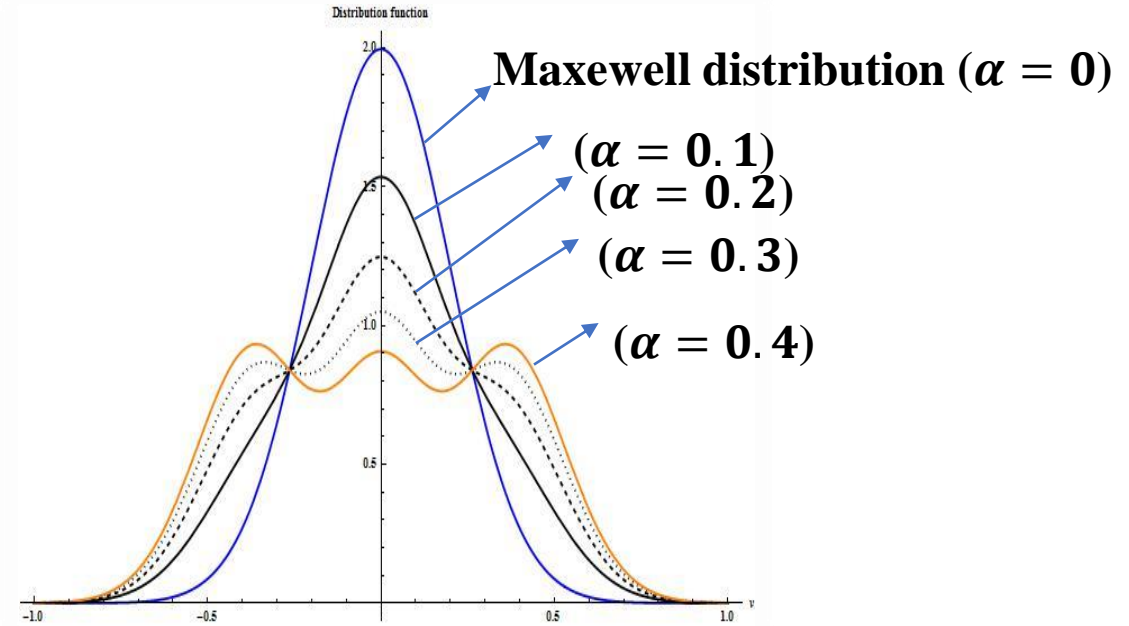
$$n_e = n_{e0} \left\{ 1 - \frac{4\alpha}{1+3\alpha} \left( \frac{e\phi}{T_e} \right) + \frac{4\alpha}{1+3\alpha} \left( \frac{e\phi}{T_e} \right)^2 \right\} \exp\left(\frac{e\phi}{T_e}\right)$$

$$\frac{\partial n_i}{\partial t} + \frac{\partial (n_i v_i)}{\partial x} = 0,$$

$$\frac{\partial v_i}{\partial t} + v_i \frac{\partial v_i}{\partial x} = -\frac{\partial \phi}{\partial x} - \frac{1}{m_i n_i} \frac{\partial p_i}{\partial x}$$

$$\frac{\partial^2 \phi}{\partial x^2} = 4\pi e (n_e - n_i + Z_d n_{d-0} - Z_d n_{d+0})$$

$$\frac{n_{e0}}{n_{i0}} = 1 - \delta z, \quad \longrightarrow \quad \delta = n_{d-0}/n_{i0} + n_{d+0}/n_{i0}$$



## Normalization Scheme

Time ( $t$ )  $\rightarrow \omega_{pi}^{-1} = (\epsilon_0 m_i / n_{i0} e^2)$ ,  
 Space parameter ( $x$ )  $\rightarrow \lambda_i = (\epsilon_0 T_e / n_{i0} e^2)$ ,  
 $n_i, n_e \rightarrow n_0, \quad v_i \rightarrow c_s = \sqrt{T_e / m_i}, \quad \phi \rightarrow T_e / e$



# Introduction and Motivation

$$\Omega^2 = \left( \frac{1}{k^2 + (1 - \delta z) \left(1 - \frac{4\alpha}{1+3\alpha}\right)} + \sigma \right) k^2$$

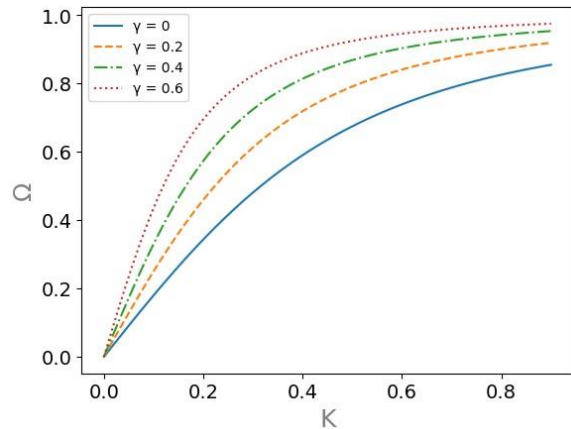


FIG. 1: The effect of nonthermal electrons on the dispersion relation configured at  $\delta z = 0.59$  and  $\sigma = 0.01$ .

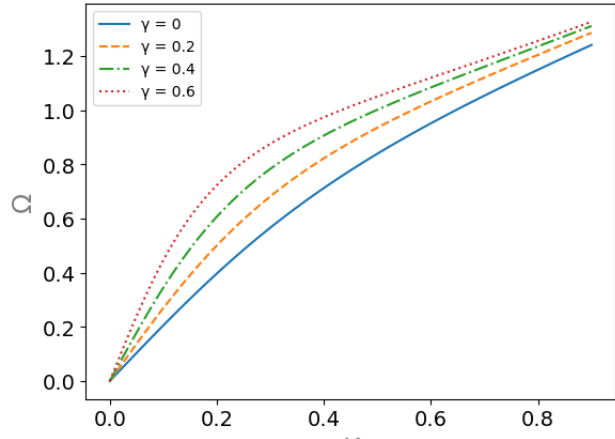


FIG. 1: The effect of nonthermal electrons on the dispersion relation configured at  $\delta z = 0.59$  and  $\sigma = 1$ .

## Small Amplitude Nonlinear Double Structures

$$X = \epsilon(x - \lambda t), \quad \tau = \epsilon^3 t,$$

$$\begin{aligned} n_i &= 1 + \epsilon n_i^{(1)} + \epsilon^2 n_i^{(2)} + \dots, \\ v_i &= \epsilon v_i^{(1)} + \epsilon^2 v_i^{(2)} + \dots, \\ \Phi &= \epsilon \Phi^{(1)} + \epsilon^2 \Phi^{(2)} + \dots \end{aligned}$$



$$\lambda = \sqrt{\frac{1}{(1 - \delta z) \left(1 - \frac{4\alpha}{1+3\alpha}\right)} + \sigma}.$$

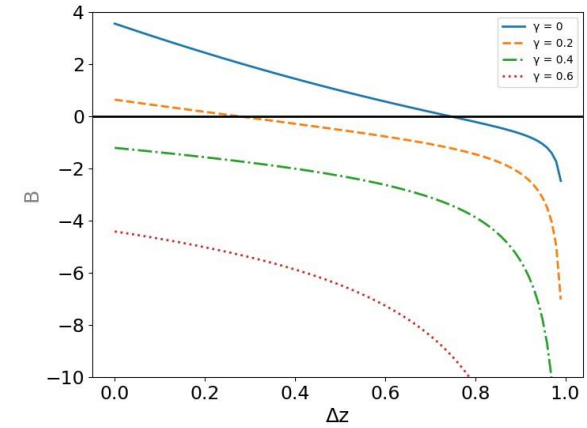
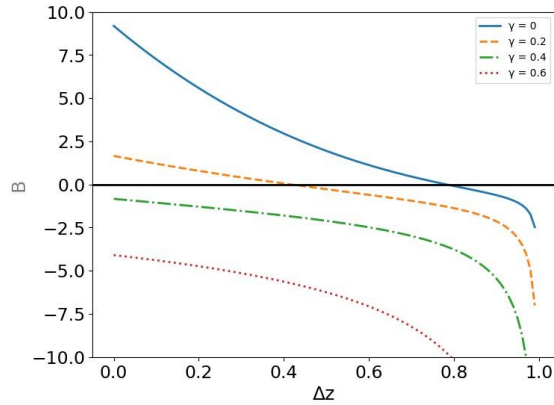
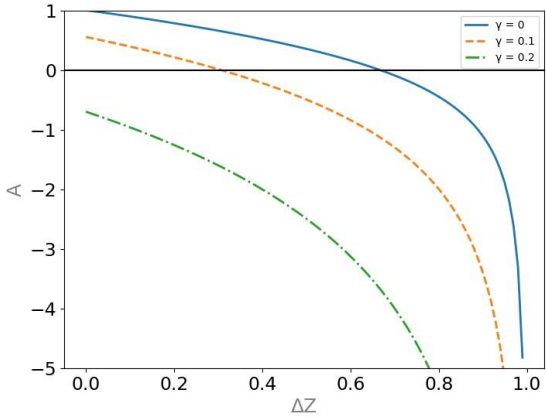
$$\frac{\partial}{\partial \tau} \Phi^{(1)} + A \Phi^{(1)} \frac{\partial}{\partial X} \Phi^{(1)} + B \Phi^{(1)^2} \frac{\partial}{\partial X} \Phi^{(1)} + C \frac{\partial^3}{\partial X^3} \Phi^{(1)} = 0.$$



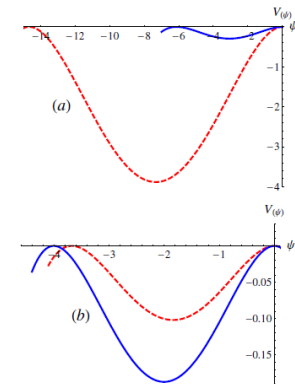
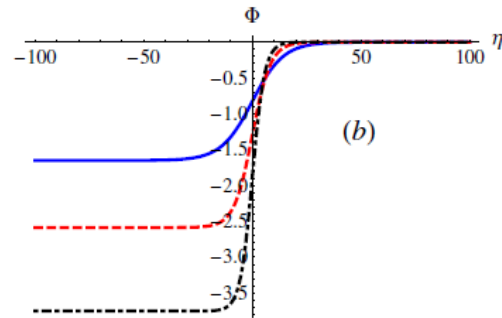
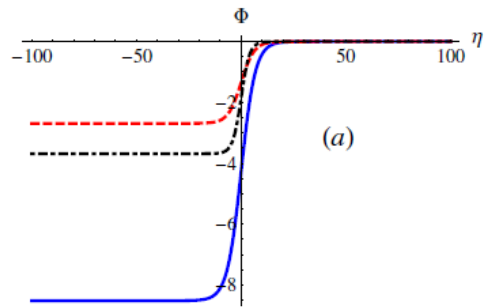
$$\psi = \Phi^{(1)} = \left(\frac{\psi}{2}\right) \left[1 - \tanh\left(\frac{2\eta}{W}\right)\right]$$



# Nonlinear Shock waves/Double layers



Nonlinear Coefficients “A” and “B” against Dust charge concentration with various non-thermality and ionic temperature configurations



Shock wave/double layer profiles for various Dust charge concentration, non-thermality and the corresponding Sagdeev Potential for the formation regime.





# Conclusions

- The present model support both positive and negative polarity shock structures (double layer) and switches sign with changing the dust charge concentration.
- Strong non-thermality favors only rarefactive structures irrespective of the ion temperature compared to the electron temperature.
- The existence regimes for the formation of ion-acoustic double layer in a dusty plasma is traced out
- It is found that increasing the nonthermal electrons in the system, the width of the double layer increases, furthermore, the shock structure forms with small dust charge concentrations.
- Our present work may be useful for understanding the electrostatic feature of DIAWs in both experiments and cosmic dusty plasmas such as discussed Ref. [4]
- The double layers (Shock structures) are also observed in dusty plasma in Q-machine experiments for  $\sigma = 1$  and dust charge concentrations  $\delta z \geq 0.75$  as mentioned in Ref. [5]





## References

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