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# $e/\gamma$ reconstruction & identification in CMS

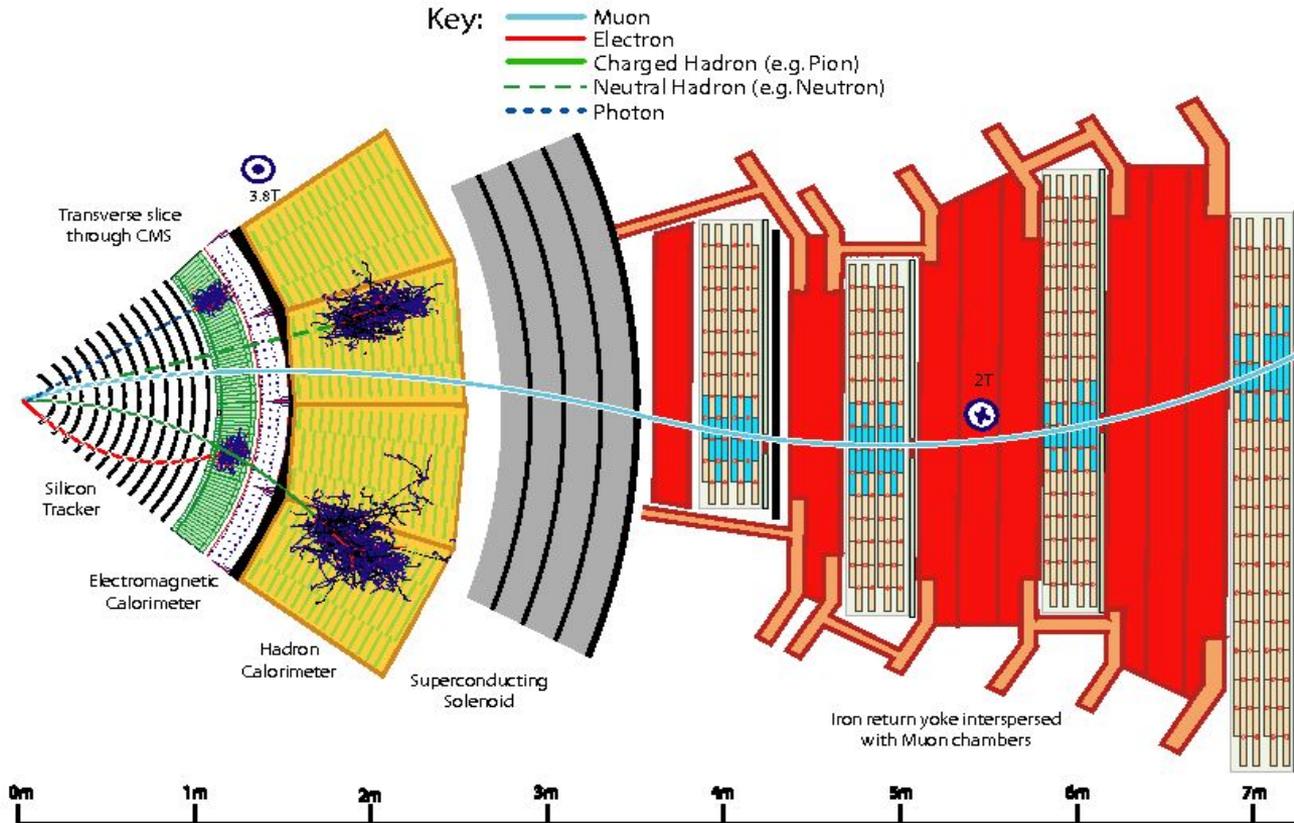
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# Introduction

- $e/\gamma$  are reconstructed with high purity & efficiency in CMS.
- Distinctive signal in electromagnetic calorimeter (ECAL) as an isolated energy deposit associated with a trace in the 'Si' tracker in case of electrons.
- The main topics that will be covered in this talk:
  - Short description of the CMS detector
  - Offline  $e/\gamma$  reconstruction
  - Online  $e/\gamma$  reconstruction
  - Difference between offline & online reconstruction
  - Electron & photon selection
  - Photon identification
  - Electron identification

# Cross-sectional view of the CMS detector



# Short description of the CMS detector

- **The magnet:** central feature of the CMS detector is a large superconducting solenoid magnet. It delivers an axial & uniform magnetic field of 3.8T which helps to separate the calorimeter energy deposits of charged & neutral particles.
- **The silicon inner tracker:** provides a pure & efficient charged-particle trajectory reconstruction in jets with  $p_T$  upto around 1 TeV.
- **Highly-segmented ECAL:** homogeneous calorimeter made of lead tungstate ( $\text{PbWO}_4$ ) crystals, sufficient to contain more than 98% of the energy of electrons & photons upto 1 TeV. Two thirds of the hadrons also start showering in the ECAL before entering the HCAL. A much finer-grained detector, called **preshower**, is installed in front of each endcap disk. **The aim of the preshower**  $\Rightarrow$  (1) to discriminate between prompt photons & photons coming from  $\pi^0$  decay (2) to indicate the presence of electron or photon by requiring an associated signal in the preshower.
- **HCAL:** sampling calorimeter consisting of brass absorber & plastic scintillator. Charged & neutral hadrons may initiate a hadronic shower in ECAL, which is fully absorbed in the HCAL.
- **The muon detectors:** muons produce hits in additional tracking layers located outside the calorimeters.

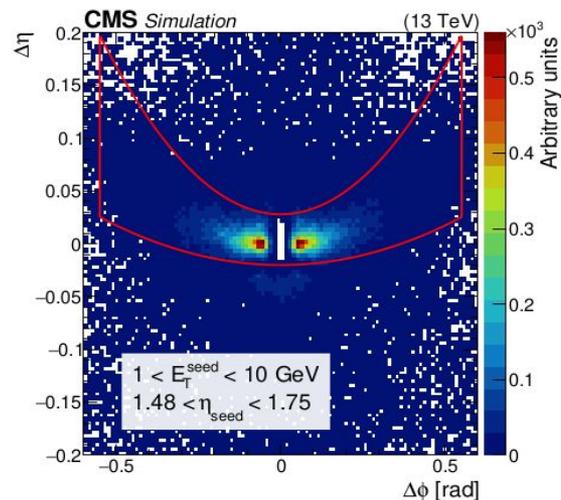
# Offline $e/\gamma$ reconstruction

# Overview

- Electrons & photons deposit almost all of their energy in ECAL, whereas hadrons deposit in HCAL.
- In addition, **electrons produce hits in the tracker.**
- **Electrons** ⇨ bremsstrahlung photons, **photons** ⇨ electron positron pairs; thus when electron or photon reaches the ECAL, it may consist of a **shower of particles.**
- A dedicated algorithm to combine clusters of particles into a single object to recover the energy of primary electron or photon ⇨ **Gaussian Sum Filter (GSF).**
- The energy reconstruction algorithm starts with a formation of clusters (> 80 MeV in EB, > 300 MeV in EE). Seed consist of most energy deposits (> 1 GeV).
- Some neighbouring clusters around the SC to include pair production & brem. loss ⇨ **superclustering**
- All objects with an associated GSF track ⇨ electrons; without a GSF track ⇨ photons

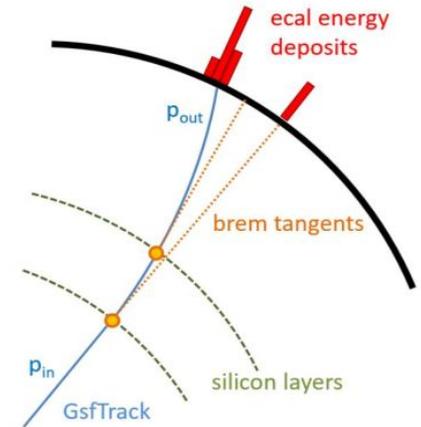
# Superclustering in ECAL

- Because of showering the original object may consist of several  $e/\gamma$  produced from pair production/brem.
- The multiple ECAL clusters need to be combined into a single SC  $\Rightarrow$  **superclustering**
- There are mainly two algorithms  $\Rightarrow$ 
  - “Mustache” algorithm
  - “Refined” algorithm
- “Mustache” algorithm  $\Rightarrow$  useful for low-energy deposits. It uses information only from ECAL & preshower. The size of the mustache region depends on  $E_T$
- “Refined” algorithm  $\Rightarrow$  It uses the tracking information. At each tracker layer, the trajectory of the GSF track is extrapolated to form a “brem. tangent”, which can be linked to a compatible ECAL cluster. Clusters linked to the “brem. tangent” are then added to the refined SC.



# Electron track reconstruction

- Electrons use the GSF tracking algorithm to include radiative losses from brem.
- The GSF track fitting algorithm is **CPU intensive**, so can't be run over all reconstructed hits in the tracker. **The electron trajectory seed can be either "ECAL-driven" or "tracker-driven"**.
- ECAL driven seeding first selects mustache SCs with  $E_{SC} > 4 \text{ GeV}$  &  $H/E_{SC} < 0.15$ . It performs better for high  $E_T$  electrons.
- Tracker driven seeding is performed only in offline reconstruction, not in HLT, as it is computationally expensive to reconstruct all tracks in an event. This approach is designed for low  $E_T$  electrons.
- The final collection of selected electron seeds is used to initiate the reconstruction of electron tracks.
- We have to use track-SC matching variables to associate a GSF track to an ECAL cluster.  $|\Delta\eta| = |\eta_{SC} - \eta_{track}| < 0.02$  &  $|\Delta\varphi| = |\varphi_{SC} - \varphi_{track}| < 0.15$



# Online $e/\gamma$ reconstruction

# Overview

- $e/\gamma$  candidates at L1 are based on ECAL trigger towers defined by arrays of 5x5 crystals.
- The central trigger tower with highest  $E_T$  ( $> 2$  GeV) is designated as the seed tower.
- To recover energy loss from brem, clusters are built from surrounding towers with  $E_T > 1$  GeV to form L1 candidates.
- No tracker information is available at L1, so electrons & photons are indistinguishable at this stage.
- The HLT electron & photon candidates are reconstructed from energy deposits in ECAL crystals grouped into clusters around L1 seed.
- Electron & photon selection at HLT relies on the identification & isolation criteria that will be discussed later.

# Difference between online & offline reconstruction

- The differences between the HLT & offline reconstruction are mainly driven by:
  - The limited CPU time available at the HLT.
  - The lack of final calibrations which are not yet computed during the data-taking period.
  - More conservative selection criteria to avoid rejecting interesting events.
- Every electron candidate reconstructed at the HLT is ECAL driven.
- Offline tracker driven reconstruction is useful only for low energy electrons, which is not easy to trigger.

$e/\gamma$  selection

# Overview

- Two different techniques are used in CMS for the identification of electrons & photons.
  - **Cut-based approach**
  - **MVA-based approach** (will not be discussed in this talk)
- Different strategies are used to identify prompt electrons & photons.
- For prompt electrons, bkg can originate from photon conversion, hadrons misidentified as electrons, secondary electrons coming from semi-leptonic decays of b-quark.
- The main bkg to prompt photons are **photons coming from  $\pi^0$  decay**.

# Isolation criteria

- One of the methods to reject electron & photon bkg is the use of isolation energy sums.
- At first a cone is defined in  $\eta$ - $\varphi$  plane; the distance with respect to the reconstructed electron or photon direction is defined by  $\Delta R$  ( $= 0.3$ ).
- The electron or photon itself is excluded from the isolation sum.

$$I_{\text{PF}} = \frac{1}{p_{\text{T}}} \left( \sum_{h^{\pm}} p_{\text{T}}^{h^{\pm}} + \sum_{\gamma} p_{\text{T}}^{\gamma} + \sum_{h^0} p_{\text{T}}^{h^0} \right)$$

# Shower shape criteria

- Three main shower shape variable that will be discussed:

- H/E

- $\sigma_{i\eta i\eta}$

- $R_9$

- H/E  $\Rightarrow$  ratio between the energy deposited in HCAL in a cone of  $\Delta R = 0.15$  around the SC & the energy of the  $e/\gamma$  candidate. HCAL noise, pileup, leakage of electrons or photons through the inter-modular gaps  $\Rightarrow$  'H'.  $H < X + Y\rho + JE$ ; X & Y represent the noise & pileup (PU) terms, J is scaling term for high energy electron or photon.

- $\sigma_{i\eta i\eta}$   $\Rightarrow$

$$\sigma_{i\eta i\eta} = \sqrt{\frac{\sum_i^{5 \times 5} w_i (\eta_i - \bar{\eta}_{5 \times 5})^2}{\sum_i^{5 \times 5} w_i}}. \quad w_i = \max(0, 4.7 + \ln(E_i/E_{5 \times 5}))$$

- $R_9$   $\Rightarrow$  defined as  $E_{3 \times 3}/E_{SC}$ . Showers of photons that convert before reaching the calorimeter have wider transverse profiles & lower values of  $R_9$  than unconverted photons.

# Photon identification

# Cut-based photon identification

- The “loose” working point (WP) has an average signal efficiency of about 90%, used when bkg is low.
- The “medium” & “tight” WP have an average efficiency of about 80% & 70% respectively, used when bkg is larger.

Variable	Barrel (tight WP)	Endcap (tight WP)
$H/E$	$<0.021$	$<0.032$
$\sigma_{i\eta i\eta}$	$<0.0099$	$<0.027$
$I_{\text{ch}}$	$<0.65 \text{ GeV}$	$<0.52 \text{ GeV}$
$I_{\text{n}}$	$<0.32 \text{ GeV} + 0.015E_{\text{T}} +$ $2.26 \times 10^{-5}E_{\text{T}}^2/\text{GeV}$	$<2.72 \text{ GeV} + 0.012E_{\text{T}} +$ $2.3 \times 10^{-5}E_{\text{T}}^2/\text{GeV}$
$I_{\gamma}$	$<2.04 \text{ GeV} + 0.0040E_{\text{T}}$	$<3.03 \text{ GeV} + 0.0037E_{\text{T}}$

# Electron identification

# Cut-based electron identification

Variable	Barrel (tight WP)	Endcaps (tight WP)
$\sigma_{i\eta i\eta}$	$<0.010$	$<0.035$
$ \Delta\eta_{\text{in}}^{\text{seed}} $	$<0.0025$	$<0.005$
$ \Delta\phi_{\text{in}} $	$<0.022 \text{ rad}$	$<0.024 \text{ rad}$
$H/E$	$<0.026 + 1.15 \text{ GeV}/E_{\text{SC}}$ $+0.032\rho/E_{\text{SC}}$	$<0.019 + 2.06 \text{ GeV}/E_{\text{SC}}$ $+0.183\rho/E_{\text{SC}}$
$I_{\text{combined}}/E_{\text{T}}$	$<0.029 + 0.51 \text{ GeV}/E_{\text{T}}$	$<0.0445 + 0.963 \text{ GeV}/E_{\text{T}}$
$ 1/E - 1/p $	$<0.16 \text{ GeV}^{-1}$	$<0.0197 \text{ GeV}^{-1}$
Number of missing hits	$\leq 1$	$\leq 1$

# Summary

- First of all, a brief overview of CMS detector has been discussed.
- Then I discussed about the offline electron & photon reconstruction.
- Next online  $e/\gamma$  reconstruction has been described briefly.
- After that I discussed how  $e/\gamma$  can be selected using isolation & shower shape criteria.
- At last only the cut-based approach of  $e/\gamma$  identification has been discussed.

