

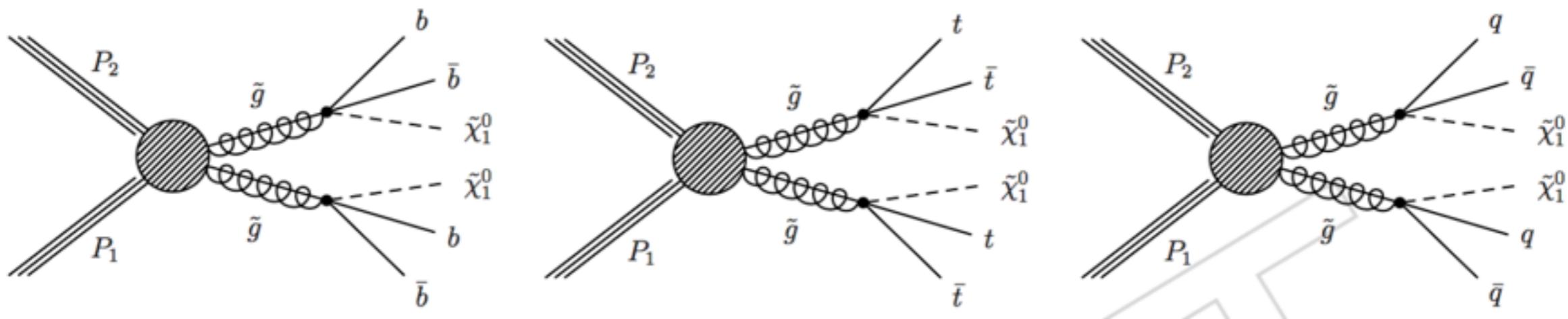


# Inclusive search for supersymmetry in hadronic final states at CMS

By: S. Norberg  
University of Puerto Rico Mayagüez  
(On behalf of CMS Collaboration)

# Motivation

- Large branching fraction to hadronic final states is typical for signals
- Final states involve events with high jet multiplicity and missing transverse momentum
- This method of search could eventually decipher the mysteries surrounding the Higgs boson and the hierarchy problem
- Two such searches covered use a different object definition to explore different gluino searches as well as squark and other direct stop production models.
- Analysis cuts & search regions are defined so that the analysis is broadly applicable to many decay modes and many bins meaning a wide range of phase space is covered

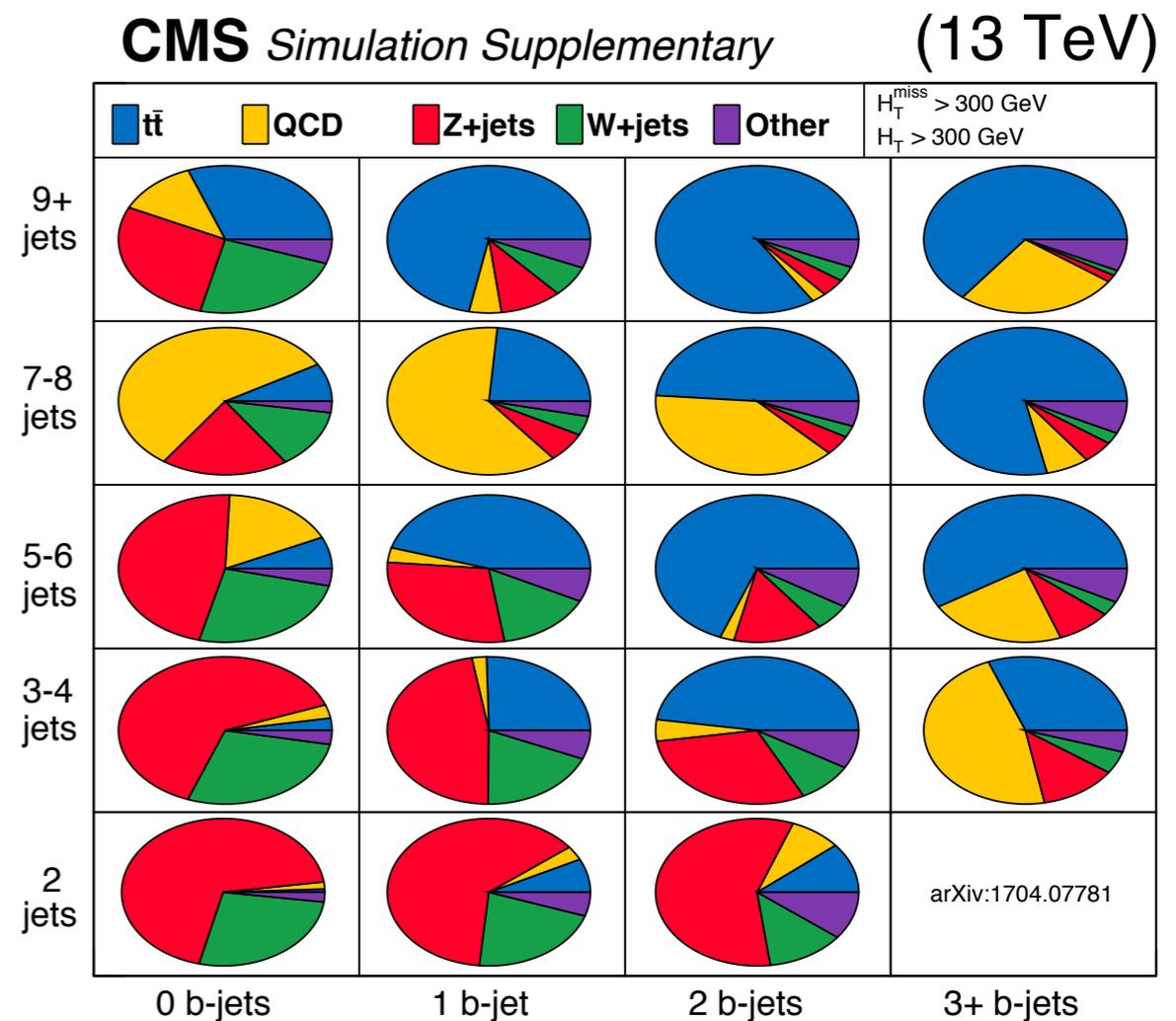


Direct Gluino production

# Common Object and Cuts

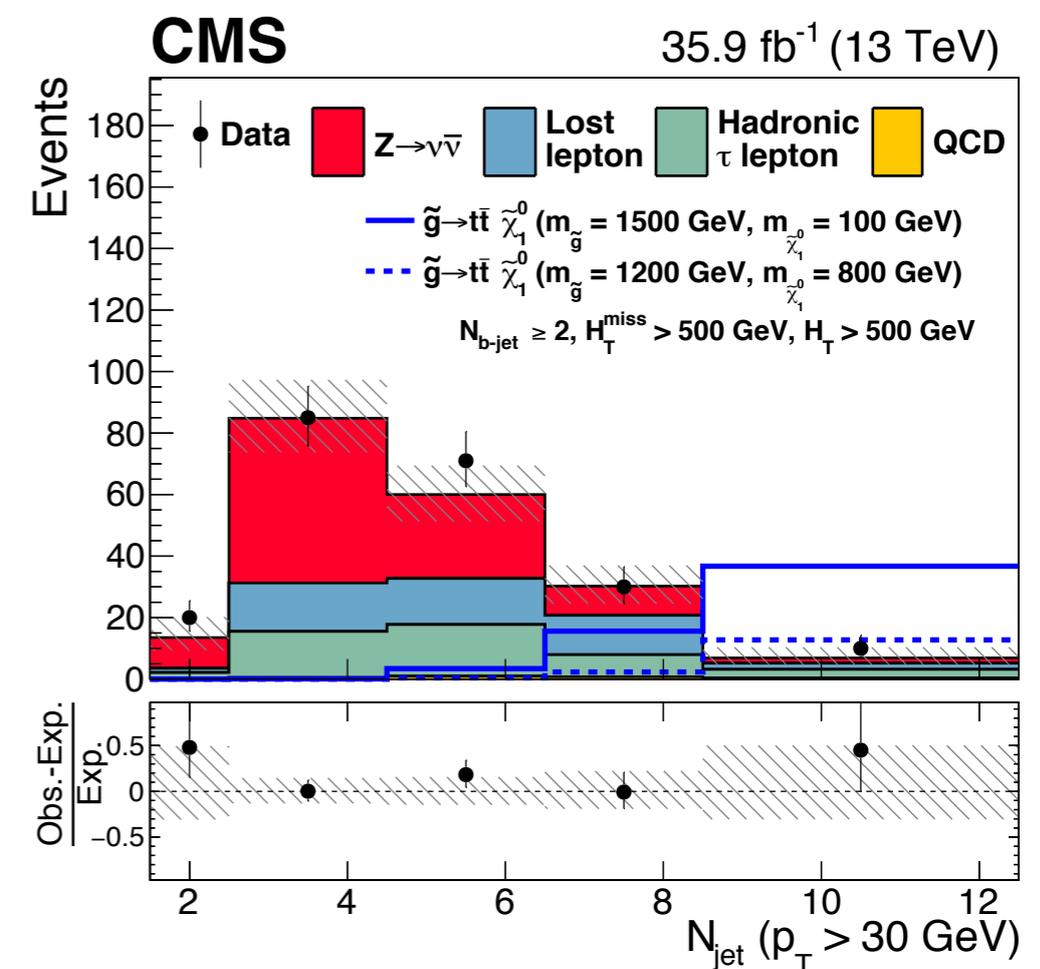
- Objects:
  - Number of b jets
  - Number of Jets
  - The scalar sum of the jet  $p_T$  which is called  $H_T$ 

$$H_T = \sum_{jets} |p_T^{\vec{p}}|$$
- To suppress backgrounds:
  - Lepton
    - lepton veto is used
  - QCD
    - Delta phi  $\Delta\phi(P^{\text{miss}}_{T,\text{jet}}) > 0.3, 0.5$



# Common Backgrounds

- $Z(\nu\nu)+\text{jets}$  ( $Z$  Invisible)
  - Irreducible background, with genuine  $P^{\text{miss}}_{\text{T}}$  from  $\nu$ 's
- Lost lepton
  - Genuine  $P^{\text{miss}}_{\text{T}}$  from leptonic  $W$  decay
  - Lepton fails acceptance of the detector, the reconstruction and identification efficiency
- QCD multijet
  - Fake  $P^{\text{miss}}_{\text{T}}$  due to either significant jet momentum mis-measurements, or sources of anomalous noise



# Search for supersymmetry in multijet events with missing transverse momentum in proton-proton collisions at 13 TeV

Uses  $M_{H_T}$

CMS-SUS-16-033

arXiv:1704.07781

Phys. Rev. D 96, 032003 (2017)

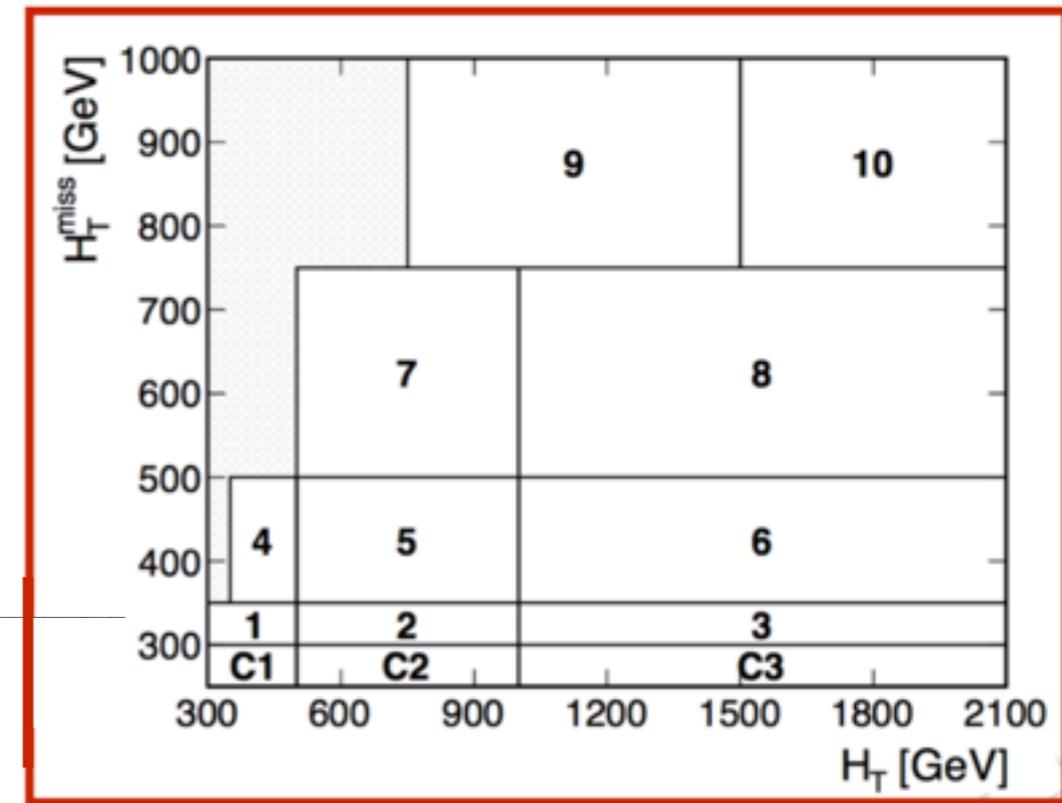
10.1103/PhysRevD.96.032003

<https://arxiv.org/abs/1704.07781>

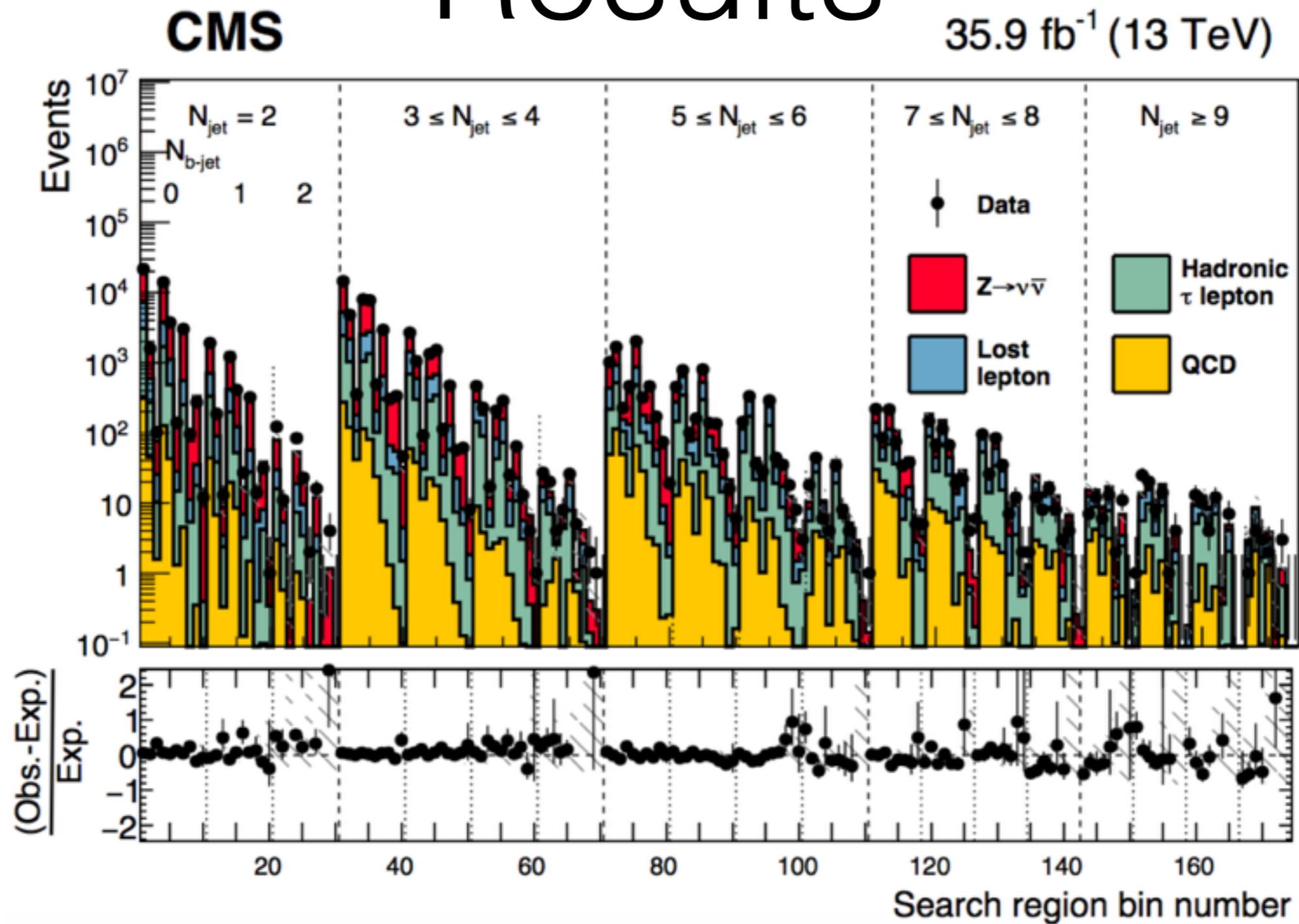
# Search Strategy

- 174 signal regions are used
- $M_{H_T}$  which is defined as the negative sum of the missing jet  $P_T$ 

$$H_T^{miss} = | - \sum_{jets} \vec{p}_T |$$
- Lost lepton
  - Relate observed counts in control region to expected counts in signal region through Transfer factors
  - single lepton control region
- Z Invisible
  - Extrapolate from  $\gamma$  control region using  $p_T$
  - $\gamma$  as a proxy for  $Z(\nu\nu)$  momentum
- QCD
  - QCD multi-jet events have no intrinsic  $P_T^{miss}$ , only instrumental  $P_T^{miss}$  due to detector response is used to estimate this background

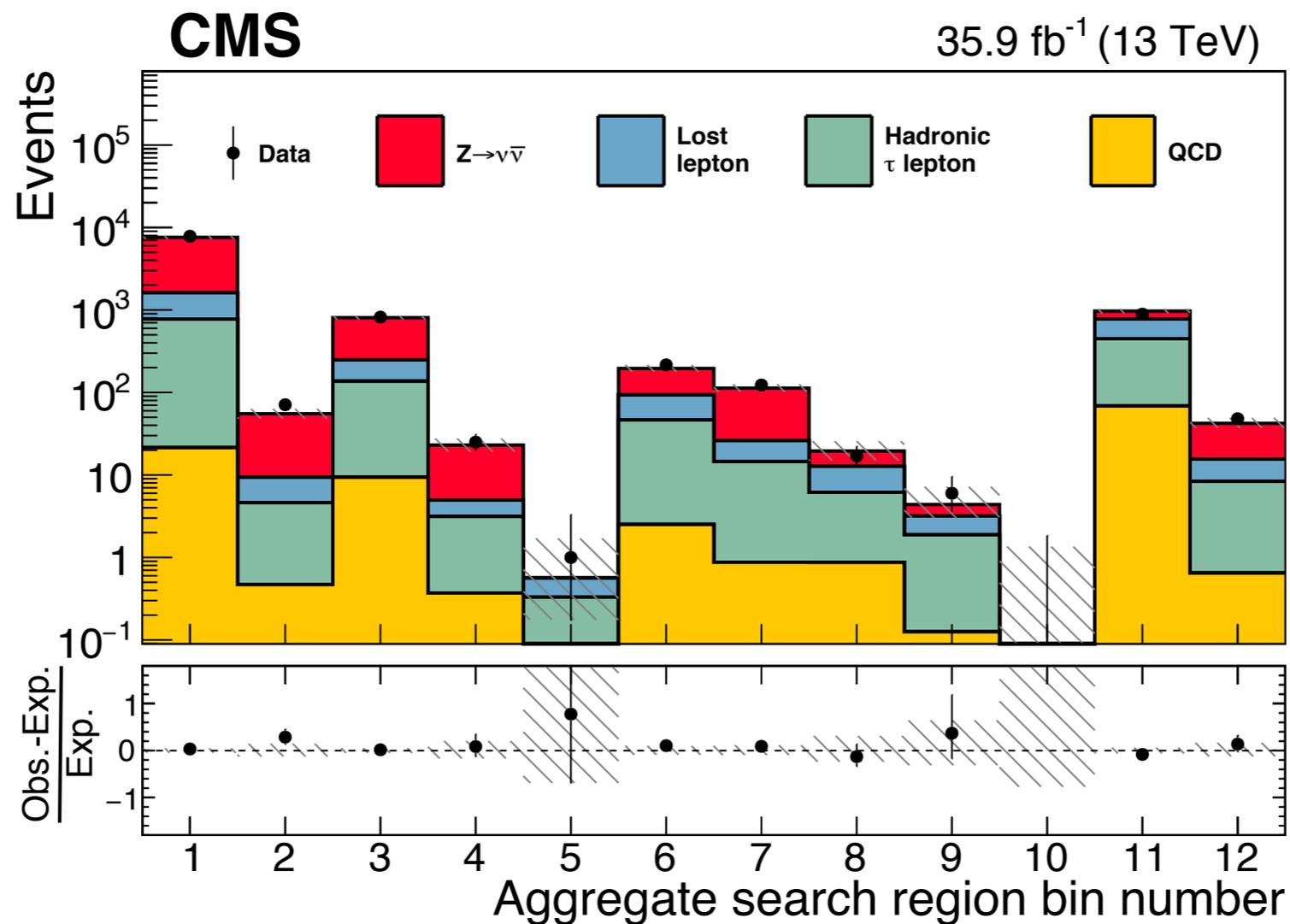


# Results



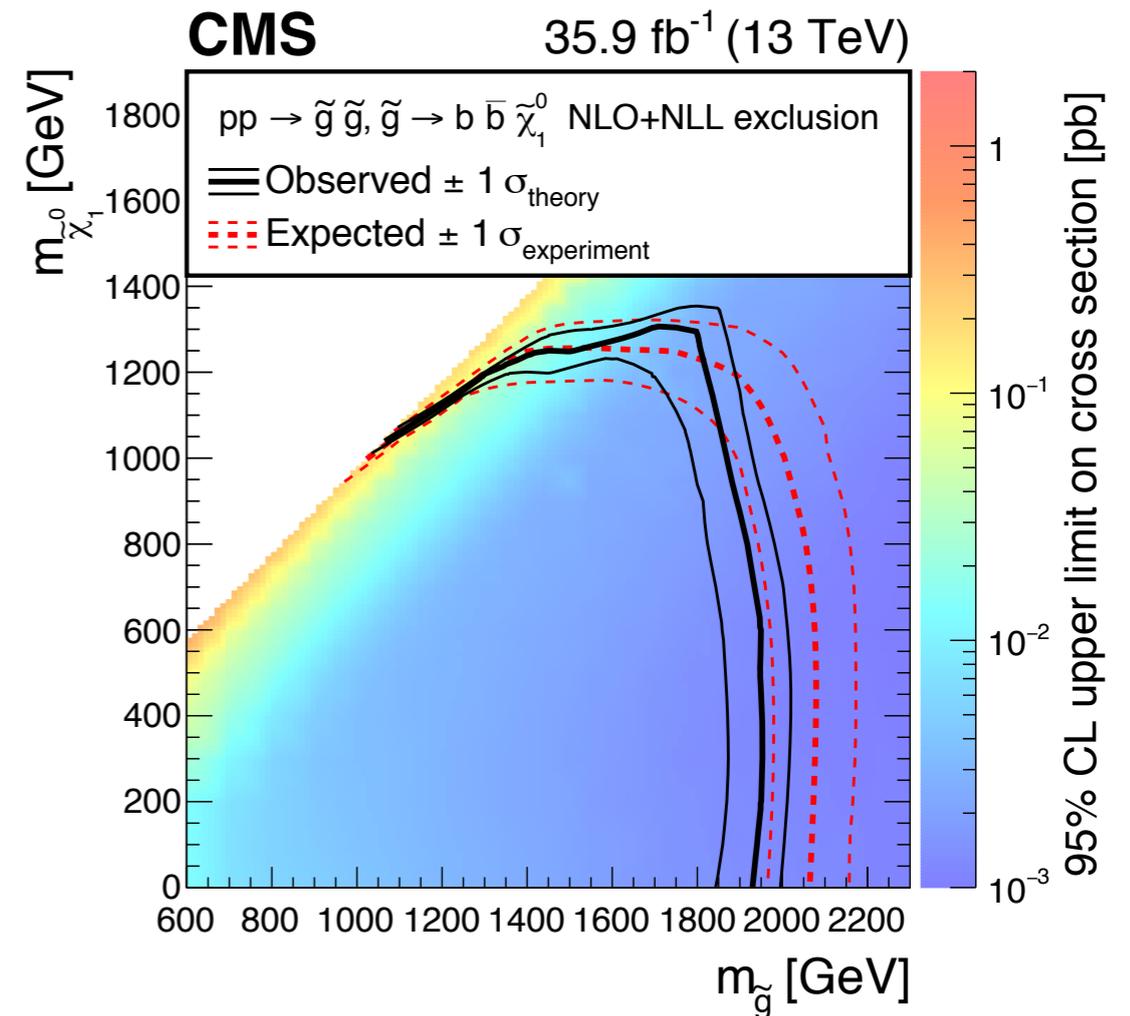
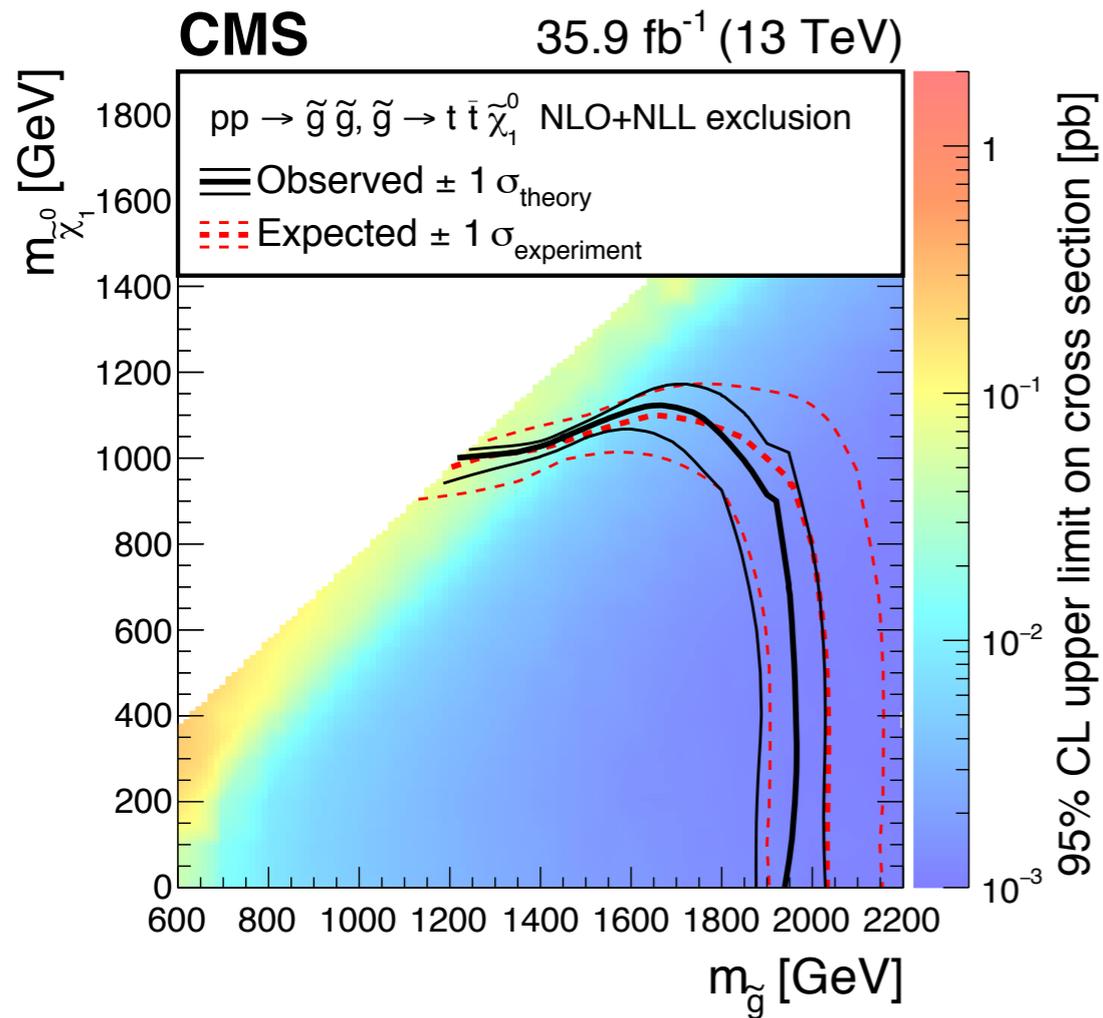
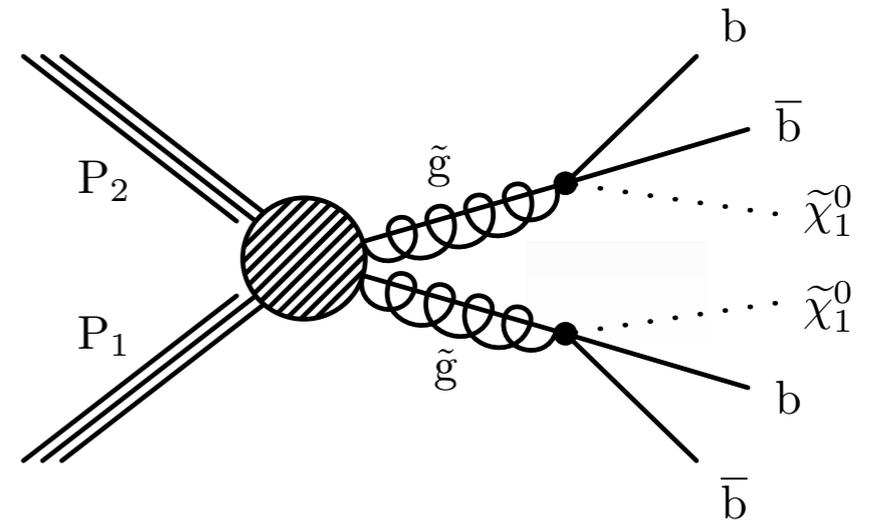
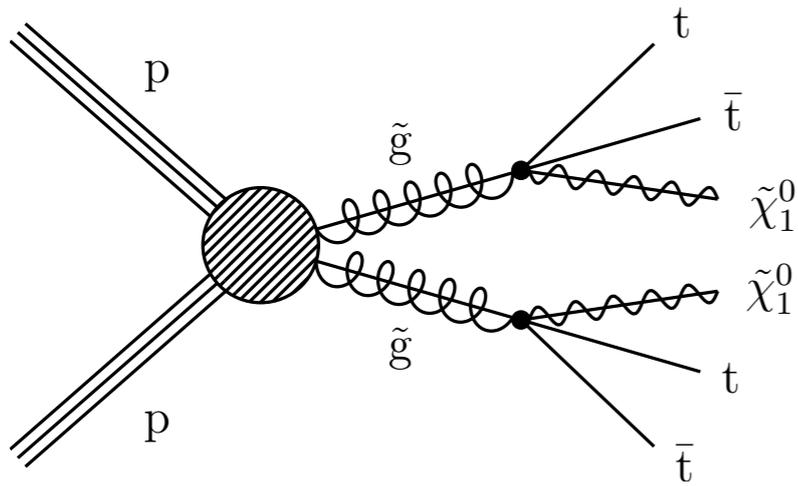
- No significant excess is seen

# Facilitating Reinterpretation



- Release small number bins that are more inclusive that give an overall view of sensitivity

# Limits





Search for new phenomena with the  $M_{T2}$  variable in the all-hadronic final state produced in proton-proton collisions at a center mass energy of 13 TeV

Uses  $M_{T2}$

CMS-SUS-16-036

arXiv:1705.04650v2

Published in the European Physical Journal C

Eur. Phys. J C 77 (2017) 710

doi:10.1140/epjc/s10052-017-5267-x

<https://arxiv.org/pdf/1705.04650.pdf>

# MT2 Variable

- $M_{T2}$  is a generalization of the transverse mass  $M_T$  for decay chains with two unobserved particles

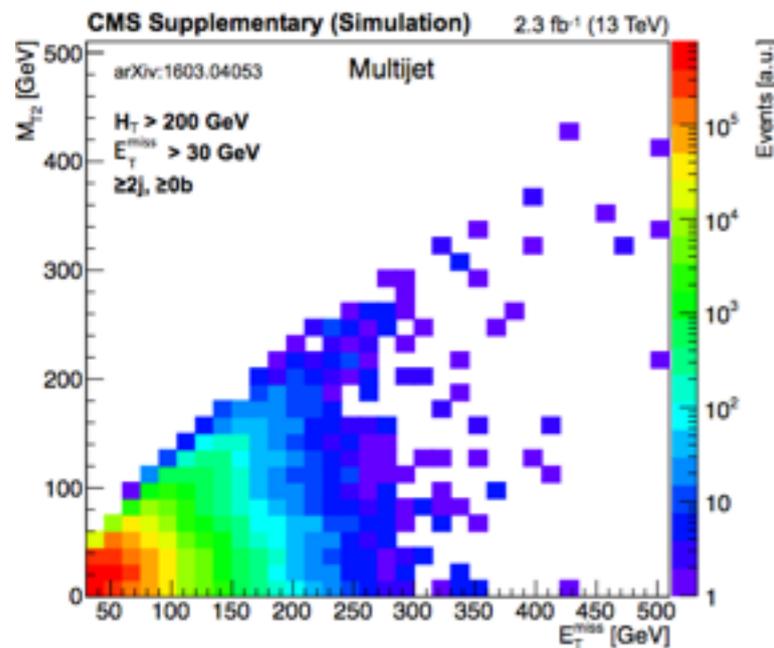
$$M_{T2}(m_X) = \min_{\vec{p}_T^{X(1)} + \vec{p}_T^{X(2)} = \vec{p}_T^{\text{miss}}} \left[ \max(M_T^{(1)}, M_T^{(2)}) \right]$$

- Where  $M_T$  is the mass between two particles where one is unobserved

$$M_T^2 = 2E_{T,1}E_{T,2}(1 - \cos\theta)$$

**Multijet:** back-to-back

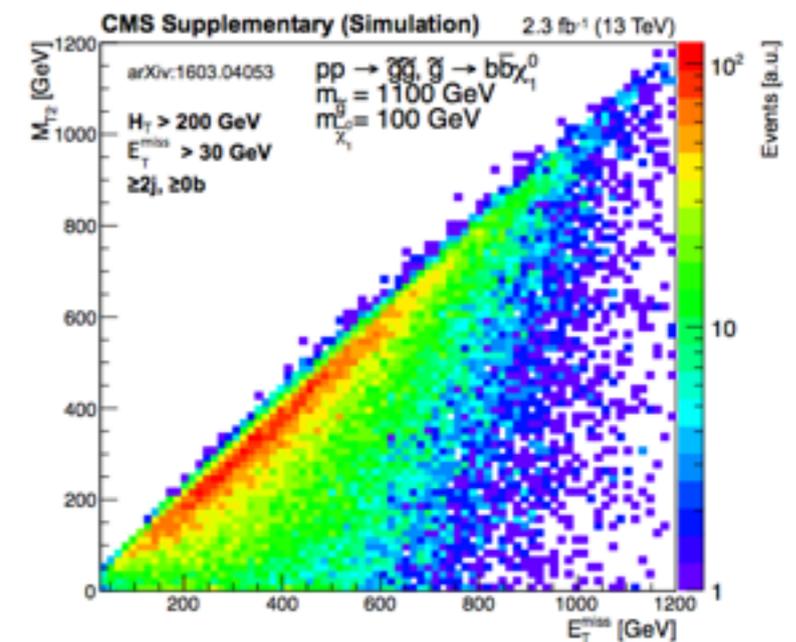
→  $M_{T2} \ll \mathbf{p}_{\text{miss}T}$



Useful because the Topology for SUSY particles is different than what is seen in the MultiJet events

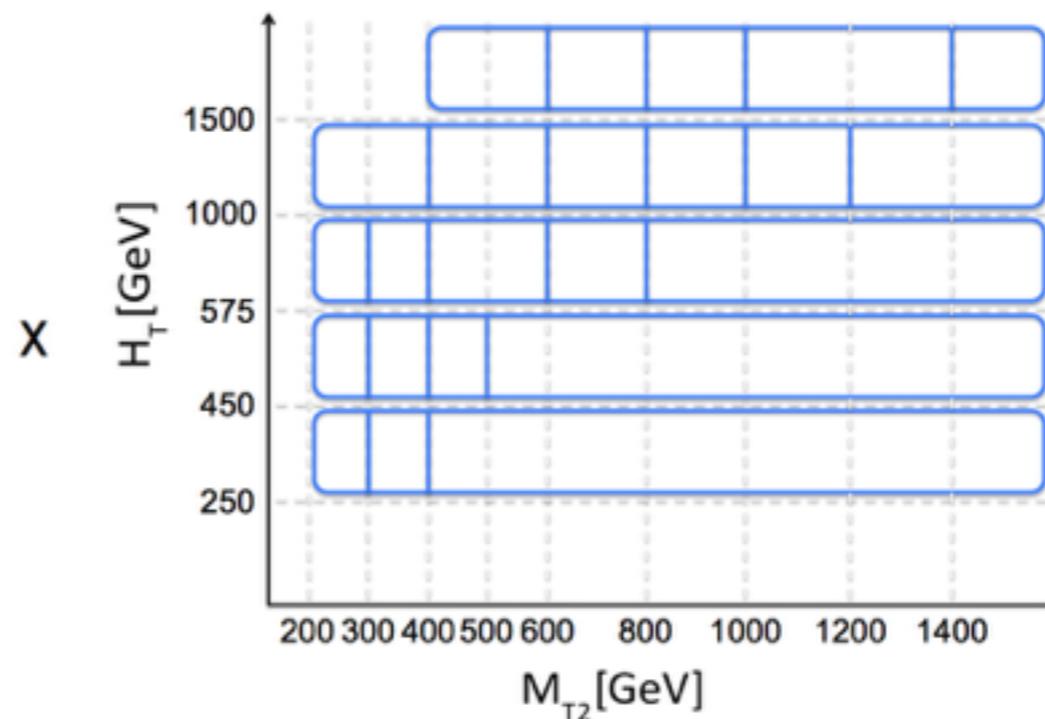
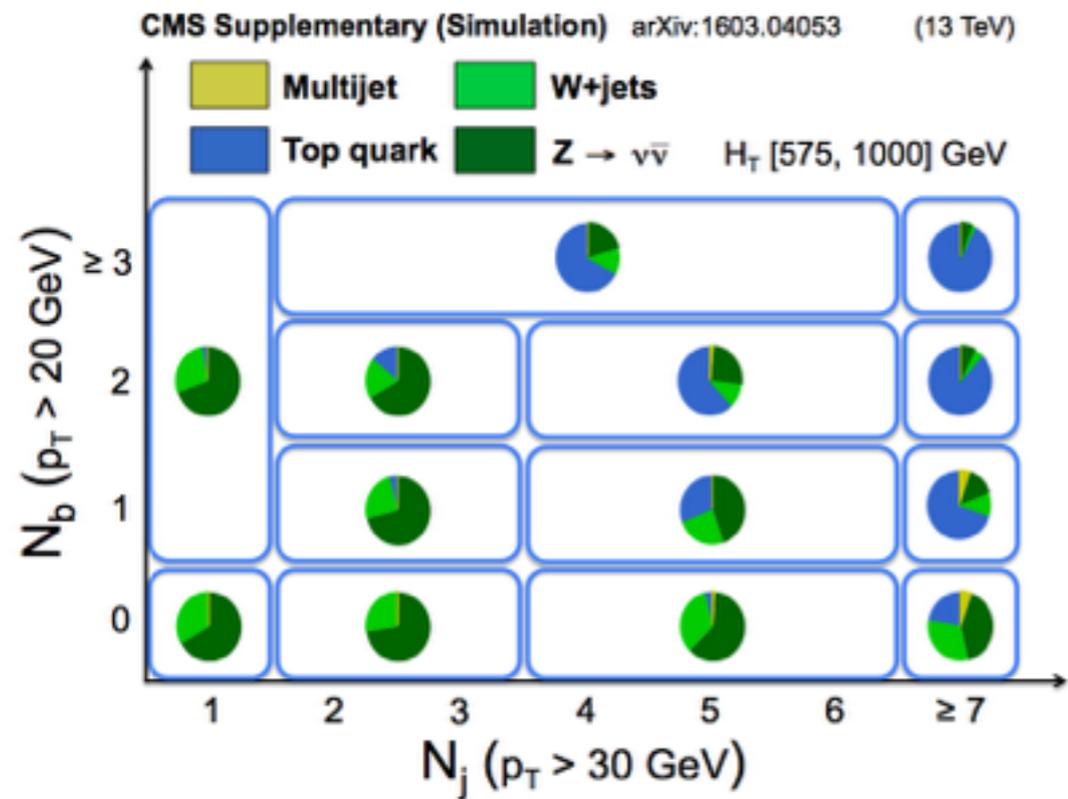
**SUSY:** symmetric topology

→  $M_{T2} \sim \mathbf{p}_{\text{miss}T}$



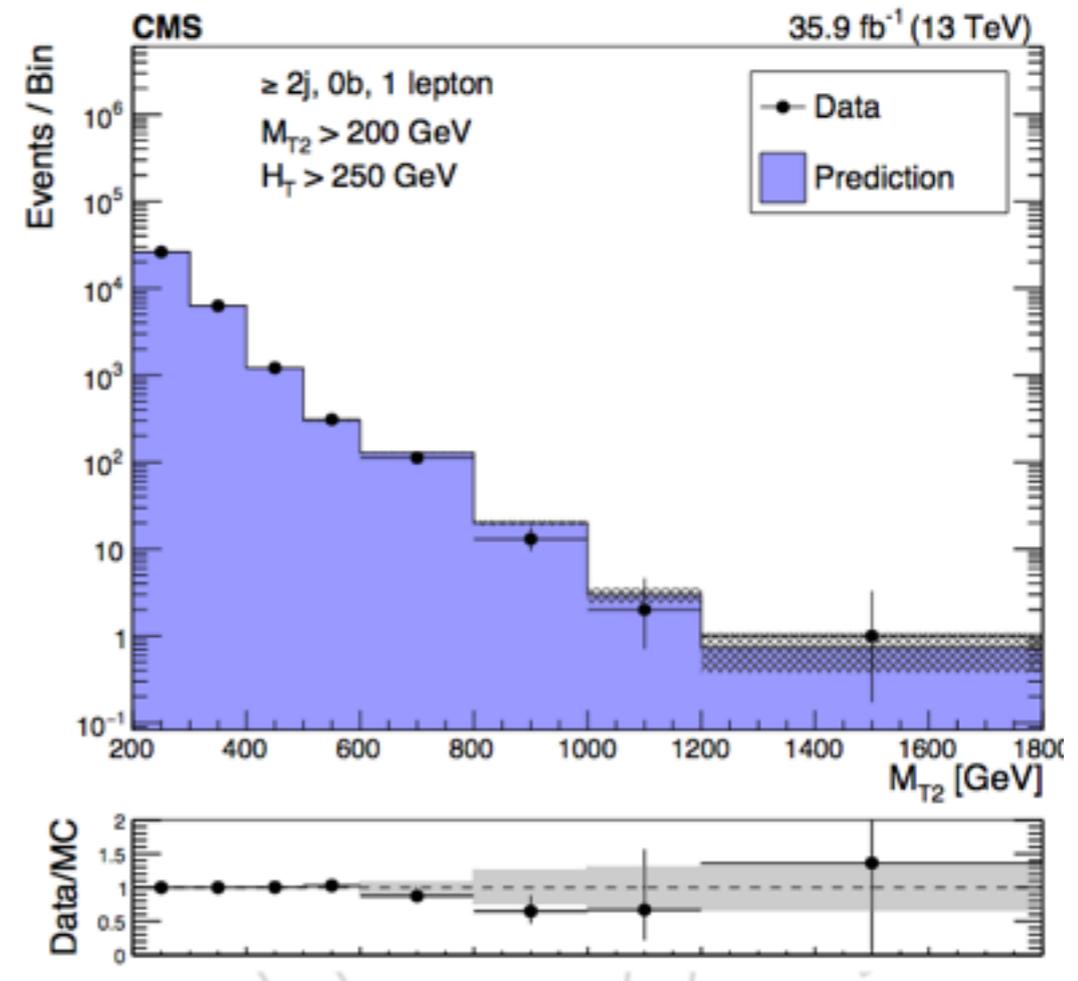
# Search strategy

- 64 search bins used



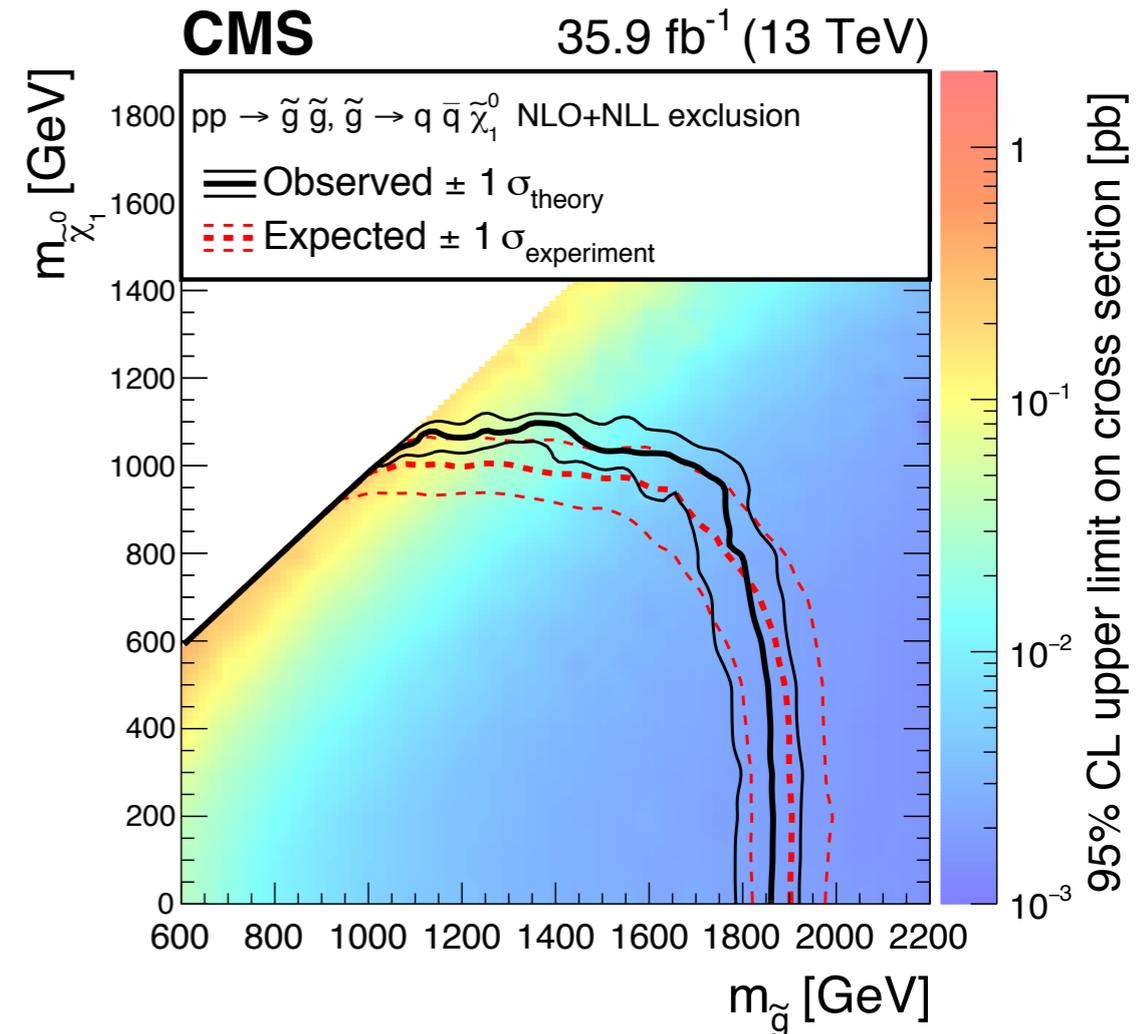
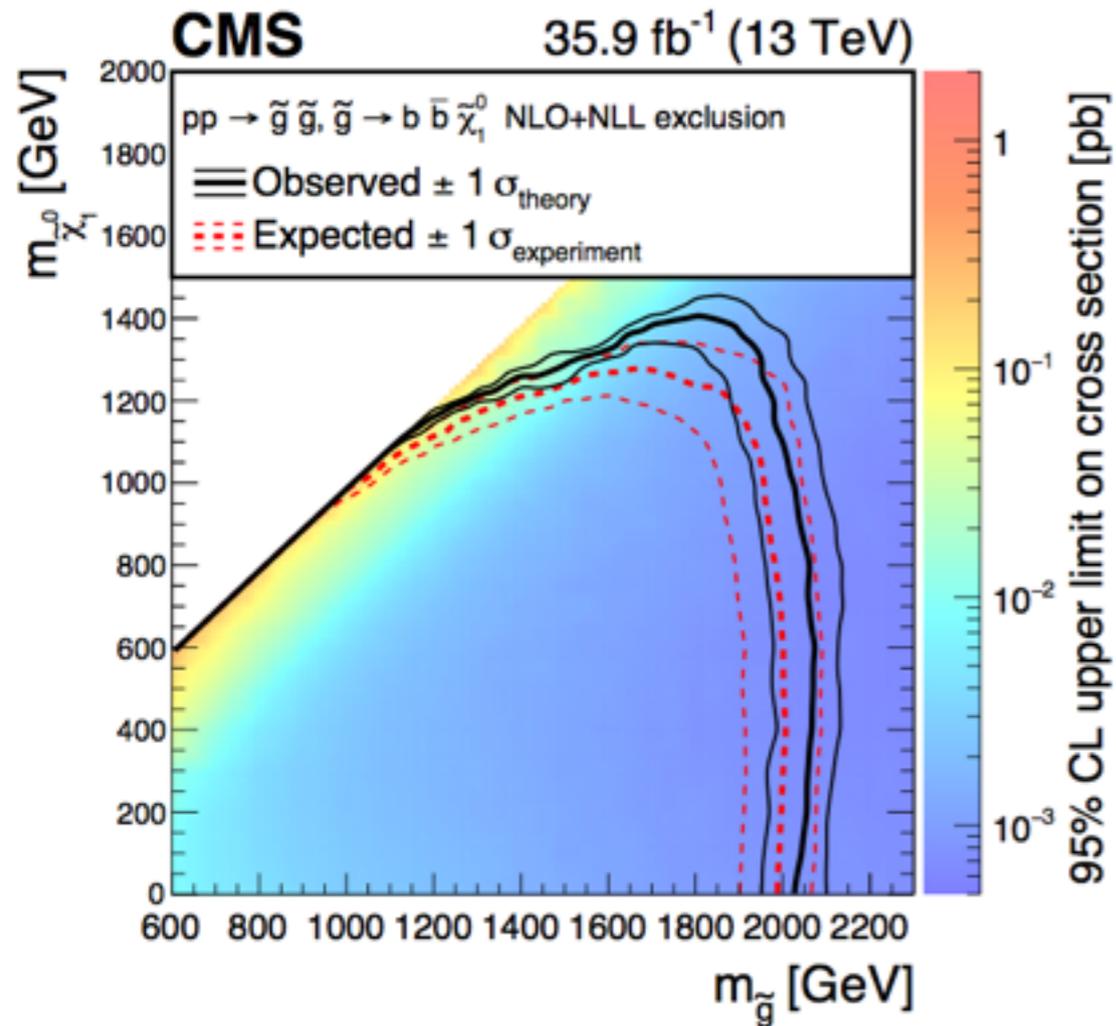
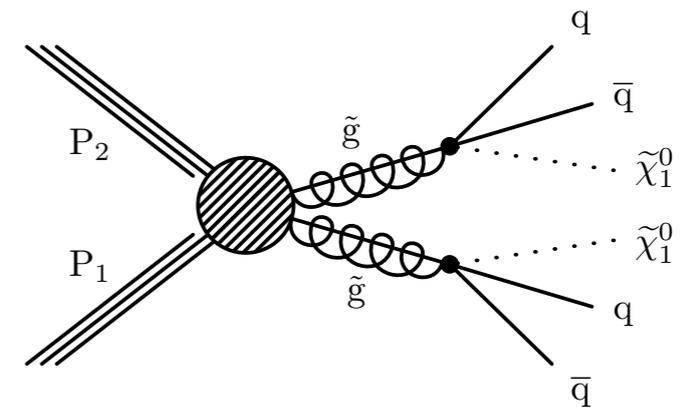
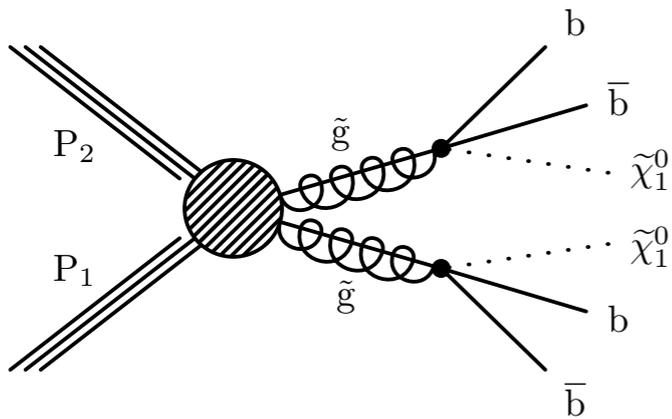
# Background Methods

- Lost Lepton
  - Transfer Factor related to observed counts in control region to expected counts in signal region
  - single lepton control region
  - Use the transfer factor to estimate background with modeling in  $M_{T2}$
- $Z(\nu\nu)+\text{Jets}$ 
  - Dilepton control region is used
  - Pro: Same process: small uncertainty on transfer factor
  - Con: small branching fraction
  - Photons Pro: Largest cross section
  - con: Different mass/couplings, larger uncertainty on transfer factor
- QCD
  - Residual QCD evaluated from control regions





# Limits



# Conclusion

- Two SUSY hadronic analysis are presented here: CMS-SUS-16-033, CMS-SUS-16-036
- Both analysis can serve as cross checks for the other and have found to give similar results in overlapping models
- No significant excess is seen
- The results seen in 2016 surpass what was seen Run1
- Stay tuned for 2017 and 2018!
- References:
  - <https://arxiv.org/pdf/1705.04650.pdf>
  - <https://arxiv.org/abs/1704.07781>

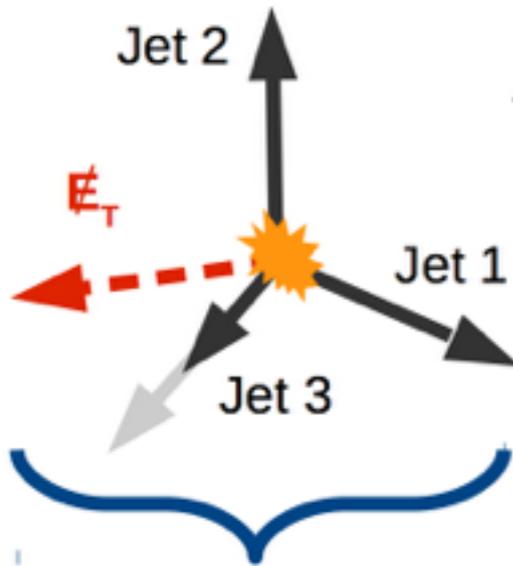


Thank you  
Back up

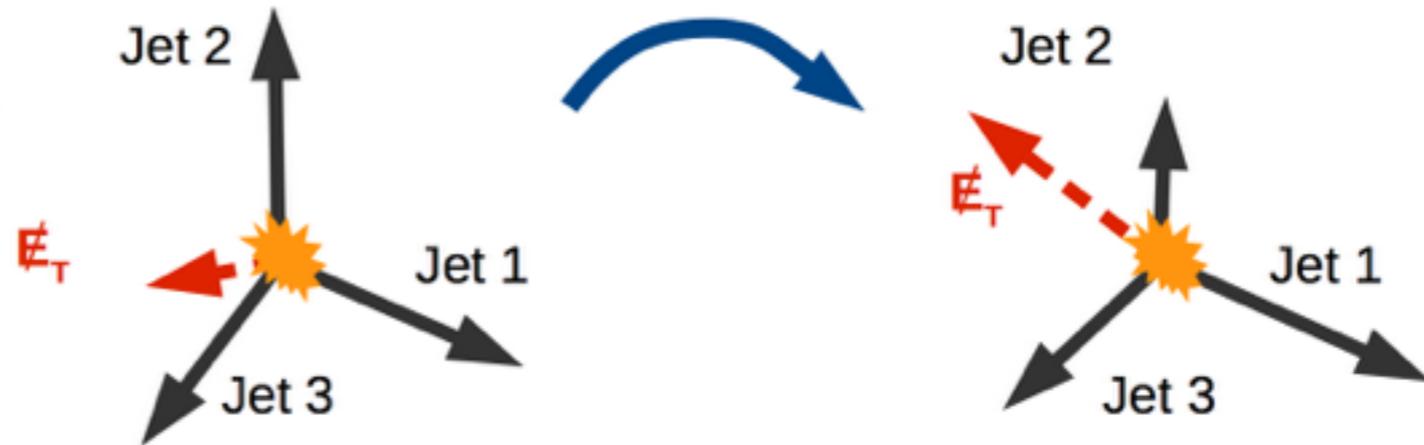


# QCD Estimate: Rebalance and Smear

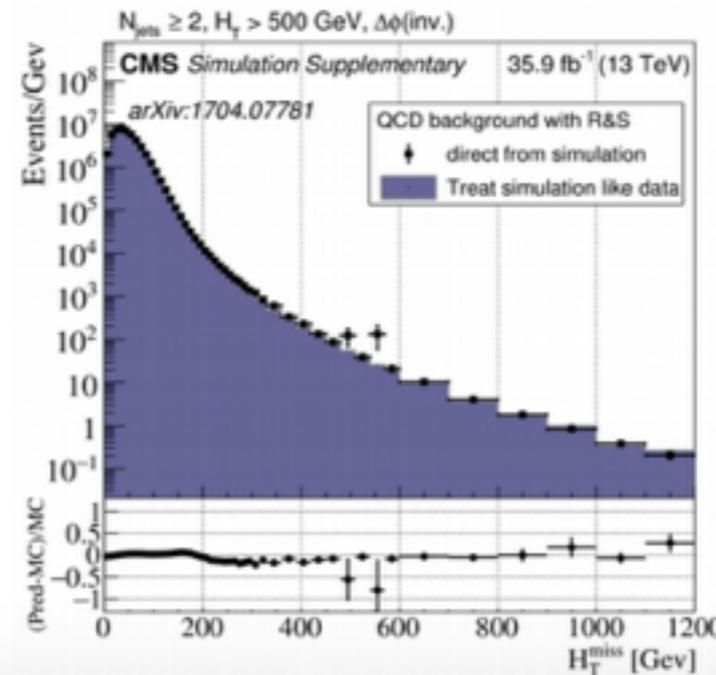
**Rebalance** jets to true hard scatter event with  $ME_T \approx 0$



**Smear** jets according to response



QCD multi-jet events have no intrinsic  $ME_T$ , only **instrumental  $ME_T$**  due to detector response that depends on  $\eta$  &  $p_T$  of jets

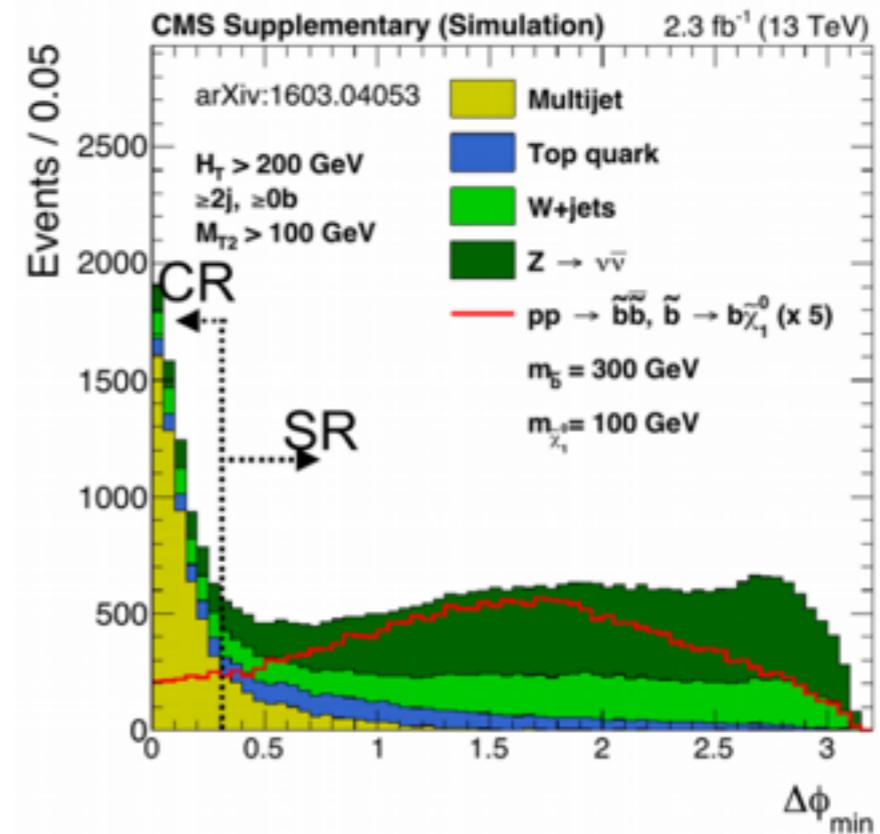


Good agreement with out of the box simulation

Myriam Schönenberger, ETH Zürich

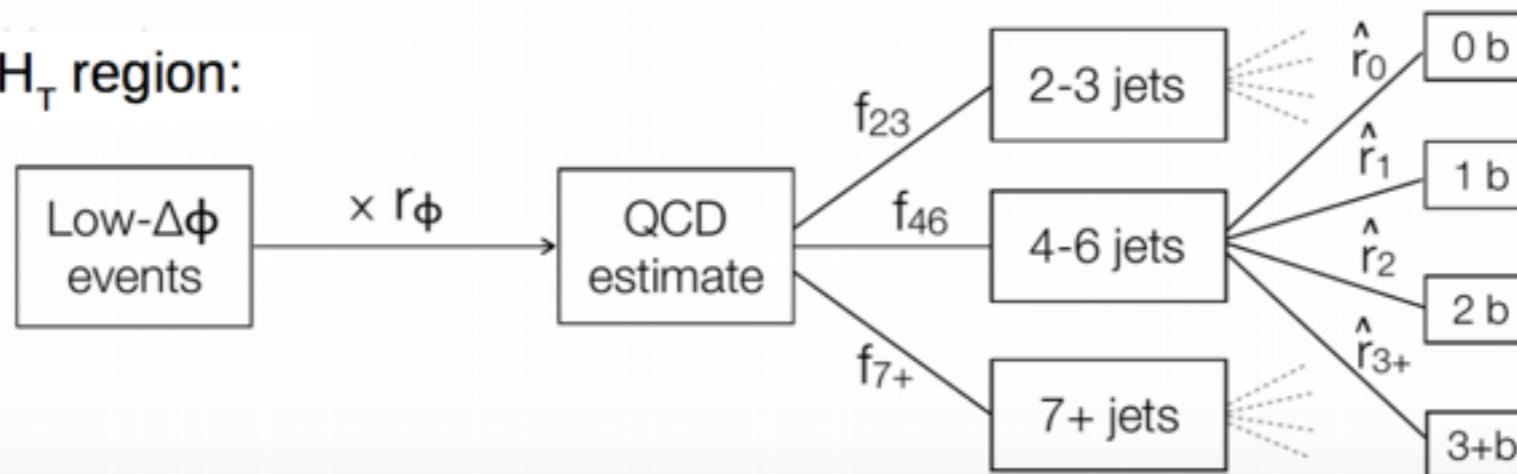
# QCD background estimate via delta Phi

- Invert  $\Delta\phi(ME_T, jets)$  cut
 
$$r_\phi = \frac{N(\Delta\phi_{min}(jets, E_T^{miss}) > 0.3)}{N(\Delta\phi_{min}(jets, E_T^{miss}) < 0.3)}$$
- Fit  $r_\phi$  at low  $M_{T2}$  & extrapolate to signal region inclusively in each  $H_T$  region
  - Then split among  $N_j/N_b$  with data based transfer factors
- $N_{CR}$  coming from signal triggers



$$N_{QCD}^{SR} = N^{CR}(H_T, M_{T2}) \cdot r_\phi(M_{T2}) \cdot f_j(H_T) \cdot r_b(N_j)$$

In each  $H_T$  region:



Myriam Schönenberger, ETH Zürich