

Key measurements for New-Physics in Flavor

J. Martin Camalich



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Outline

1 Intro

2 The $b \rightarrow c\tau\nu$ transitions

- Angular analyses
- The B_c -lifetime bound on 2HDM

3 The $b \rightarrow s\ell\ell$ decays

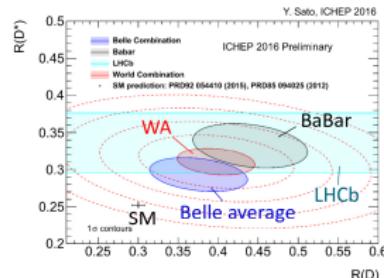
- Lepton universality violating observables
- Other observables

4 Decays of light hadrons

- Global fit to New Physics

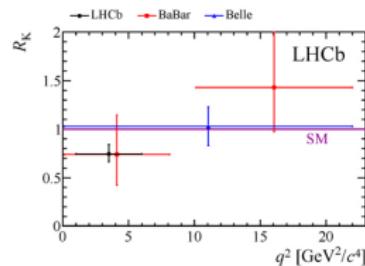
(Lepton universality violating) New-Physics in B decays?

- “ $R_{D^{(*)}}$ anomaly” in $B \rightarrow D^{(*)}\ell\nu$!



- “ R_K anomaly” in $B \rightarrow K\ell\ell$ (FCNC)!

LHCb PRL113(2014)151601



- Anomalies addressed in many models of NP (see e.g. V. Sudhir, J. Zupan's, S. Fajfer, ... talks)

- **Excesses** observed at $\sim 4\sigma$ WG2 on Th.
- Other “anomalies” in $b \rightarrow (u, c)\ell\nu$
 - Inclusive vs. Exclusive V_{ub} and V_{cb}
- $\Lambda_{\text{NP}} \sim 2 \text{ TeV}$

- Tension with SM $\sim 2.6\sigma$ WG3 on Tue.
- Other anomalies in $b \rightarrow s\mu\mu$
 - Branching fractions
 - Angular analysis $B \rightarrow K^* \mu\mu$
- Up to 4σ in global fits Javi Virto's talk
- $\Lambda_{\text{NP}} \sim 10 \text{ TeV}$

EFT: The bottom-up approach to new physics in flavor

No evidence of new-particles at colliders up to $E \simeq 1 \text{ TeV}$

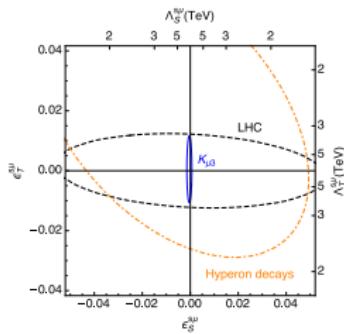
Construct the most general effective operators \mathcal{O}_k built with **all** the SM fields and subject to the strictures of $SU(3)_c \times SU(2)_L \times U(1)_Y$

Buchmuller *et al.*'86, Grzadkowski *et al.*'10

① Dictionary between low-energy flavor observables and high-energy (SMEFT) Wilson coefficients

- ▶ Complementary between **kaon decays** and LHC

$$\mathcal{L}_{\text{c.c.}} \supset -\frac{G_F V_{us}}{\sqrt{2}} \left[\epsilon_S (\bar{u} s) (\bar{e}(1-\gamma_5)\nu_e) + \epsilon_T (\bar{u}\sigma^{\mu\nu}s) (\bar{e}\sigma_{\mu\nu}(1-\gamma_5)\nu_e) \right]$$



Gonzalez-Alonso and JMC, arXiv 1605.07114

② Model-independent relations in the Wilson coefficients of the low-energy EFT

- ▶ Caveat in HEFT's Cata&Jung PRD92(2015)no.5,055018

① Charged-Currents (no RH ν):

$$\begin{aligned}\mathcal{L}_{\text{eff}}^{\ell} = & -\frac{G_F V_{cb}}{\sqrt{2}} [(1+\epsilon_L^{\ell}) \bar{\ell} \gamma_{\mu} (1-\gamma_5) \nu_{\ell} \cdot \bar{c} \gamma^{\mu} (1-\gamma_5) b + \epsilon_R^{\ell} \bar{\ell} \gamma_{\mu} (1-\gamma_5) \nu_{\ell} \cdot \bar{c} \gamma^{\mu} (1+\gamma_5) b \\ & + \bar{\ell} (1-\gamma_5) \nu_{\ell} \cdot \bar{c} [\epsilon_S^{\ell} + \epsilon_P^{\ell} \gamma_5] b + \epsilon_T^{\ell} \bar{\ell} \sigma_{\mu\nu} (1-\gamma_5) \nu_{\ell} \cdot \bar{c} \sigma^{\mu\nu} (1-\gamma_5) b] + \text{h.c.},\end{aligned}$$

- ▶ Matching to the SMEFT

$$\mathcal{O}_{Hud} = \frac{i}{\Lambda_{\text{NP}}^2} (\tilde{H}^\dagger D_\mu H)(\bar{u}_R \gamma^\mu d_R)$$

☛ RHC is lepton universal: $\epsilon_R^{\ell} \equiv \epsilon_R + \mathcal{O}(\frac{v^4}{\Lambda_{\text{NP}}^4}) \Rightarrow$ It cannot explain LUV in $R_{D^{(*)}}$!

Bernard, Oertel, Passemar & Stern PLB638(2006)480, Alonso, Grinstein and JMC JHEP 1510 (2015) 184

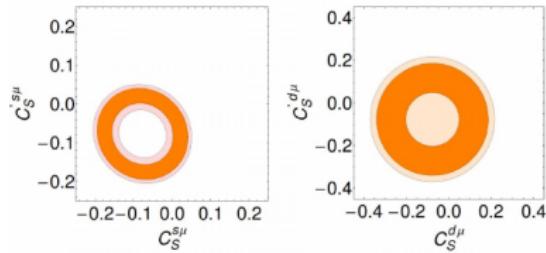
② Flavor Changing Neutral Currents:

- ▶ Important consequences in $B_s \rightarrow \mu\mu$

- ▶ Constraints LR operators: from **6** to **2**

★ No tensor operators!

Alonso, Grinstein, JMC, PRL113(2014)241802



☛ Only $\mathcal{O}_{9,10}^{(\prime)}$ can explain LUV in R_K !

Alonso, Grinstein, JMC, PRL113(2014)241802 (see also Hiller&Schmaltz'14, ...)

The $b \rightarrow c\tau\nu$ decays

New physics in τ decays?

- New physics..
 - ▶ could **look different** in different $\tau \rightarrow X\nu$ decay channels
 - ▶ should **appear universally** in all $Y \rightarrow Z\tau (\rightarrow X\nu)$ decays
- Leptonic τ decays: **Michel parameters**

A. Pich PPNP75(2014)41

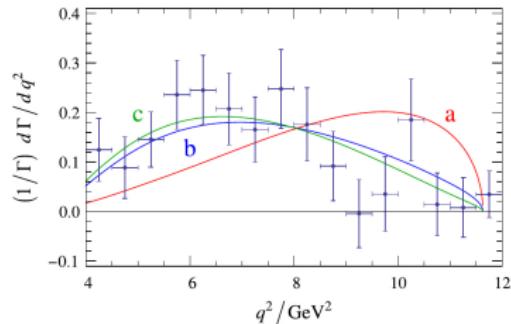
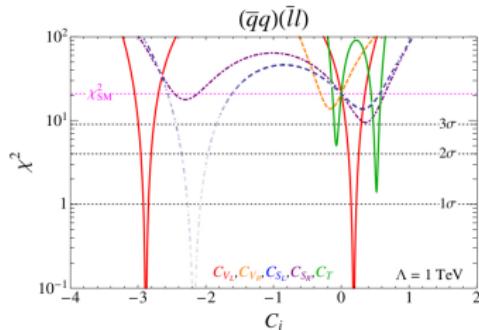
Bounds on the $g_{\ell\alpha}^V$ couplings, assuming that (non-standard) W -exchange is the only relevant interaction. The τ -decay (μ -decay) limits are at 95% CL (90% CL). Numbers within parentheses use μ -decay data through cross-channel identities.

	$ g_{RR}^V $	$ g_{LR}^V $	$ g_{RL}^V $	$ g_{LL}^V $
$\mu \rightarrow e$	<0.0004	<0.023	<0.017	>0.999
$\tau \rightarrow \mu$	<0.017 (0.003)	<0.12	<0.14 (0.023)	>0.983
$\tau \rightarrow e$	<0.017 (0.002)	<0.13	<0.13 (0.017)	>0.983

- ▶ 4-lepton (**pseudo**)scalar interactions less constrained
- Hadronic τ decays: **Lepton universality ratios**
$$R_{\tau/P} \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau P^-)}{\Gamma(P^- \rightarrow \mu^- \bar{\nu}_\mu)} = \left| \frac{g_\tau}{g_\mu} \right|^2 \frac{m_\tau^3}{2m_P m_\mu^2} \frac{(1 - m_p^2/m_\tau^2)^2}{(1 - m_\mu^2/m_p^2)^2} (1 + \delta R_{\tau/P})$$
$$\delta R_{\tau/\pi} = (0.16 \pm 0.14)\%$$
- ▶ New physics constrained at **subpercent level** in hadronic modes

Fits to new physics

- The $R_{D(*)}$ and spectrum do not provide enough discriminating power



Freytsis *et al.*, PRD92(2015)no.5, 054018; also Sasaki *et al.* PRD91(2015)no.11, 114028, ...

- Looking for LUR in new decay modes

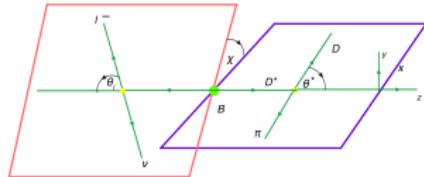
- ▶ Semileptonic $B_c \rightarrow J/\Psi \ell \nu$ decays Lytle *et al.* Wed. WG2 (arXiv: 1605.05645) and C. Bozzi Th. WG2
- ▶ $\Lambda_b \rightarrow \Lambda_c^{(*)} \ell \nu$ S. Meinel @ HC2NP-Tenerife
- ▶ $B_s \rightarrow D_s^{(*)} \ell \nu$ A. Bhol, EPL106(2014)31001

- Measure new observables

Look at the full kinematic (angular) distributions!

A first approach: Full angular analysis of the $B \rightarrow D^{(*)}\tau\nu$

- Many angular observables as functions of helicity amplitudes (like $B \rightarrow K^{(*)}\mu\mu$!)



$$\frac{d^4\Gamma}{dq^2 d\cos\theta_l d\cos\theta_{D^*} d\chi} = \frac{9}{32\pi} NF \{ \cos^2\theta_{D^*} (V_1^0 + V_2^0 \cos 2\theta_l + V_3^0 \sin 2\theta_l) + \sin^2\theta_{D^*} (V_1^T + V_2^T \cos 2\theta_l + V_3^T \sin 2\theta_l) \\ + V_4^T \sin^2\theta_{D^*} \sin^2\theta_l \cos 2\chi + V_5^{0T} \sin 2\theta_{D^*} \sin 2\theta_l \cos\chi + V_5^{0T} \sin 2\theta_{D^*} \sin\theta_l \cos\chi \\ + V_5^T \sin^2\theta_{D^*} \sin^2\theta_l \sin 2\chi + V_6^{0T} \sin 2\theta_{D^*} \sin\theta_l \sin\chi + V_6^{0T} \sin 2\theta_{D^*} \sin 2\theta_l \sin\chi \}$$

$$V_1^0 = 2 \left[\left(1 + \frac{m_l^2}{q^2} \right) (|A_0|^2 + 16|A_{0T}|^2) + \frac{2m_l^2}{q^2} |A_{lP}|^2 - \frac{16m_l}{\sqrt{q^2}} \text{Re}[A_{0T} A_0^*] \right],$$

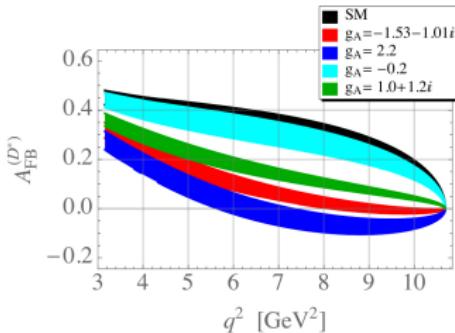
$$V_2^0 = 2 \left(1 - \frac{m_l^2}{q^2} \right) [-|A_0|^2 + 16|A_{0T}|^2],$$

$$V_3^0 = -8 \text{Re} \left[\frac{m_l^2}{q^2} A_{lP} A_0^* - \frac{4m_l}{\sqrt{q^2}} A_{lP} A_{0T}^* \right]$$

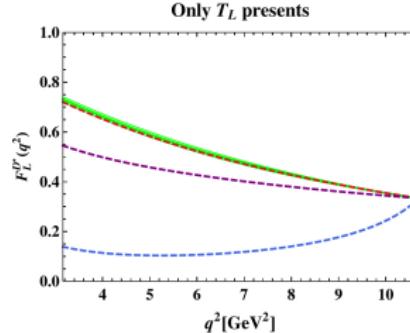
...

Duraisamy *et al.*, PRD90, 074013 (2014)

► Leptonic forward-backward asymmetry



► Longitudinal polarization of the D^*



Becirevic *et al.*, arXiv:1602.03030.

Duraisamy *et al.*, PRD90, 074013 (2014)

- **However** the τ 's lifetime is $\sim 10^{-13}$ s

- ▶ It is not *observed* but *reconstructed* from decay products with **missing neutrinos!**
(kinematics might be reconstructed from 3-prong τ decays) Bozzi's talk Th. WG2

- **Alternatively**, kinematic distributions of the **observable** τ decay products!

- ▶ Maximize the coverage of the τ 's lifetime

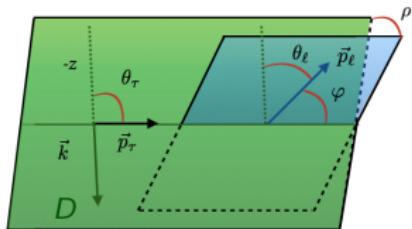
Channel	$\tau \rightarrow \mu \nu \nu$	$\tau \rightarrow e \nu \nu$	$\tau \rightarrow \pi \nu$	$\tau \rightarrow \rho \nu$	$\tau \rightarrow 3\pi \nu$	TOTAL
\mathcal{B}	17.4%	17.8%	10.82%	25%	9%	$\sim 80\%$

- ▶ **Different τ decay modes are subject to very different backgrounds!**
- ▶ Two different strategies in the literature
 - ★ **Analytical:** Nierste *et al.* PRD78,015006 '08 ($BD-\pi\nu$), Alonso, Kobach & JMC, PRD94 (2016) no.9, 094021 ($BD^{(*)}-\ell\nu\nu$), Alonso, JMC & Westhoff, in preparation ($BD^{(*)}-\pi\nu$ or $\rho\nu$)
 - ★ **Montecarlo:** Hagiwara *et al.* PRD89, 094009 (2014) ($BD-3\pi\nu$), Bordone *et al.* EPJC76 (2016) no.7, 360 ($BD-\ell\nu\nu$), Ligeti *et al.* arXiv:1610.02045 ($BD^*(\rightarrow D\pi)-\ell\nu\nu, \dots$)

Test new physics and understand better systematics!

$$B \rightarrow D^{(*)} \tau^- (\rightarrow \ell^- \bar{\nu}_\ell \nu_\tau) \bar{\nu}_\tau$$

Alonso, Kobach, JMC, arXiv: 1602.07671



- Integrate analytically τ angular phase-space: (nontrivial)

$$\frac{d^3\Gamma_5}{dq^2 dE_\ell d(\cos \theta_\ell)} = \mathcal{B}[\tau_\ell] \frac{G_F^2 |V_{cb}|^2 \eta_{EW}^2}{32\pi^3} \frac{|\vec{k}|}{m_B^2} \left(1 - \frac{m_\tau^2}{q^2}\right)^2 \frac{E_\ell^2}{m_\tau^3} \times [I_0(q^2, E_\ell) + I_1(q^2, E_\ell) \cos \theta_\ell + I_2(q^2, E_\ell) \cos \theta_\ell^2]$$

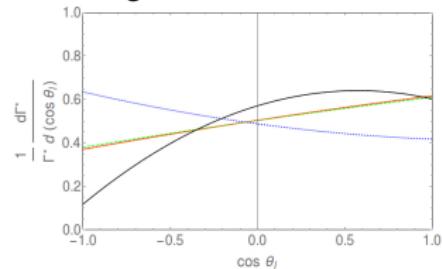
- ▶ $\cos \theta_\ell$ defined as for the normalization mode (w.r.t recoiling $D^{(*)}$ in the q rest frame)
- ▶ $I_{0,2}(q^2, E_\ell)$ accessed in $R_{D^{(*)}}$
- ▶ $I_1(q^2, E_\ell)$ accessible only with a FB leptonic asymmetry!

$$\frac{d^2 A_{FB}(q^2, E_\ell)}{dq^2 dE_\ell} = \left(\int_0^1 d(\cos \theta_\ell) - \int_{-1}^0 d(\cos \theta_\ell) \right) \frac{d^3\Gamma_5}{dq^2 dE_\ell d(\cos \theta_\ell)}$$

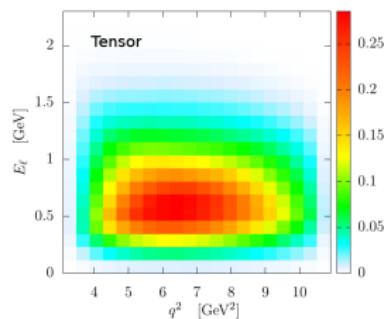
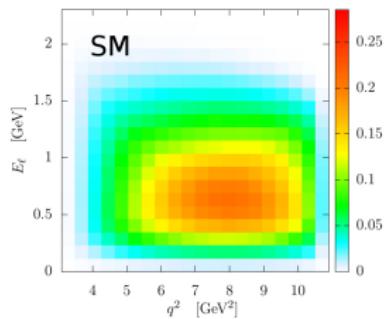
$$R_{FB}^{(*)} = \frac{1}{\mathcal{B}[\tau_\ell]} \frac{1}{\Gamma_{\text{norm}}} A_{FB},$$

- Angular distributions can help to discriminate signal vs. normalization

	R_{D^*}	R_{FB}^*
SM	0.252(4)	0.0310(7)
Current	0.333	0.0410
Scalar	0.315	0.0363
Tensor	0.346	-0.0377
Expt.	0.322(18)(12)	-



- E_ℓ and double (E_ℓ, q^2) spectra can also be studied



Ligeti *et al.* arXiv:1610.02045

$B \rightarrow D^{(*)}\tau^- (\rightarrow \pi^-\nu_\tau)\bar{\nu}_\tau$: $\tau^- \rightarrow \pi^-\nu_\tau$ as a τ polarimeter

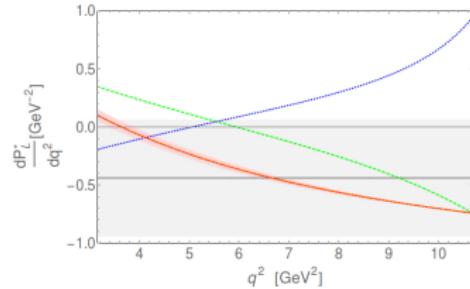
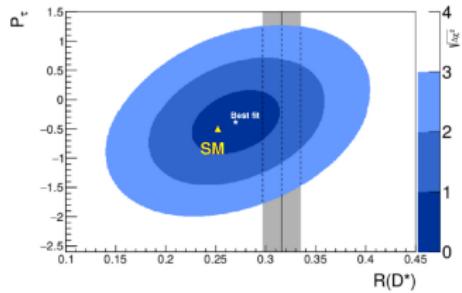
$$\frac{d^2\Gamma_4}{dq^2 dE_\pi} = \frac{\mathcal{B}[\tau\pi] m_\tau^2}{|\vec{p}_\tau|(m_\tau^2 - m_\pi^2)} \frac{d\Gamma_B}{dq^2} \left[1 + \xi(E_\pi, q^2) \frac{dP_L}{dq^2} \right], \quad \xi(E_\pi, q^2) = \frac{1}{\beta_\tau} \left(2 \frac{E_\pi}{E_\tau} - 1 \right)$$

Tanaka&Watanabe, PRD82, 034027 (2010)

Slope in E_π of $d\Gamma_4 \Rightarrow$ Longitudinal Polarization

$$\frac{dP_L}{dq^2} = \frac{d\Gamma_{B,+}/dq^2 - d\Gamma_{B,-}/dq^2}{d\Gamma_B/dq^2}$$

- Applied to the BD^* channel by **Belle**



Belle, arXiv: 1608.06391, K. Adamczyk's talk Th. WG2

$B \rightarrow D^{(*)}\tau^- (\rightarrow \pi^-\nu_\tau)\bar{\nu}_\tau$: $\tau^- \rightarrow \pi^-\nu_\tau$ as a τ polarimeter

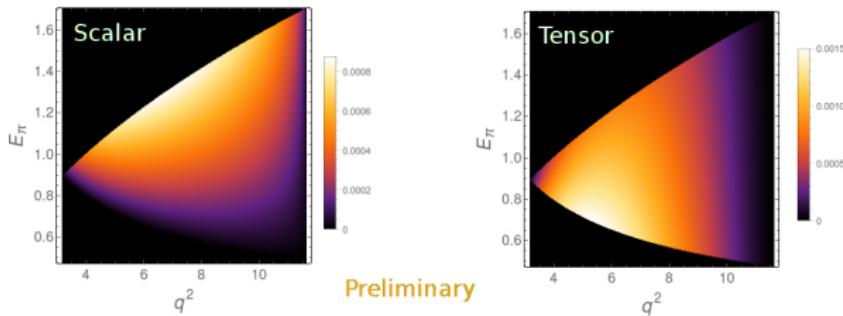
$$\frac{d^2\Gamma_4}{dq^2 dE_\pi} = \frac{\mathcal{B}[\tau_\pi] m_\tau^2}{|\vec{p}_\tau|(m_\tau^2 - m_\pi^2)} \frac{d\Gamma_B}{dq^2} \left[1 + \xi(E_\pi, q^2) \frac{dP_L}{dq^2} \right], \quad \xi(E_\pi, q^2) = \frac{1}{\beta_\tau} \left(2 \frac{E_\pi}{E_\tau} - 1 \right)$$

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$$\frac{dP_L}{dq^2} = \frac{d\Gamma_{B,+}/dq^2 - d\Gamma_{B,-}/dq^2}{d\Gamma_B/dq^2}$$

- Scalar and Tensor modify $d^2\Gamma_4$ up to 50% in some regions of phase-space

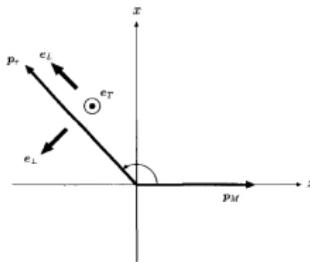


Alonso, JMC & Westhoff, to appear

More $B \rightarrow D^{(*)}\tau\nu$ observables from FB asymmetries

PRELIMINARY: Alonso, JMC & Westhoff to appear

- Decay rate into a τ polarized along a given direction \hat{s}



$$d\Gamma_B(\hat{s}) = d\Gamma + \frac{1}{2} d\Gamma \left(dP_L \hat{z}' + dP_{\perp} \hat{x}' + dP_T \hat{y}' \right) \cdot \hat{s}$$

- dP_L measured by Belle
- dP_T (T -odd): Not accessible without τ direction
- dP_{\perp} accessible from the pionic FB asymmetry!

Tanaka Z. Phys. C 67, 321

- Master formula for the BD mode

$$\frac{d^2\Gamma_{FB}^{N_X}}{dq^2 dE_\pi} = -\frac{\mathcal{B}[\tau\pi] m_\tau^2}{|\vec{p}_\tau|^3} \frac{d\Gamma_B}{dq^2} \left(\xi_{FB}^X(E_\pi, q^2) \frac{dA_{FB}}{dq^2} + \xi_{\perp}(E_\pi, q^2) \frac{dP_{\perp}}{dq^2} \right)$$

- $\xi_{FB,\perp}^X(q^2, E_X)$ depend on the decay channel $\tau \rightarrow X \nu_\tau$

FB asymmetries allows us to measure 3-body decay observables A_{FB} and P_{\perp} !

The lifetime of the B_c meson and new physics

- 4 Wilson coefficients can explain R_{D^*} anomalies: $\epsilon_L, \epsilon_S, \epsilon_P, \epsilon_T$
 - ▶ 2HDM with discrete symmetries:

$$\epsilon_S \simeq \epsilon_P \simeq \xi^b \xi^\ell \frac{m_b m_\ell}{m_{H^\pm}^2}$$

- BD^* depends on ϵ_P Browder's talk
- $B_c \rightarrow \tau \nu$ also receives a chiral enhanced contribution from ϵ_P !

$$\text{Br}(B_c^- \rightarrow \tau \bar{\nu}_\tau) = \tau_{B_c^-} \frac{m_{B_c} m_\tau^2 f_{B_c}^2 G_F^2 |V_{cb}|^2}{8\pi} \left(1 - \frac{m_\tau^2}{m_{B_c}^2}\right)^2 \left|1 + \epsilon_L + \frac{m_{B_c}^2}{m_\tau(m_b + m_c)} \epsilon_P\right|^2$$

The constraint from the B_c lifetime

- ▶ QCD: "Most of the B_c lifetime comes from $\bar{c} \rightarrow \bar{s}$ ($\sim 65\%$) and $b \rightarrow c$ ($\sim 30\%$) transitions"

Bigi PLB371 (1996) 105, Beneke *et al.* PRD53(1996)4991, Kiselev *et al.* NPB585 (2000) 353

👉 $\text{BR}(B_c \rightarrow \tau \nu) \leq 30\% \text{ (accounting conservatively for errors)}$

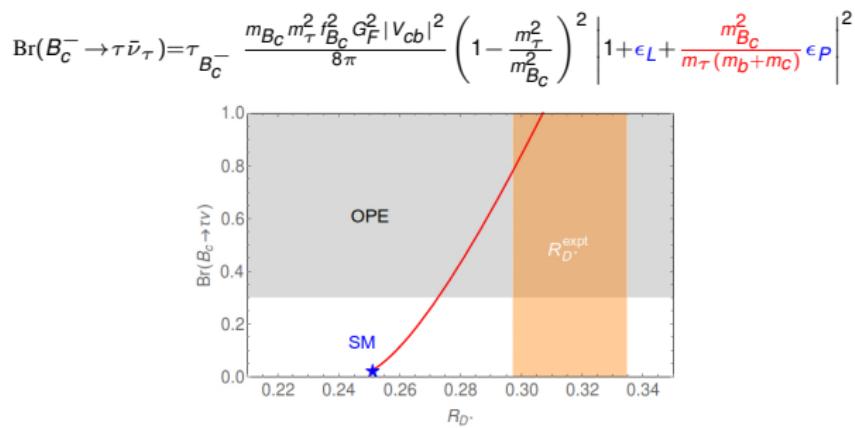
Alonso, Grinstein&JMC, arXiv: 1611.06676 (see also Xin-Qiang Li *et al.*, JHEP 1608 (2016) 054)

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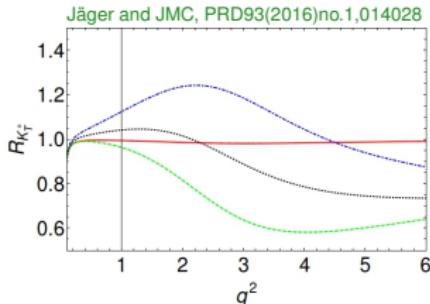


τ_{B_c} makes highly implausible ANY “scalar solution”
(e.g. 2HDM) to the R_{D^*} anomaly!

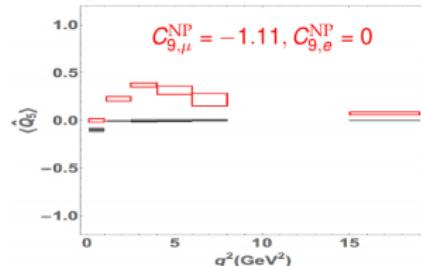
The $b \rightarrow s\ell\ell$ decays

Lepton Universality in $b \rightarrow s\ell\ell$

- **Theory:** Plenty of LUR's in the $B_{(s)} \rightarrow K^*(\Phi)\ell\ell$ angular analysis!

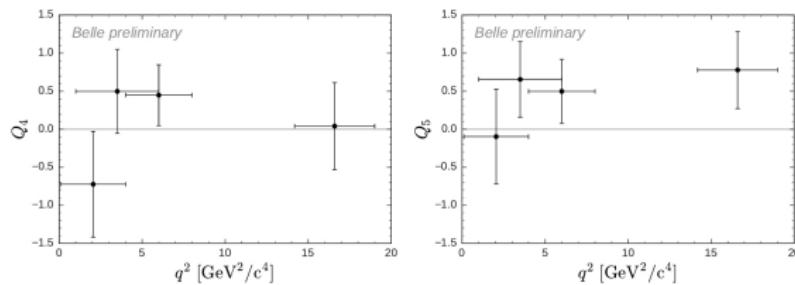


Quim's talk, Capdevila et al. JHEP 1610 (2016) 075



- **Experiment:** New results from Belle in angular observables!

S. Wehle's talk



Eagerly waiting for the new LHCb data on $R_{K^{(*)}}$ and R_Φ (other angular observables?)

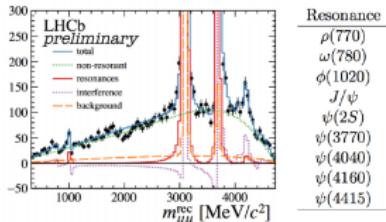
Paula Alvarez's talk

Angular analyses of $B_{(s)} \rightarrow K^*(\phi)\mu\mu$ and radiative

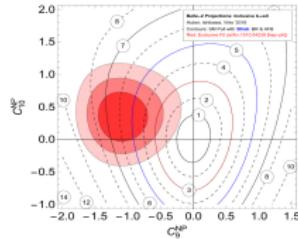
Nature and size of hadronic corrections still a **controversial** topic

- New experimental information and **prospects**
 - ▶ Phases and fudge factors of resonances
 - ▶ Future interplay with $B \rightarrow X_s \ell \ell$ at Belle2

Ponci's talk ($B \rightarrow K\mu\mu$)



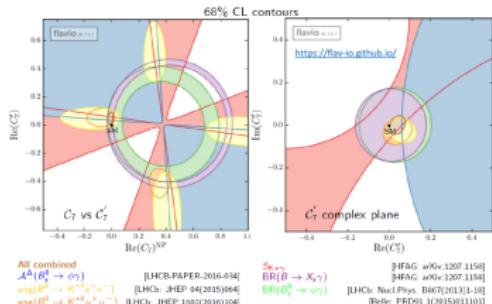
Virto's talk



- Radiative and $B \rightarrow K^* ee$ at very low q^2 very clean! Jäger and JMC, PRD93(2016)no.1,014028

- ▶ New data on $\text{BR}(B_s \rightarrow \Phi\gamma)$ and \mathcal{A}_Δ !
Arantxa Oranguren's and Polci's talk
 - ▶ Stringent constraints on $C_7^{(\prime)}$

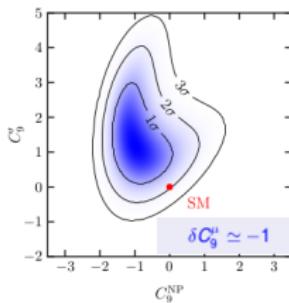
Paul & Straub arXiv:1608.02556



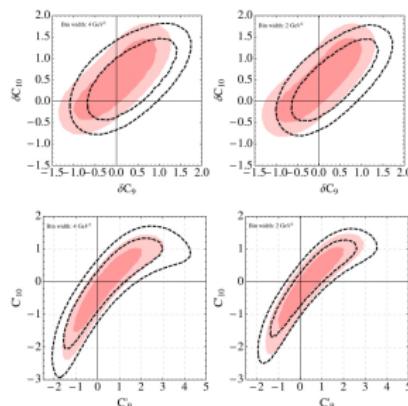
What about the high q^2 region?

No satisfactory (model-independent) solution (yet?)

- FFs in LQCD+OPE Wingate's talk



- Duality violation Braß et al. arXiv:1606.00775



- Some analyses obtain large RHC

Rusa Mandal's talk

- New ideas: Analyse $B_s^* \rightarrow \ell\ell$ and $\ell\ell \rightarrow B_s^* \rightarrow B_s \gamma$! B. Grinstein and JMC PRL116(2016)no.14,141801

- Cleanest (non-LUR) observables sensitive to C_9

- Matrix elements of local quark currents are described by decay constants
- $m_{B_s}^2 \sim 30 \text{ GeV}^2$ well above charm resonances (use OPE!)

- Very challenging experimentally

- $\text{BR} \sim 10^{-11}$ not far from $B_s \rightarrow \mu\mu$; X-sections receive resonant enhancement $\sigma \sim 10 \text{ fb}$
- Might be useful in the charm sector Khodjamirian et al. JHEP 1511 (2015) 142

New physics in semileptonic decays of light hadrons

New-physics in light quark charged-current transitions?

- ① **Experimental data:** $K_{\ell 2}$ and $K_{\ell 3}$ boast an extremely rich database

Cecucci's, Moulson's talks, FlaviaNet Kaon Working Group, Antonelli *et al.* EPJC69, 399 (2010)

- ② **Hadronic matrix elements:** Flagship quantities in χ PT and LQCD

FLAG collaboration, Simula's talk

- ③ Radiative and isospin-breaking corrections understood!

Cirigliano *et al.* Rev.Mod.Phys. 84 (2012) 399

LUR's predicted at the sub-permille level!

$$R_P = \frac{\Gamma(P_{e2(\gamma)})}{\Gamma(P_{\mu 2(\gamma)})}$$

$$R_\pi^{\text{SM}} = 1.2352(1) \times 10^{-4}$$

$$R_\pi = 1.2344(30) \times 10^{-4}$$

$$R_K^{\text{SM}} = 2.477(1) \times 10^{-5}$$

$$R_K = 2.488(9) \times 10^{-5}$$

Cirigliano & Rosell PRL99 (2007) 231801

A Global NP analysis of the light-quark meson decay data is called for!

$K_{\ell 3}$

- Neglecting contributions $\mathcal{O}(\frac{v^4}{\Lambda_{\text{NP}}^4})$ (terms quadratic in ϵ_i^2)

$$\Gamma(K_{\ell 3(\gamma)}) = \underbrace{\frac{G_F^2 m_K^5}{192\pi^3} C S_{\text{EW}} |\tilde{V}_{us}^\ell|^2 f_+(0)^2}_{\left(1 + \epsilon_L^{s\ell} + \epsilon_R^s - \tilde{\nu}_L\right) V_{us}^{\text{SM}}} \underbrace{l_K^\ell(\lambda_{+,0}, \epsilon_S^{s\ell}, \epsilon_T^{s\ell})}_{\text{Phase-space Int.}} \underbrace{\left(1 + \delta^c + \delta_{\text{em}}^{c\ell}\right)^2}_{\text{Rad. and isosp. corr.}}$$

- $f_+(0)$, δ^c and $\delta_{\text{em}}^{c\ell}$ th. inputs (LQCD and χ PT)
- $\epsilon_{S,T}^{s\ell}$ accessible through the spectra/angular distribution

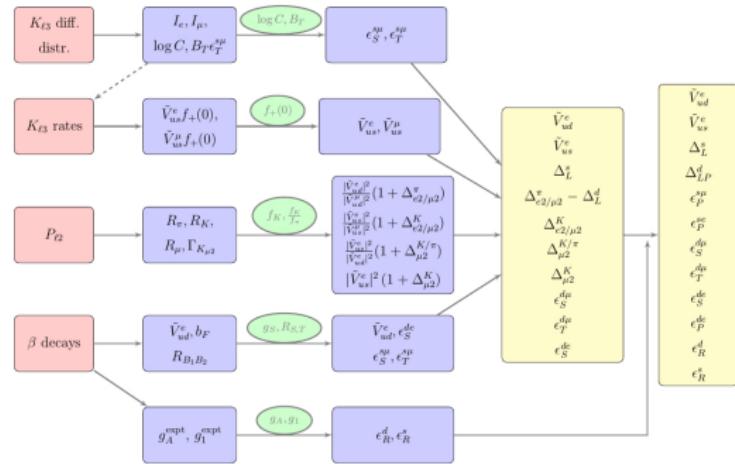
Interference with SM is $\propto m_\ell$!

- K_{e3} spectra is SM-like! (sensitivity to $|\epsilon_{S,T}^{se}|^2$)
- $K_{\mu 3}$ sensitive \Rightarrow Simultaneous fit of $\lambda_{+,0}$, $\epsilon_S^{s\mu}$, $\epsilon_T^{s\mu}$

- $|\tilde{V}_{us}^\ell|$ only accessible through CKM unitarity and LUV tests
 - Less NP-polluted for K_{e3}
 - Cross-contamination from NP in μ decays

Flowchart of data, theoretical inputs and outputs of the fit

M. Gonzalez-Alonso & JMC, arXiv:1605.07114, JHEP in print



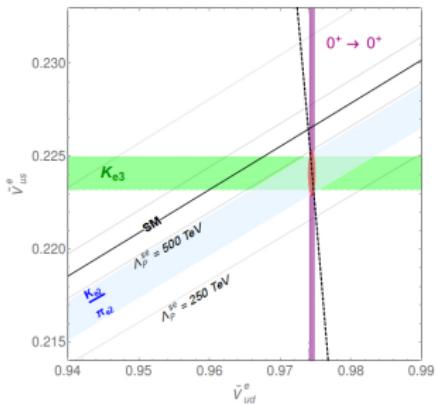
- Nuclear, neutron and hyperon β decay data essential!
 - ▶ Nuclear decays provide the most precise determination of $|\tilde{V}_{ud}^e|$ Hardy's talk
 - ▶ Axial nucleon/hyperon couplings important probes of **RHC**
 $g_1(0) = (1 - 2\epsilon_R^s)g_1(0)^{\text{QCD}}$
 - ▶ Hyperon decays very sensitive to **tensor currents** Chang *et al.* PRL114 (2015) no.16, 161802
 - ▶ ...
- **Very old data base:** Plenty of room for improvement at LHCb and NA62!

● Results of the fit

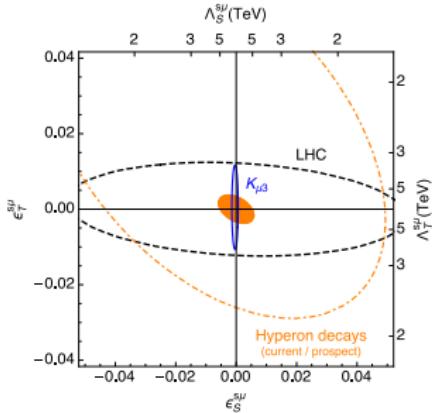
$$\begin{pmatrix} \hat{V}_{ud}^e \\ \hat{V}_{us}^e \\ \hat{V}_{ub}^e \\ \Delta_L^e \\ \Delta_{LF}^e \\ \epsilon_{\bar{K}}^{ee} \\ \epsilon_{\bar{K}}^{es} \\ \epsilon_{\bar{K}}^{eu} \\ \epsilon_{\bar{K}}^{us} \\ \epsilon_{\bar{K}}^{ub} \\ \epsilon_{\bar{S}}^{eu} \\ \epsilon_{\bar{S}}^{us} \\ \epsilon_{\bar{S}}^{ub} \end{pmatrix} = \begin{pmatrix} 0.97451 \pm 0.00038 \\ 0.22408 \pm 0.00087 \\ 1.1 \pm 3.2 \\ 1.9 \pm 3.8 \\ 4.0 \pm 7.8 \\ -1.3 \pm 1.7 \\ -0.4 \pm 2.1 \\ -0.7 \pm 4.3 \\ 0.1 \pm 5.0 \\ -3.9 \pm 4.9 \\ 0.5 \pm 5.2 \\ 1.4 \pm 1.3 \end{pmatrix} \times 10^{-8}$$

$\rho = \begin{pmatrix} 1. & 0. & 0. & 0.01 & 0.01 & 0. & 0. & 0. & 0. & 0. & 0. & 0. & 0. & 0.82 \\ -1. & -0.12 & 0. & 0. & 0. & 0.04 & 0.04 & 0. & -0.26 & 0. & 0. \\ - - & 1. & 0. & 0. & 0. & 0. & 0.03 & 0. & 0. & 0.72 & 0. \\ - - & - & 1. & 0.9995 & -0.87 & 0.09 & 0.09 & 0. & 0.04 & 0. & 0.01 \\ - - & - & - & 1. & -0.87 & 0.09 & 0.09 & 0. & 0.04 & 0. & 0.01 \\ - - & - & - & - & 1. & 0. & 0. & 0. & 0. & 0. & 0. \\ - - & - & - & - & - & 1. & 0.9993 & -0.98 & -0.01 & 0. & 0. \\ - - & - & - & - & - & - & 1. & -0.98 & -0.01 & 0.02 & 0. \\ - - & - & - & - & - & - & - & 1. & 0. & 0. & 0. \\ - - & - & - & - & - & - & - & - & 1. & 0. & 0. \\ - - & - & - & - & - & - & - & - & - & 1. & 0. \\ - - & - & - & - & - & - & - & - & - & - & 1. \end{pmatrix}$

- Re-interpretation of *Classic* $|V_{ud}| - |V_{us}|$ plot



- Interplay kaons, hyperons and LHC



Conclusions

① The $b \rightarrow c\tau\nu$ decays

- ▶ Angular analyses and new decay modes will settle this issue
- ▶ The lifetime of the B_c meson makes highly implausible 2HDMs interpretations
- ▶ Lack of time:
 - ★ NP in $B \rightarrow D^{(*)}\ell\nu$ (angular) observables and $|V_{cb}|$ Tayduganov's talk, Colangelo *et al.* arXiv:1611.07387, ...
 - ★ EW mixing and interplay with colliders Ferruglio *et al.* arXiv:1606.00524, Faroughy *et al.* PLB764(2017)126

② The $b \rightarrow s\ell\ell$ decays

- ▶ Lepton Universality ratios
- ▶ New interesting data on charm and radiative
- ▶ Lack of time:
 - ★ LFV; $b \rightarrow s\tau\tau$ very enhanced in models aligning R_K and $R_{D^{(*)}}$ with flavor

③ (Semi)leptonic decays of light quarks

- ▶ Boasts an exquisite data base and good theoretical understanding
- ▶ Global fits to NP → Reinterpretation of $|V_{us}|-|V_{ud}|$ plot, interplay with colliders, ...
- ▶ Lack of time:
 - ★ Extensions to global fits to (semi)leptonic charm meson decays Fajfer's talk
 - ★ Rare kaon and hyperon decays (LHCb upcoming program in strange physics!)