Rare Kaon Decays

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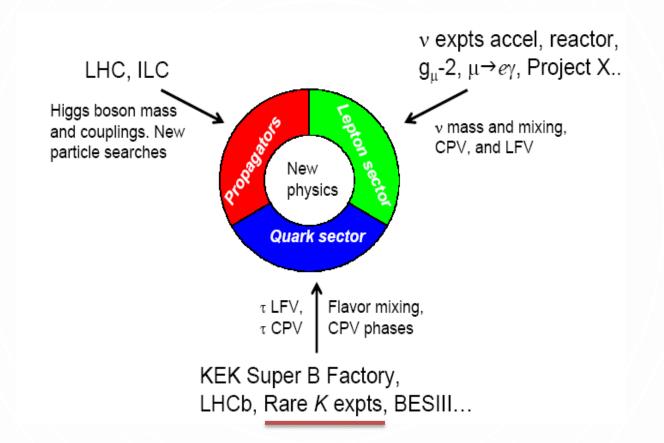
Project X meeting



Jan13-14, 2011









- Search for explicit violation of Standard Model
 Lepton Flavor Violation
- Probe the flavor sector of the Standard Model FCNC
- Test fundamental symmetriesCP, CPT
- ➤ Study the strong interactions at low energy
 Chiral Perturbation Theory, K structure
- Exploring lepton mass matrix
 Unique possibility of measuring double beta decay analogue for μ.



Rare decays: K_L decays

 $\Delta S = 1$

 $= (2.7 \pm 0.4) \times 10^{-11}$

Mode Expt. value

•
$$\not\simeq K_L \to \pi^0 \nu \overline{\nu}$$
 < 6.7 \times 10⁻⁸

•
$$K_L \rightarrow \pi^0 \pi^0 \nu \overline{\nu}$$
 $< 4.7 \times 10^{-5}$

•
$$\not\approx K_L \to \pi^+\pi^-e^+e^- < (3.11 \pm 0.19) \times 10^{-7}$$

•
$$K_L \to \pi^0 \pi^0 e^+ e^- < 6.6 \times 10^{-9}$$

•
$$K_L \rightarrow \mu^+ \mu^-$$
 (6.84 ± 0.11) × 10⁻⁹

•
$$K_L \rightarrow e^+e^ (9^{+6}_{-4}) \times 10^{-12}$$

•
$$K_L \to \pi^0 \mu^+ \mu^- < 3.8 \times 10^{-10}$$

•
$$K_L \to \pi^0 e^+ e^- < 2.8 \times 10^{-10}$$

•
$$K_L \to e^{\pm} \mu^{\mp}$$
 < 4.7 × 10⁻¹²

•
$$K_L \to e^{\pm}e^{\pm}\mu^{\mp}\mu^{\mp} < 4.12 \times 10^{-11}$$

•
$$K_L \to \pi^0 e^{\pm} \mu^{\mp}$$
 < 7. 6 × 10⁻¹¹

•
$$K_L \to \pi^0 \pi^0 e^{\pm} \mu^{\mp} < 1.7 \times 10^{-10}$$



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K⁺decays

$$\Delta S = 1$$

Lepton Family Number

Lepton number

Angular momentum

•
$$\not\cong K^+ \to \pi^+ \nu \overline{\nu}$$
 $(1.7 \pm 1.1) \times 10^{-10}$ $= (8.5 \pm 0.07) \times 10^{-11}$

•
$$K^+ \to \pi^+ \pi^0 \nu \overline{\nu}$$
 < 4.3 × 10⁻⁵

•
$$K^+ \to \pi^+ e^+ e^ (3.00 \pm 0.09) \times 10^{-7}$$

•
$$K^+ \to \pi^+ \mu^+ \mu^-$$
 (8.1 ± 1.4) × 10⁻⁸

•
$$K^+ \to \mu^- \nu e^+ e^+$$
 2.0 × 10⁻⁸

•
$$K^+ \to \pi^+ \mu^+ e^- < 1.3 \times 10^{-11}$$

• $K^+ \to \pi^+ e^+ \mu^- < 5.2 \times 10^{-10}$

$$< 5.2 \times 10^{-10}$$

•
$$K^+ o \pi^- \mu^+ e^+ < 5.0 imes 10^{-10}$$

• $K^+ o \pi^- e^+ e^+ < 6.4 imes 10^{-10}$

•
$$K^+ \rightarrow \pi^- \mu^+ \mu^+$$
 < 3.0 × 10⁻⁹

•
$$K^+ \to \pi^- \mu^+ \mu^+ < 3.0 \times 10^{-9}$$

•
$$K^+ \to \pi^+ \gamma$$
 < 2.3 × 10⁻⁹

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$

Theoretically clean mode

Buras: $K^+ \to \pi^+ \nu \overline{\nu}$, $K^0 \to \pi^0 \nu \overline{\nu}$, ratio x_d/x_s of $B_d^0 - \overline{B_d^0}$ to $B_s^0 - \overline{B_s^0}$ mixing and class of asymmetries in neutral B decays cleanest observables, being essentially free from hadronic uncertainties.

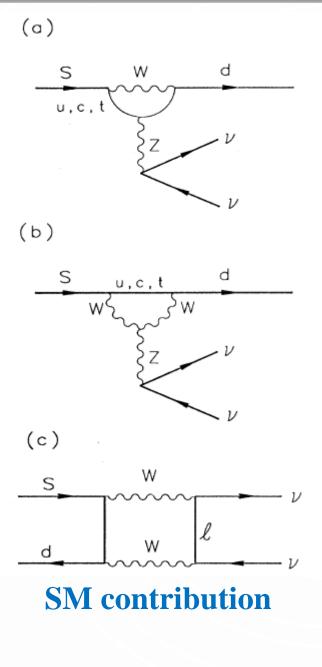
Hadronic matrix element of the operator

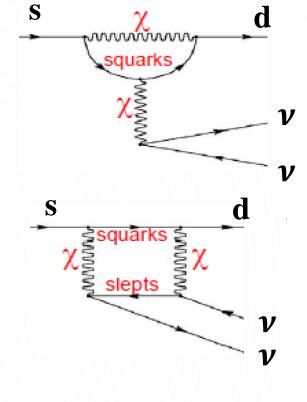
$$\overline{s}\gamma_{\mu}(1-\gamma_5)d\overline{\nu}\gamma_{\mu}(1-\nu_5)\nu$$

can be measured in the leading decay $K^+ o \pi^0 e^+ \nu$

$$B_{SD}(K^{+} \to \pi^{+} \nu \overline{\nu}) = \frac{\kappa_{+} \alpha^{2} B(K_{e3})}{2\pi^{2} \sin^{4} \theta_{W} |V_{us}|^{2}} \sum_{l} |X_{t} \lambda_{t} + X_{c} \lambda_{c}|^{2} = 8.9 \times 10^{-11} A^{4} \left[(\rho_{0} - \overline{\rho})^{2} + \overline{\eta}^{2} \right]$$







SUSY Contribution

Project X sensitive to 1000 SM events.
BSM rates 10x SM rates.



$K_L \to \pi^0 \nu \bar{\nu}$

Buras: $K_L \to \pi^0 \nu \overline{\nu}$, $K^+ \to \pi^+ \nu \overline{\nu}$, ratio x_d/x_s of $B_d^0 - B_d^0$ to $B_s^0 - \overline{B_s^0}$ mixing and class of asymmetries in neutral B decays cleanest observables, being essentially free from hadronic uncertainties.

- •Purely CP-Violating (Littenberg, 1989)
- •Totally dominated from t-quark
- •Computed to NLO in QCD (Buchalla, Buras, 1999)
- •No long distance contribution SM ~3 × 10⁻¹¹

Backgrounds:

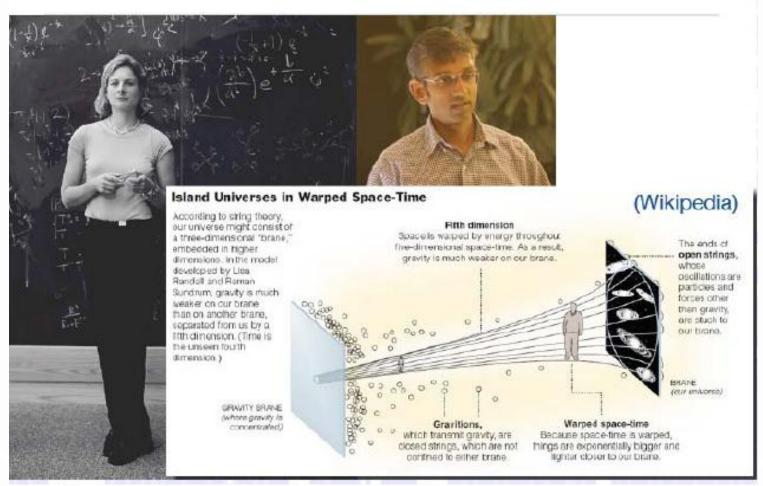
$$K_L
ightarrow 2\pi^0$$
 , $\pi^0 e^+ e^-$, $\pi^0 \gamma \gamma$

Difficult mode to measure



Rates of $K^0 \to \pi^0 \nu \overline{\nu}$ and $K^+ \to \pi^+ \nu \overline{\nu}$ sensitive to other NP models also

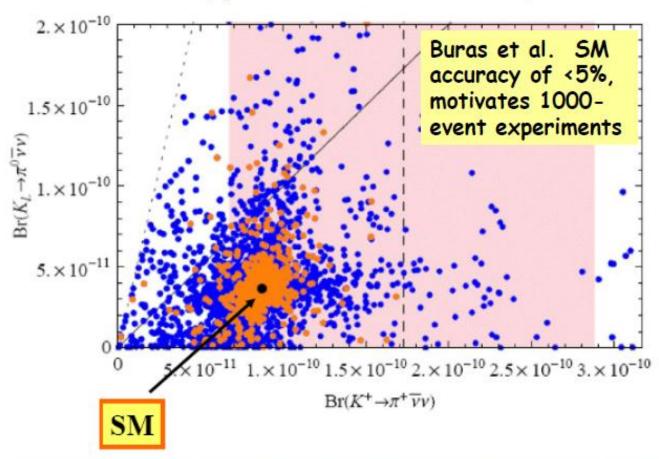
The Randall-Sundrum (RS) idea







(Up to Factor 3 and 2 Enhancements)



Effect of Warped Extra Dimension Models on Branching Fractions

SEL 2011, January 2011

R. Tschirhart - Fermilab



$K_L \to \pi \pi \ell^+ \ell^-$

$$d\Gamma = \frac{G_F^2}{2^{12}\pi^6 M_K^5} \sin^2\Theta_C X \sigma_{\pi} \left[1 - \frac{4m_l^2}{s_l} \right]^2 I(s_{\pi}, s_l, \Theta_{\pi}, \Theta_l, \Phi) ds_{\pi} ds_l d \cos\Theta_{\pi} d \cos\Theta_l d\Phi ,$$

$$+I_4\sin 2\Theta_l\cos \Phi +I_5\sin \Theta_l\cos \Phi$$

 $I = I_1 + I_2 \cos 2\Theta_l + I_3 \sin^2\Theta_l \cos 2\Phi$

$$+I_6\cos\Theta_l+I_7\sin\Theta_l\sin\Phi$$

$$+I_8\sin 2\Theta_l\sin \Phi+I_9\sin^2\Theta_l\sin 2\Phi$$
,

$$F_1 = Xf + \sigma_{\pi} s \cos\Theta_{\pi} g$$
,

$$F_2 = \sigma_{\pi}(s_{\pi}s_l)^{1/2}g$$
,

$$F_3 = \sigma_{\pi} X (s_{\pi} s_l)^{1/2} \frac{h}{M_K^2}$$
,

$$\begin{split} I_1 &= \frac{1}{4} [\{|F_1|^2 + \frac{3}{2} (|F_2|^2 + |F_3|^2) \sin^2\Theta_{\pi} \} \\ &+ (F_{1,2,3} \to \widetilde{F}_{1,2,3})] \ , \\ I_2 &= -\frac{1}{4} [\{|F_1|^2 - \frac{1}{2} (|F_2|^2 + |F_3|^2) \sin^2\Theta_{\pi} \} \\ &+ (F_{1,2,3} \to \widetilde{F}_{1,2,3})] \ , \\ I_3 &= -\frac{1}{4} [\{|F_2|^2 - |F_3|^2 \} + (F_{1,2,3} \to \widetilde{F}_{1,2,3})] \ , \\ I_4 &= \frac{1}{2} \mathrm{Re}(F_1^* F_2) \sin\Theta_{\pi} + (F_{1,2,3} \to \widetilde{F}_{1,2,3}) \ , \\ I_5 &= -\{ \mathrm{Re}(F_1^* F_3) \sin\Theta_{\pi} - (F_{1,2,3} \to \widetilde{F}_{1,2,3}) \} \ , \\ I_6 &= -\{ \mathrm{Re}(F_2^* F_3) \sin^2\Theta_{\pi} - (F_{1,2,3} \to \widetilde{F}_{1,2,3}) \} \ , \\ I_7 &= -\{ \mathrm{Im}(F_1^* F_2) \sin\Theta_{\pi} - (F_{1,2,3} \to \widetilde{F}_{1,2,3}) \} \ , \\ I_8 &= \frac{1}{2} \mathrm{Im}(F_1^* F_3) \sin\Theta_{\pi} + (F_{1,2,3} \to \widetilde{F}_{1,2,3}) \ , \\ I_9 &= -\frac{1}{2} [\mathrm{Im}(F_2^* F_3) \sin^2\Theta_{\pi} + (F_{1,2,3} \to \widetilde{F}_{1,2,3})] \ . \end{split}$$

Visiting of the state of the st

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$$\mathcal{A} = \frac{\int_{0}^{\pi/2} \frac{d\Gamma}{d\Phi} d\Phi - \int_{\pi/2}^{\pi} \frac{d\Gamma}{d\Phi} d\Phi}{\int_{0}^{\pi/2} \frac{d\Gamma}{d\Phi} d\Phi + \int_{\pi/2}^{\pi} \frac{d\Gamma}{d\Phi} d\Phi} / Strong \ phase}$$

$$= 15\% \sin[\Phi_{+}] + \delta_{0}(m_{K}^{2}) - \overline{\delta_{1}}]$$

$$\approx 14\%.$$

$$Weak \ phase$$

$$K_L \to \pi^+(p_+)\pi^-(p_-)\ell^+(k_+)\ell^-(k_-)$$

$$\begin{bmatrix}
\mathbf{p}_{\pm} \xrightarrow{CP} & \mathbf{p}_{\mp} & \cos \Theta_{\pi} & \rightarrow -\cos \Theta_{\pi} & \sin \Theta_{\pi} & \rightarrow \sin \Theta_{\pi} \\
\mathbf{U} \mathbf{n} \mathbf{d} \mathbf{e} \mathbf{r} & \mathbf{C} \mathbf{P} \\
\mathbf{k}_{\pm} \xrightarrow{CP} & \mathbf{k}_{\mp} & \cos \Theta_{\ell} & \rightarrow -\cos \Theta_{\ell} & \sin \Theta_{\ell} & \rightarrow \sin \Theta_{\ell} \\
& \cos \Phi & \rightarrow \cos \Phi & \sin \Phi & \rightarrow -\sin \Phi
\end{bmatrix}$$

Signal of T-reversal violation

Several papers supporting and several other disputing signal is genuine T-violation



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CPT is introduced through the Hamiltonian

$$\mathcal{H} = E \begin{pmatrix} \cos \theta & \sin \theta e^{-i\phi} \\ \sin \theta e^{i\phi} & -\cos \theta \end{pmatrix} - iD\mathcal{I}$$
. CPT restored if $\theta = \frac{\pi}{2}$

A complete calculation without CPT in mixing is underway.

CPT violation should be studied in K since large numbers of K mesons will be produced.



Conclusion

- After more than 60 years K meson continues to be produced in lab and is still a valuable source for understanding new physics.
- Will continue to be a studied at least until a 5σ signal is observed in $K_L\to\pi^0\nu\bar\nu$.



$$\begin{split} K_1 &= (K^0 + \overline{K}^{\ 0})/\sqrt{2} \ , \\ K_2 &= (K^0 - \overline{K}^{\ 0})/\sqrt{2} \ , \\ K_S &= (K_1 + \widetilde{\epsilon} K_2)/\sqrt{1 + |\widetilde{\epsilon}|^2} \ , \\ K_L &= (K_2 + \widetilde{\epsilon} K_1)/\sqrt{1 + |\widetilde{\epsilon}|^2} \ . \end{split}$$

