

# Hydrodynamics and fluctuations in relativistic heavy-ion collisions

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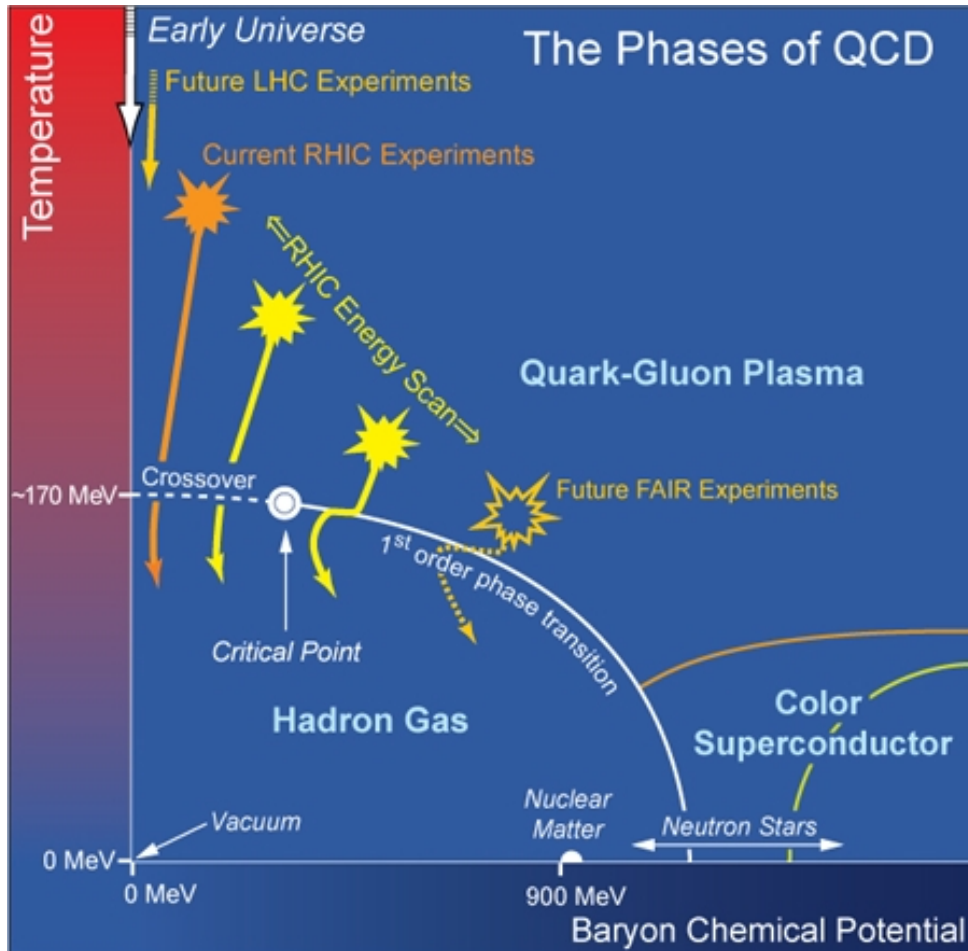
# The relativistic heavy-ion programme



RHIC at Brookhaven :  
Two beams of atomic Au nuclei, accelerated at energies up to 100 GeV per nucleon (since 2000)

LHC at CERN:  
Two beams of atomic Pb nuclei will be accelerated at energies up to 2.7 TeV per nucleon (2010?)

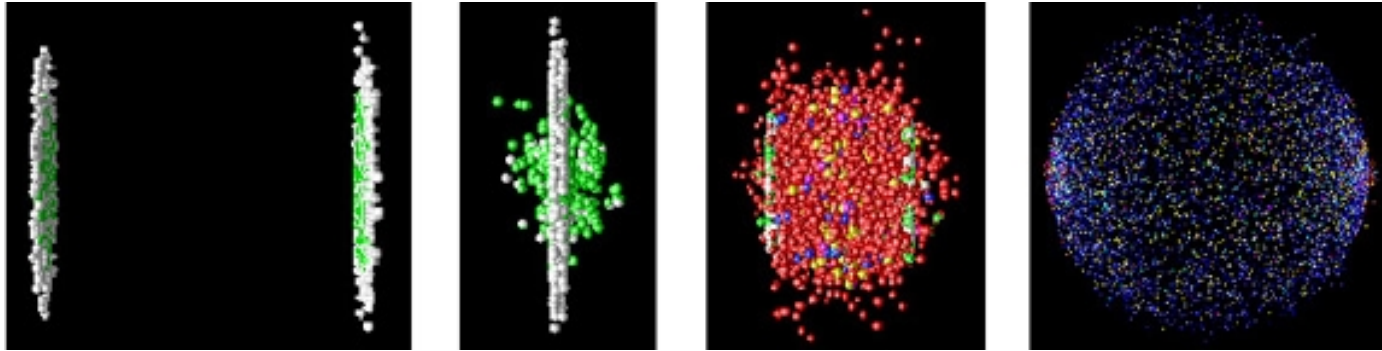
# The phase diagram of strong interactions



Can we learn anything about hot gauge theories by smashing heavy ions together?

Do heavy-ion collisions have anything to do with temperature and thermodynamics?

# What RHIC has taught us



During the expansion, the matter behaves collectively like a fluid.

This fluid has the smallest [viscosity/entropy](#) ratio ever seen: typically 2x the absolute lower bound postulated using gauge/gravity duality:  $\eta/s = \hbar/4\pi k_B$ .

[Kovtun Son Starinets hep-th/0405231](#)

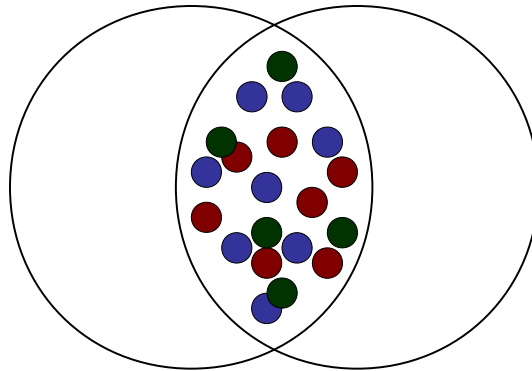
# Outline

- A close look at the pattern of emitted particles: elliptic flow and « higher harmonics »
- A simple, universal prediction from hydrodynamics
- Comparing with experimental data
- Taking into account fluctuations
- Conclusion

Gombeaud, JYO, [arXiv:0907.4664](https://arxiv.org/abs/0907.4664)

# A primer on nucleus-nucleus collisions

A typical Au-Au collision viewed in the **transverse** plane, perpendicular to beam axis



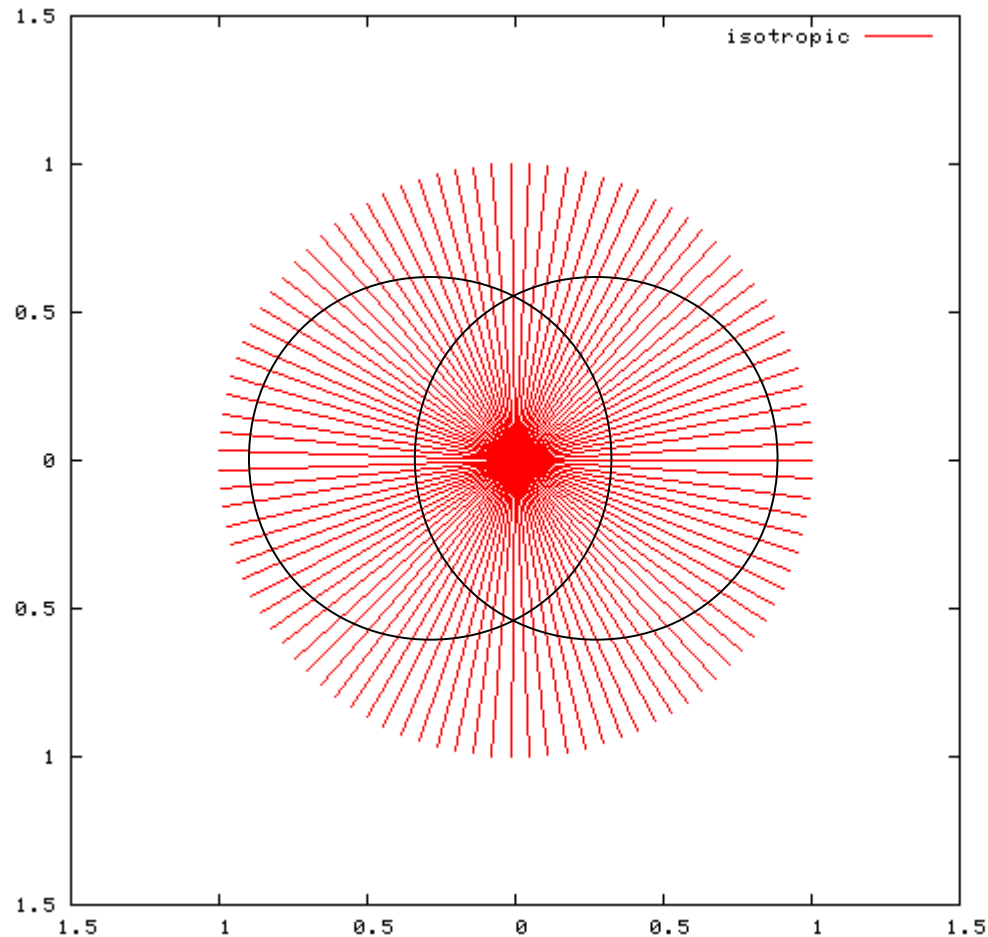
Collisions between partons/nucleons occur in the **overlap area** between the two nuclei: this is where matter is originally created

The non-overlapping parts don't interact (we call them « spectator » nucleons)

By measuring the number of spectators, or the number of participants, one estimates the **centrality** (impact parameter) of a given collision.

# What are the directions of created particles?

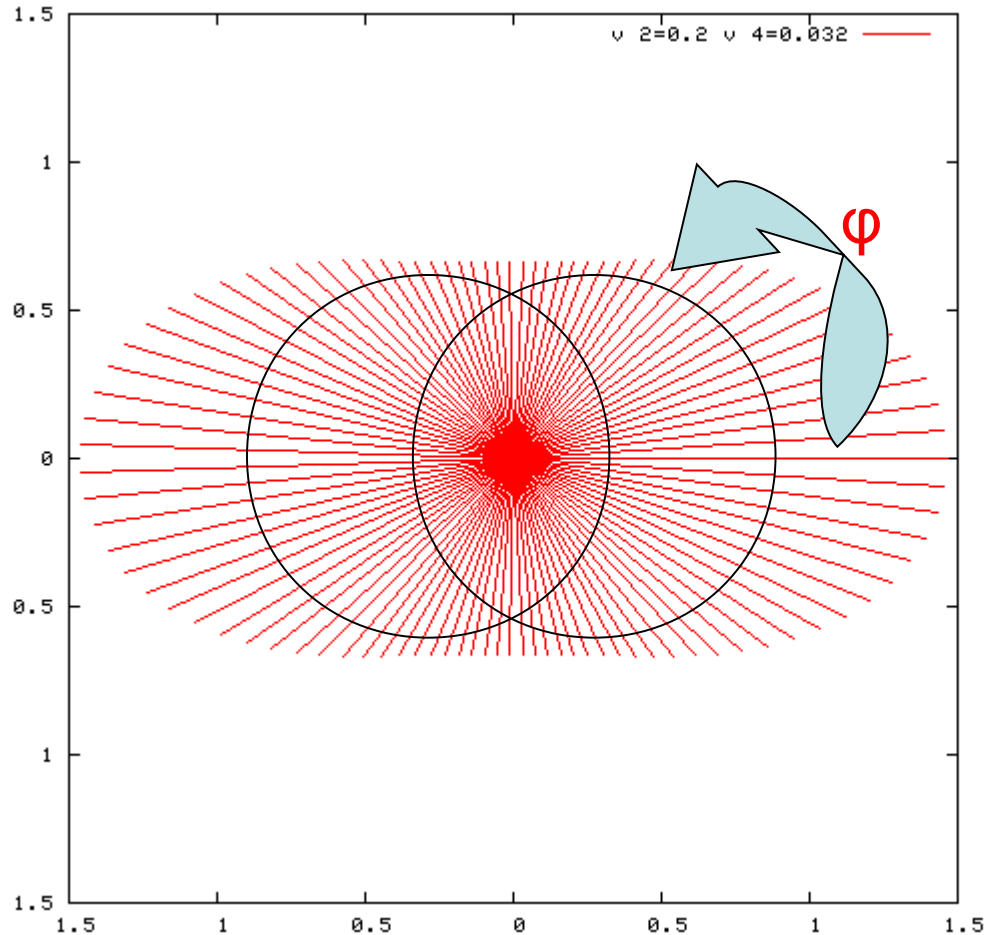
Random  
parton-parton  
collisions  
occurring on  
scales  $\ll$   
nuclear radius.  
No preferred  
direction in the  
production  
process.  
Isotropic  
azimuthal  
distribution





# What we see is :

Bar length  
= number  
of particles  
in the  
direction  
= Azimuthal  
( $\varphi$ )  
distribution  
plotted in  
polar  
coordinates



We call this  
**elliptic flow**.  
We think it  
is created  
by pressure  
gradients in  
the overlap  
area

(for pions with transverse momentum  $\sim 2$  GeV/c)



# Anisotropic flow

Fourier series expansion of the azimuthal distribution:

Using the  $\varphi \rightarrow -\varphi$  and  $\varphi \rightarrow \varphi + \pi$  symmetries of overlap area:

$$dN/d\varphi = 1 + 2v_2 \cos(2\varphi) + 2v_4 \cos(4\varphi) + \dots$$

$v_2 = \langle \cos(2\varphi) \rangle$  ( $\langle \dots \rangle$  means average value) is elliptic flow

$v_4 = \langle \cos(4\varphi) \rangle$  is a (much smaller) « higher harmonic »

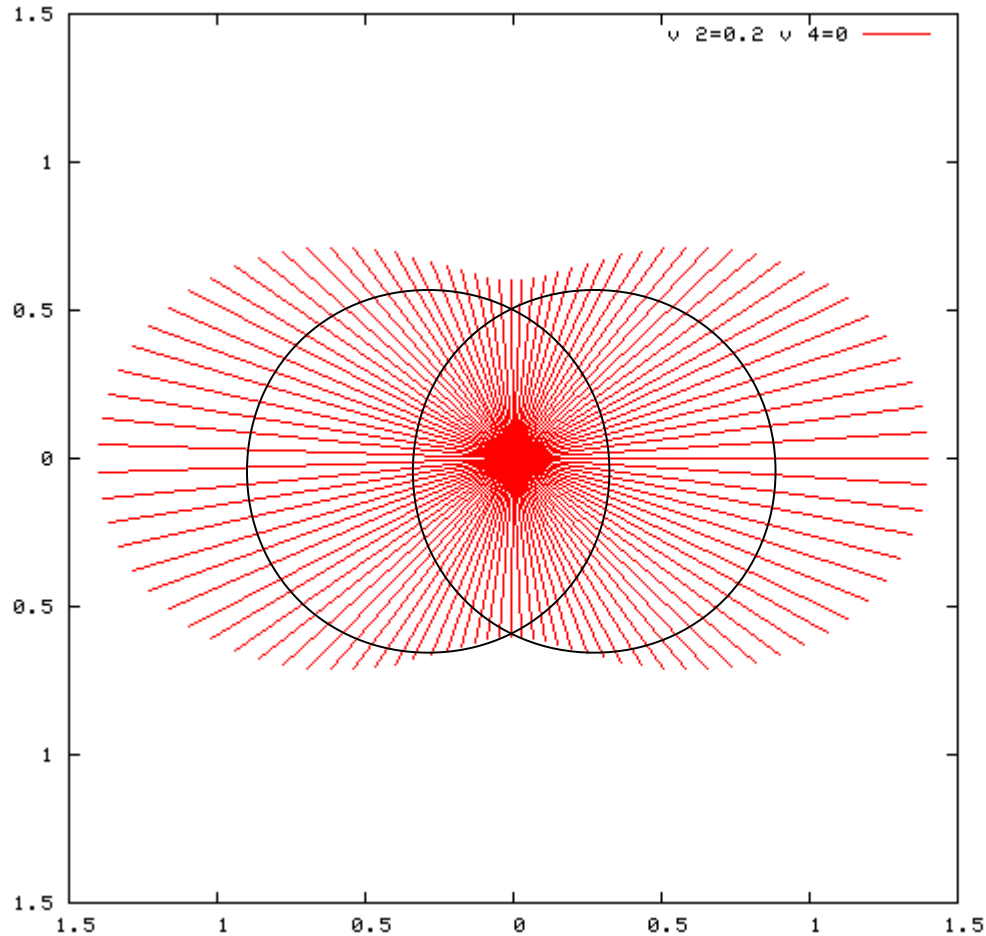
higher harmonics  $v_6$ , etc are 0 within experimental errors.

This talk really is about  $v_4$

# Azimuthal distribution without $v_4$

A small effect:  
Average value 0.3%,  
maximum value 3%

Should we care?



The beauty is in the details!

# A primer on hydrodynamics

- Ideal gas (weakly-coupled particles) in **global** thermal equilibrium. The phase-space distribution is (Boltzmann)

$$dN/d^3pd^3x = \exp(-E/T)$$

**Isotropic!**

- A fluid moving with velocity  $\mathbf{v}$  is in **(local)** thermal equilibrium in its rest frame:

$$dN/d^3pd^3x = \exp(-(E-\mathbf{p}\cdot\mathbf{v})/T)$$

**Not isotropic: Momenta parallel to  $\mathbf{v}$  preferred**

- At RHIC, the fluid velocity depends on  $\varphi$ :  
typically  $v(\varphi) = v_0 + 2\varepsilon \cos(2\varphi)$

# The simplicity of $v_4$

- Within the approximation that particle momentum  $\mathbf{p}$  and fluid velocity  $\mathbf{v}$  are parallel (good for large momenta)

$$dN/d\varphi = \exp(2\varepsilon p \cos(2\varphi)/T)$$

- Expanding to order  $\varepsilon$ , the  $\cos(2\varphi)$  term is

$$v_2 = \varepsilon p/T$$

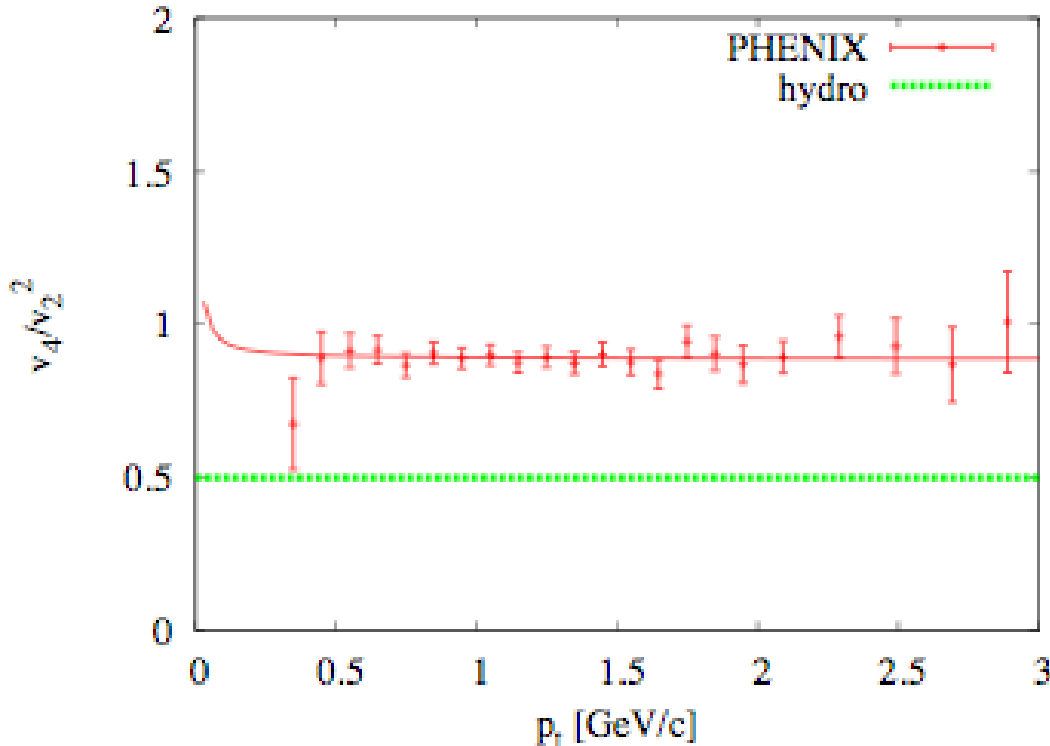
- Expanding to order  $\varepsilon^2$ , the  $\cos(4\varphi)$  term is

$$v_4 = \frac{1}{2} (v_2)^2$$

**Hydrodynamics has a universal prediction for  $v_4/(v_2)^2$  !**

Should be independent of equation of state, initial conditions, centrality, particle momentum and rapidity, particle type

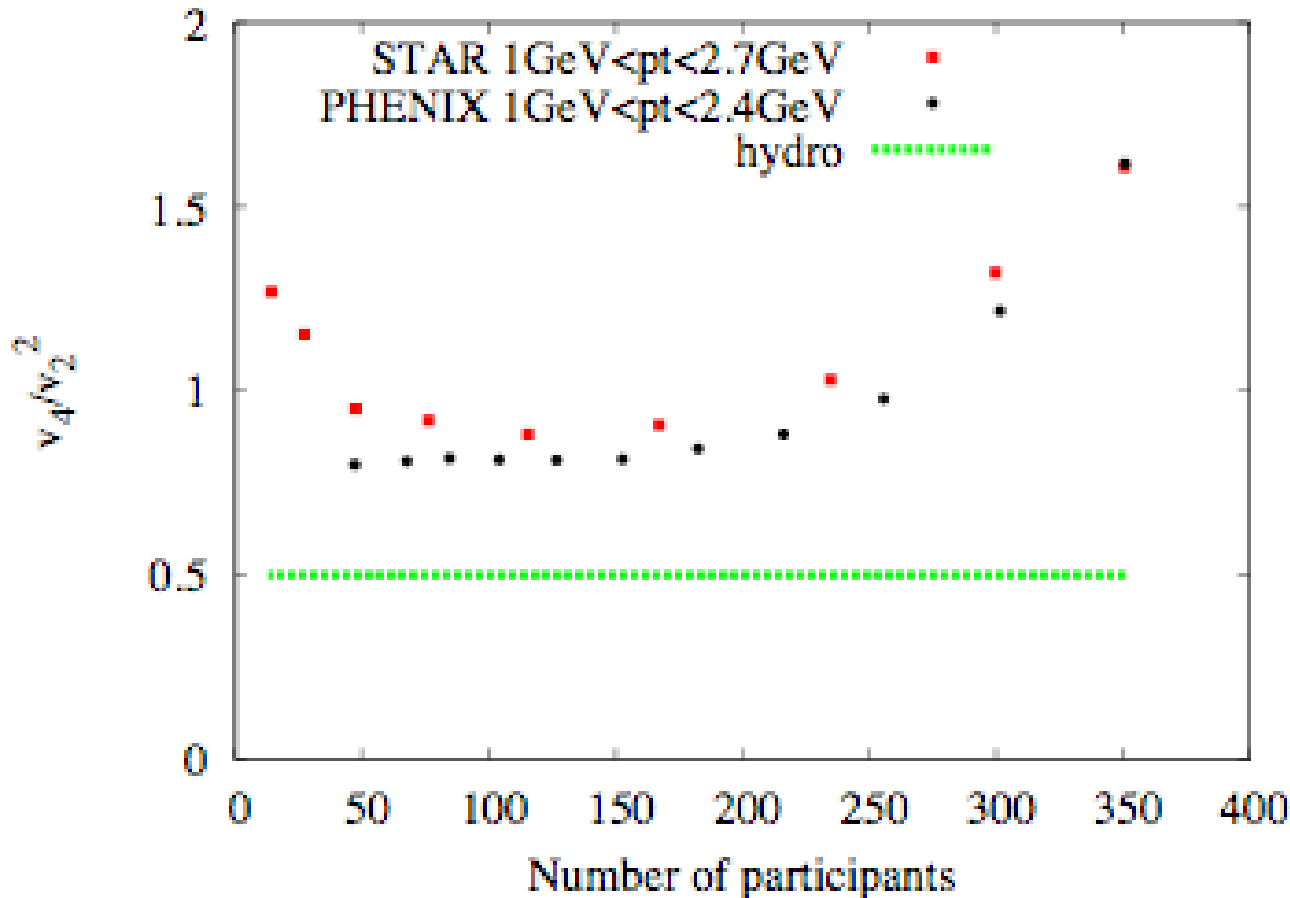
# PHENIX results for $v_4$



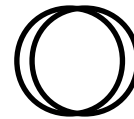
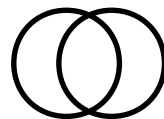
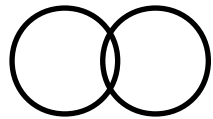
PHENIX data for charged pions  
Au-Au collisions at 100+100 GeV  
20-60% most central

The ratio is independent of  $p_T$ , as predicted by hydro.  
But... the value is significantly larger than 0.5

# More data : centrality dependence



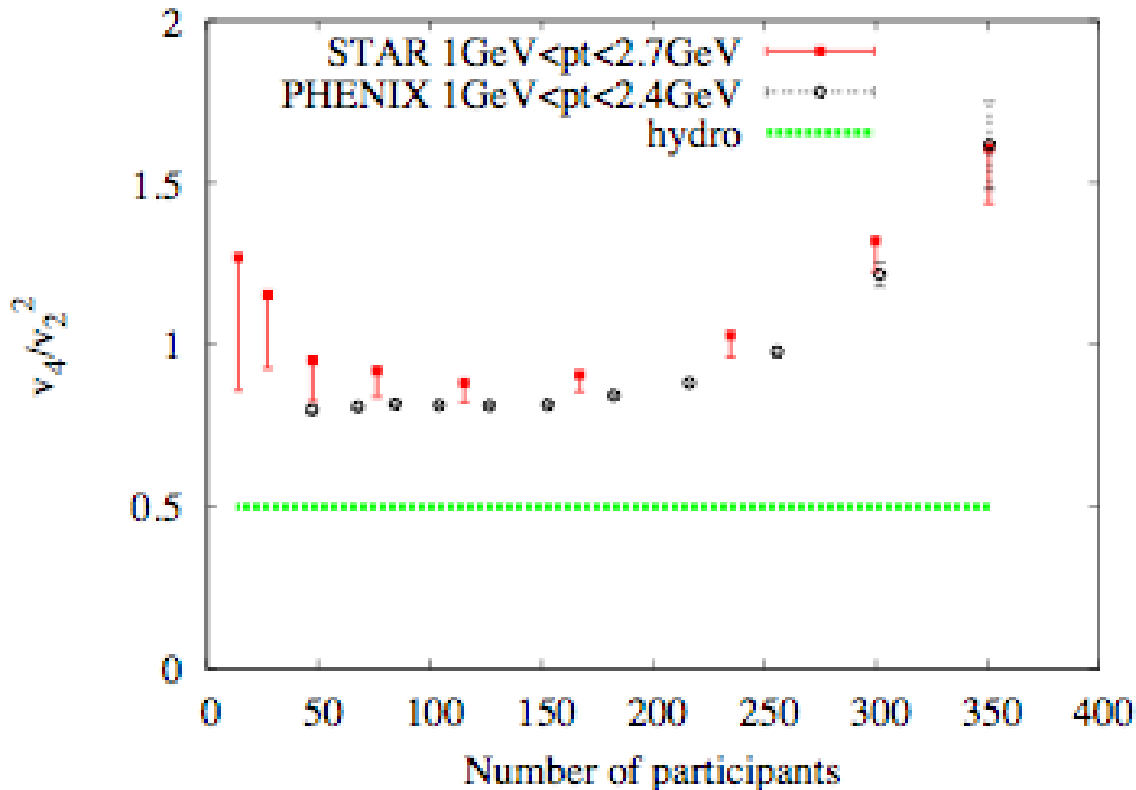
Au-Au collision  
100+100 GeV  
per nucleon



Data > hydro

Small discrepancy between STAR and PHENIX data

# Estimating experimental errors



$v_2$  and  $v_4$  are not measured directly but inferred from azimuthal correlations (more later on this). There are many sources of correlations (jets, resonance decays,...): this is the « nonflow » error which we can estimate (order of magnitude only)

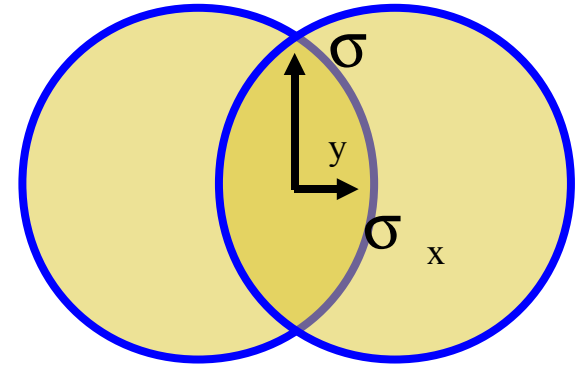
Difference between STAR and PHENIX data compatible with non-flow error

How do we understand the discrepancy with hydrodynamics??



# Eccentricity scaling

We understand elliptic flow as the consequence of the almond shape of the overlap area

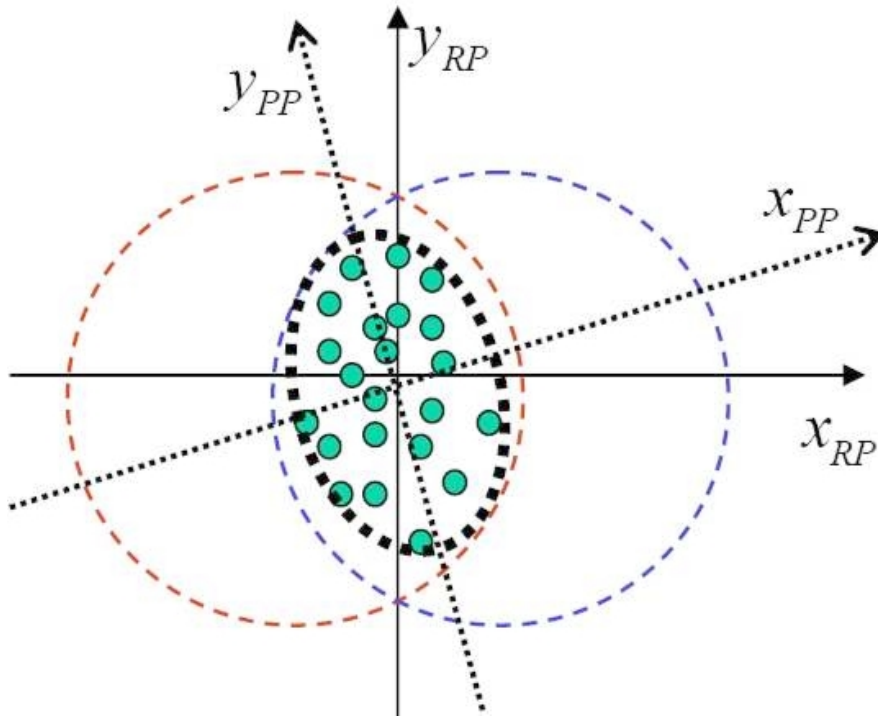


It is therefore natural to expect that  $v_2$  scales like the **eccentricity**  $\varepsilon$  of the **initial density profile**, defined as :

$$\varepsilon = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$

(this is confirmed by numerical hydro calculations)

# Eccentricity fluctuations



Depending on where the participant nucleons are located within the nucleus at the time of the collision, the actual shape of the overlap area may vary: the orientation and eccentricity of the ellipse defined by participants fluctuates.

Assuming that  $v_2$  scales like the eccentricity, eccentricity fluctuations translate into  $v_2$  (elliptic flow) fluctuations

# Why fluctuations change $v_4/(v_2)^2$

The reference directions  $x$  and  $y$  are not known experimentally:  
thus  $v_2 = \langle \cos(2\varphi) \rangle$  and  $v_4 = \langle \cos(4\varphi) \rangle$  are not measured directly

$v_2$  from 2-particle correlations:  $\langle \cos(2\varphi_1 - 2\varphi_2) \rangle = \langle v_2 \rangle^2$

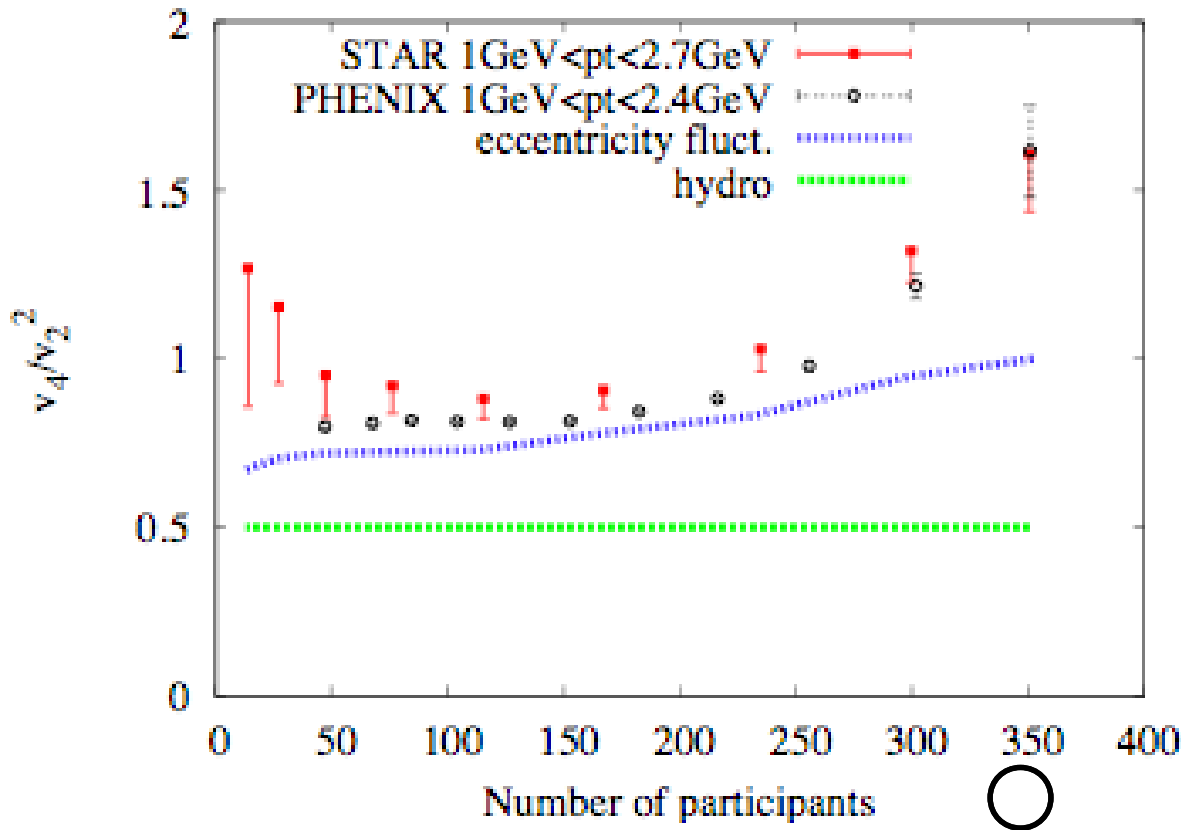
$v_4$  from 3-particle correlations:  $\langle \cos(4\varphi_1 - 2\varphi_2 - 2\varphi_3) \rangle = \langle v_4 \rangle (v_2)^2$

If  $v_2$  and  $v_4$  fluctuate, the measured  $v_4/(v_2)^2$  is really

$\langle v_4 \rangle (v_2)^2 / \langle (v_2)^2 \rangle^2$ . Inserting the prediction from hydrodynamics,

$$[v_4/(v_2)^2]_{\text{exp}} = 1/2 \langle (v_2)^4 \rangle / \langle (v_2)^2 \rangle^2$$

# Data versus eccentricity fluctuations



Eccentricity fluctuations can be modelled using a **Monte-Carlo** program provided by the PHOBOS collaboration: **Throw the dice** for the positions of nucleons, with probability given by the nuclear density:  
**No free parameter!**

Fluctuations explain most of the discrepancy between data and hydro, except for central collisions which suggest  $\langle (v_2)^4 \rangle / \langle (v_2)^2 \rangle^2 = 3$   
 By symmetry,  $v_2 = 0$  for central collisions, *except for fluctuations!*

# Conclusions

- The fourth harmonic,  $v_4$ , of the azimuthal distribution gives a further, independent indication that the matter produced at RHIC expands like a **relativistic fluid**
- We are colliding nuclei=complex quantum systems. We clearly see in the data large **fluctuations** which originate from the **wavefunction** of the colliding nuclei
- The (by now standard) model of eccentricity fluctuations fails for central collisions. We need a better understanding of fluctuations.

# Is $v_2$ linear in momentum?

Hydro by Huovinen et al.  
hydro tuned to fit central  
spectra data.

