

# Discriminating Higgs production mechanisms using jet energy profiles

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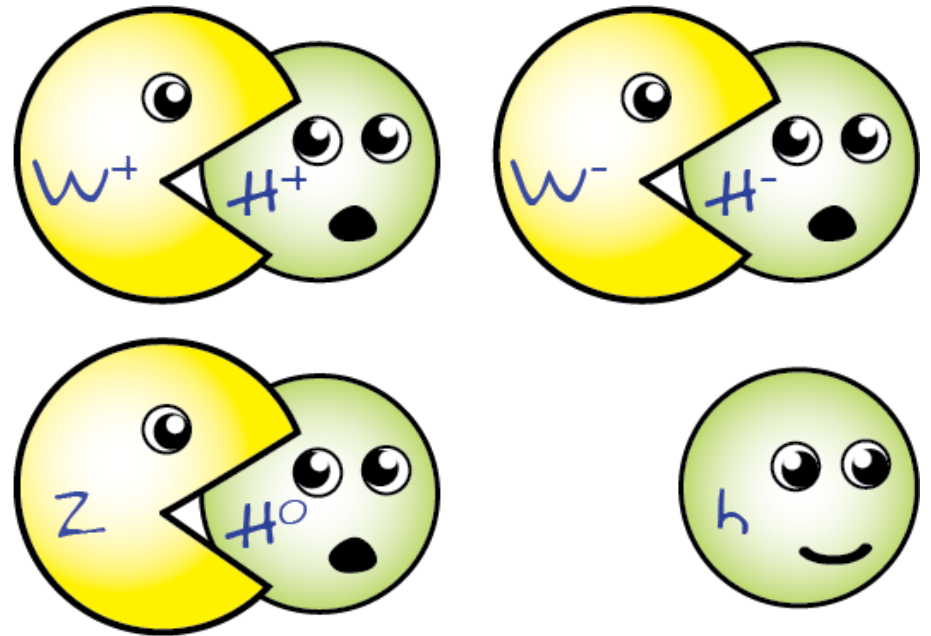
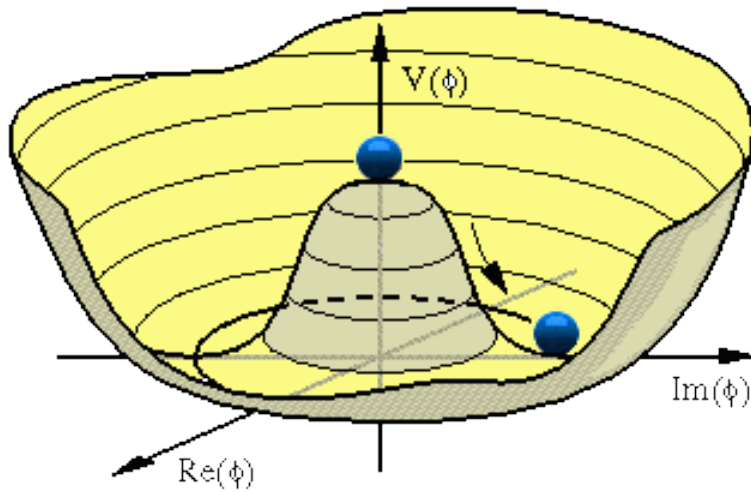


# Outline

- The Higgs boson
- Higgs observables and Higgs coupling extraction
- Breaking degeneracies by separating the production mechanisms
- Kinematic discriminants
- Jet Energy Profiles
- Results

**What is the origin of mass  
in the Standard Model?**

# Higgs mechanism



**Remnant of EWSB is the Higgs boson**

# Status so far of what the LHC has seen

- We have seen a spin-0, 125 GeV particle that has the approximate properties of the Standard Model Higgs boson.
- No other new (high mass) particles or resonances have been seen yet.

**We have discovered A Higgs boson, is  
it THE Higgs boson of the SM?**

# Key predictions from the SM

- There is a SINGLE Higgs field that acquires a vacuum expectation value
- Excitation of this Higgs field is the Higgs boson
- Couplings of the Higgs boson to all SM particles must be in proportion to their masses

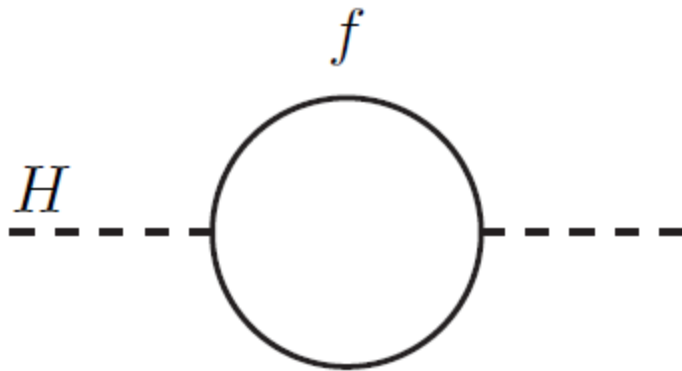
# Why go beyond the Standard Model?

## Some of the big questions

- What are dark matter/dark energy?
- What explains masses and mixings of fermions?
- What is the origin of the small neutrino masses?
- What explains matter/anti-matter asymmetry?
- What is the mechanism that causes inflation?
- Quantum gravity?
- ....
- Hierarchy problem

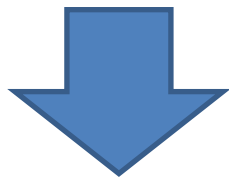


# The Hierarchy/Naturalness Problem



- Cut-off  $\sim$  TeV scale without large fine tuning

$$m_H^2 = m_0^2 - \frac{|\lambda_f|^2}{8\pi^2} \Lambda_{UV}^2 + \dots$$



125 GeV



$M_{\text{pl}} = 10^{19}$  GeV

**Generic Prediction: New physics at  
the TeV scale!**

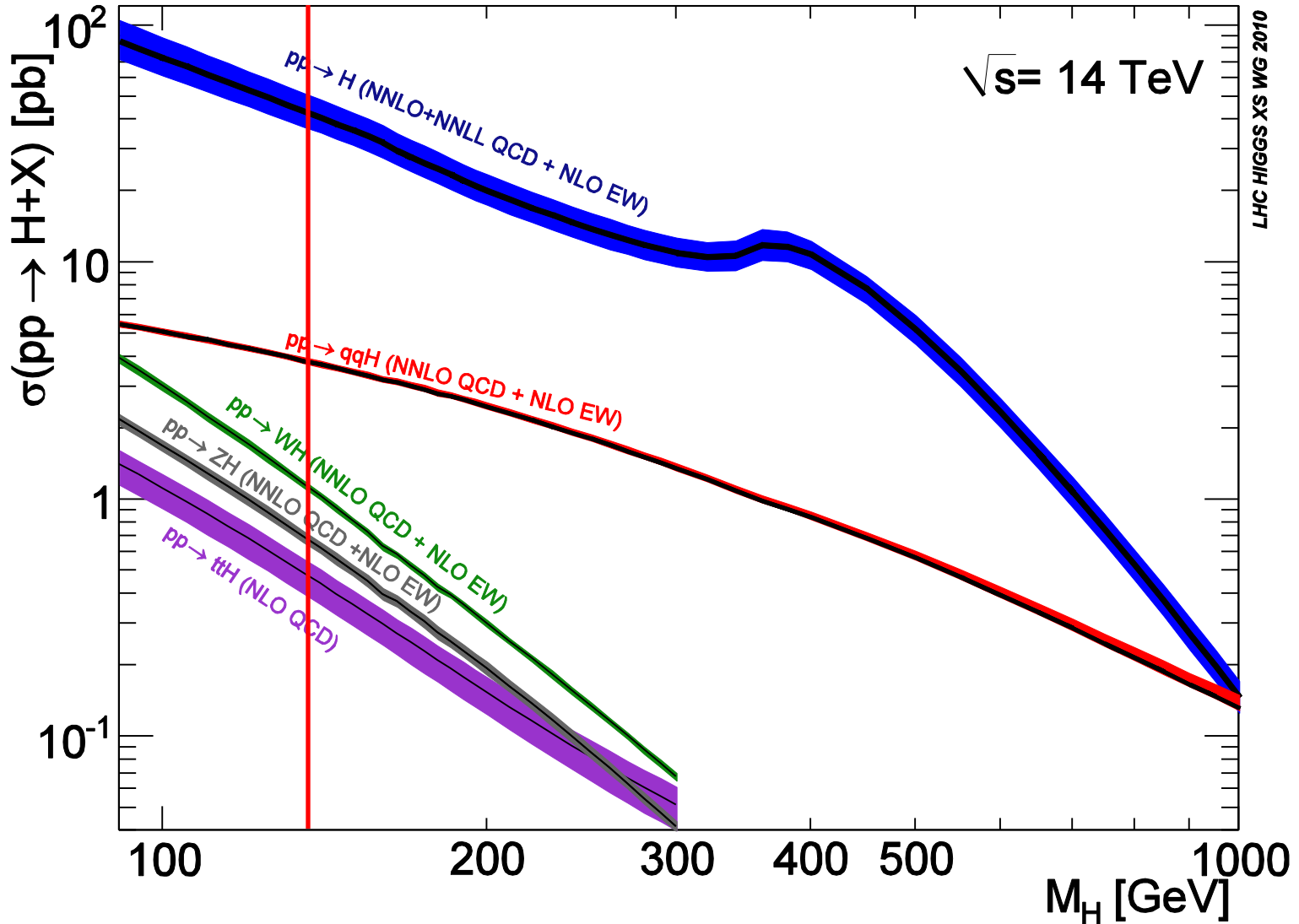
# Generic predictions of solutions to the hierarchy problem

- 1. Deviations in Higgs boson interaction as compared to the SM.**
2. New resonances of the electroweak  $W$  and  $Z$  gauge bosons.
3. Extra Higgs multiplets.
4. Partners to SM particles from models such as supersymmetry and extra dimensions.

# How much data has the LHC gathered?

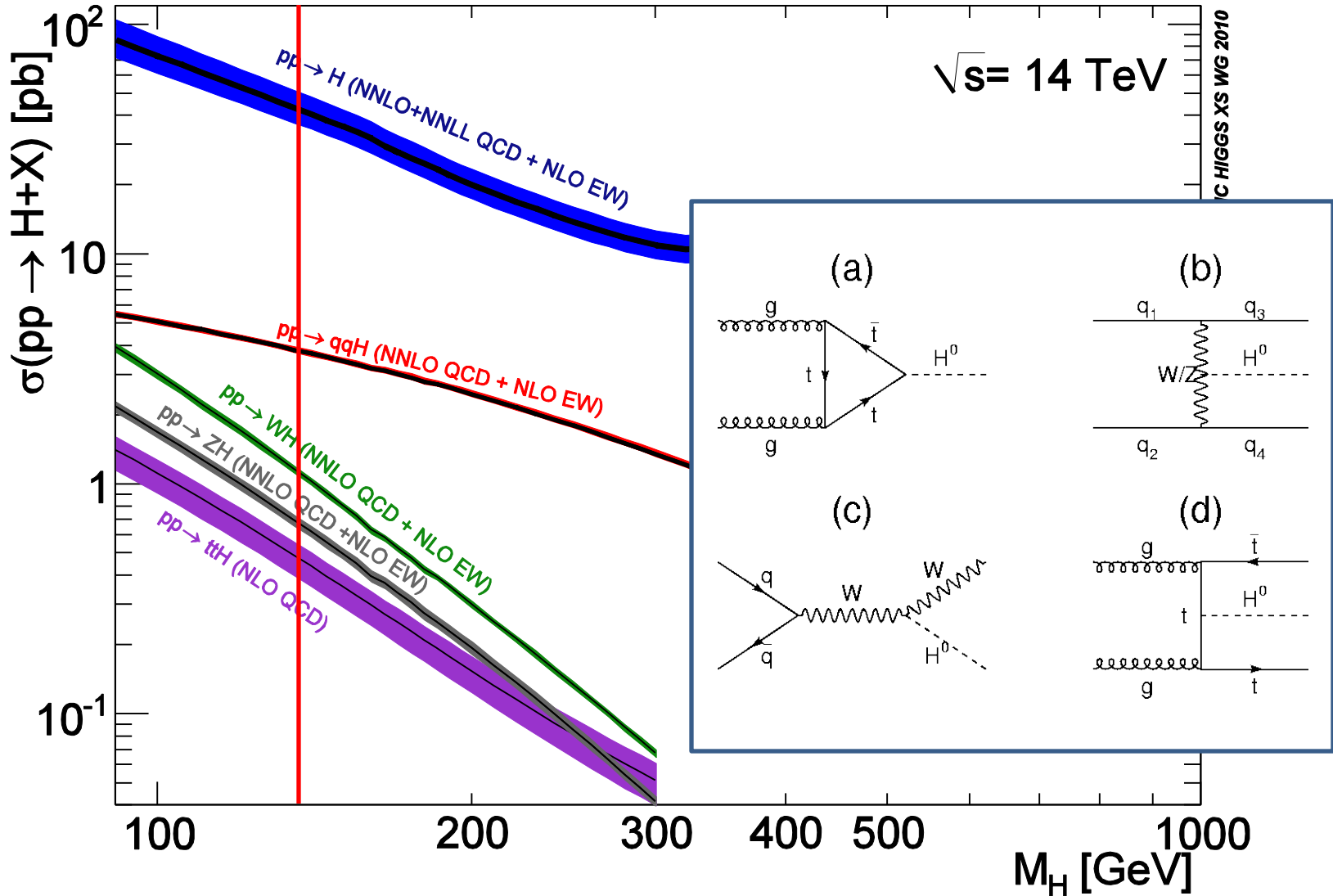
- LHC has just completed a low energy run at 7 and 8 TeV center of mass energy
- Data collected so far:
  - At 7 TeV, we have  $5 \text{ fb}^{-1}$  of data
  - At 8 TeV, we have  $20 \text{ fb}^{-1}$  of data
- LHC has restarted collisions at 13 TeV center of mass energy. We expect to collect up to  $300 \text{ fb}^{-1}$  of data in the next few years.

# Higgs production rates at LHC



Dominated by GF and VBF

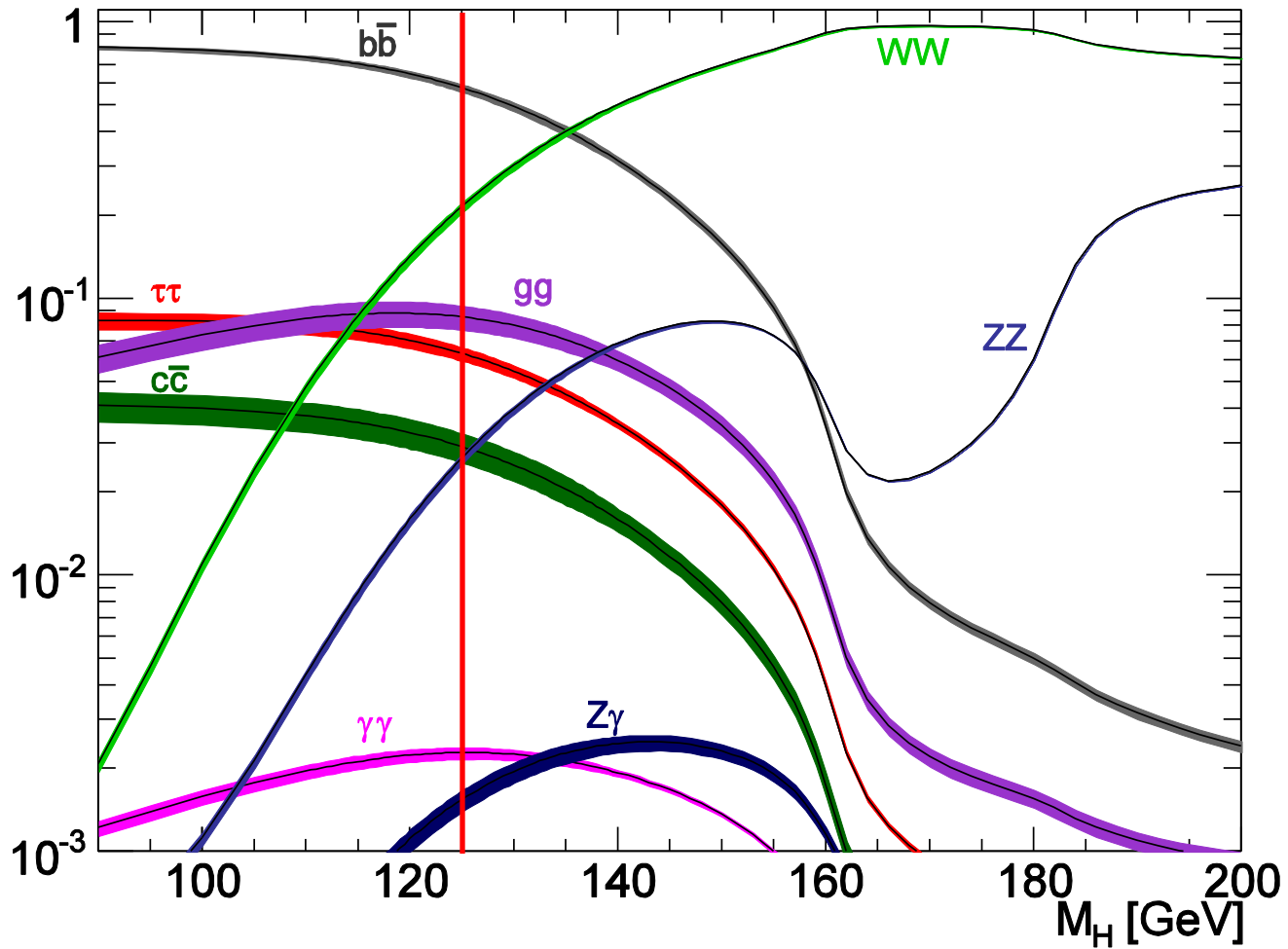
# Higgs production rates at LHC



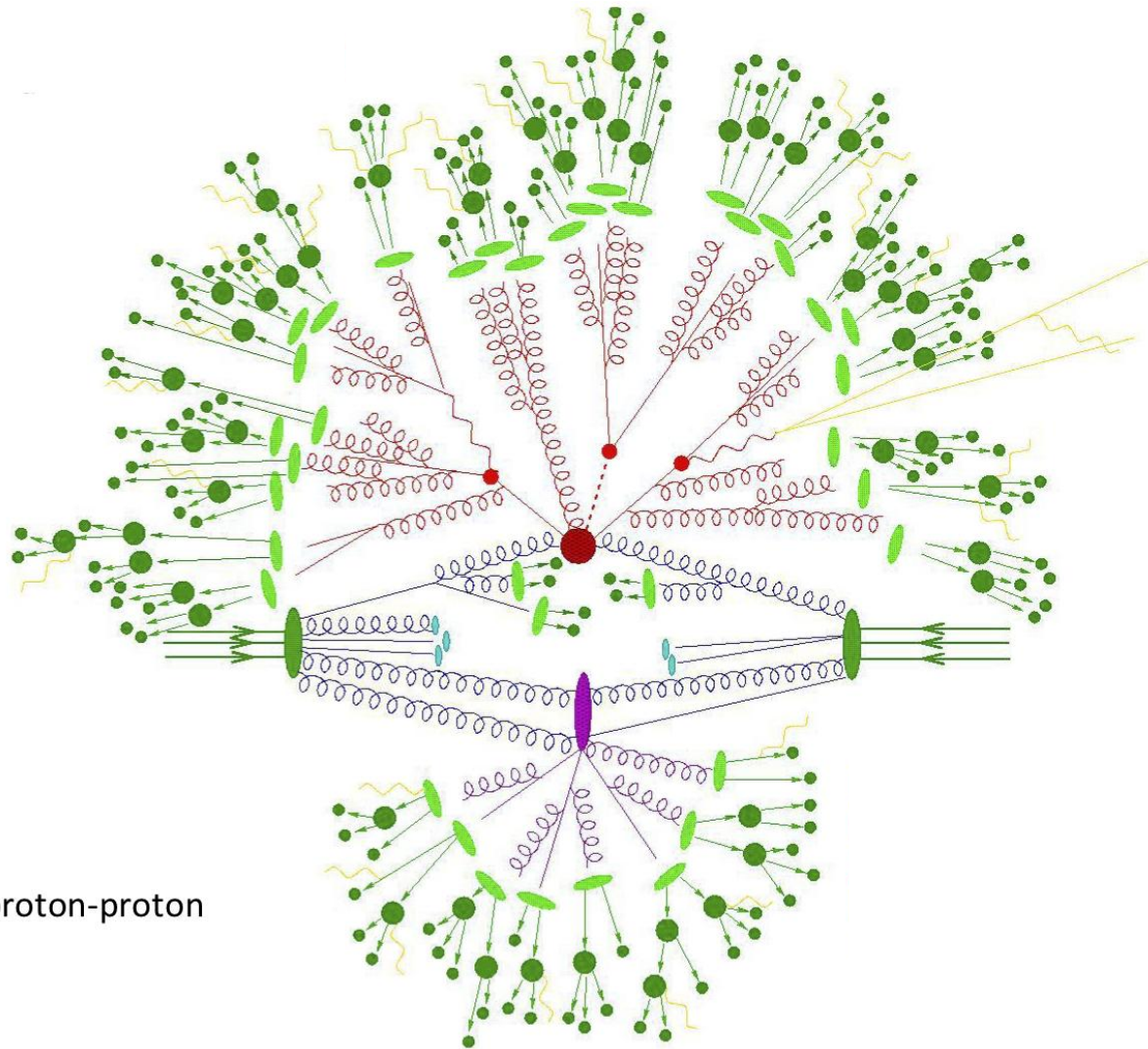
IC HIGGS XS WG 2010

Dominated by GF and VBF

# Higgs branching fractions



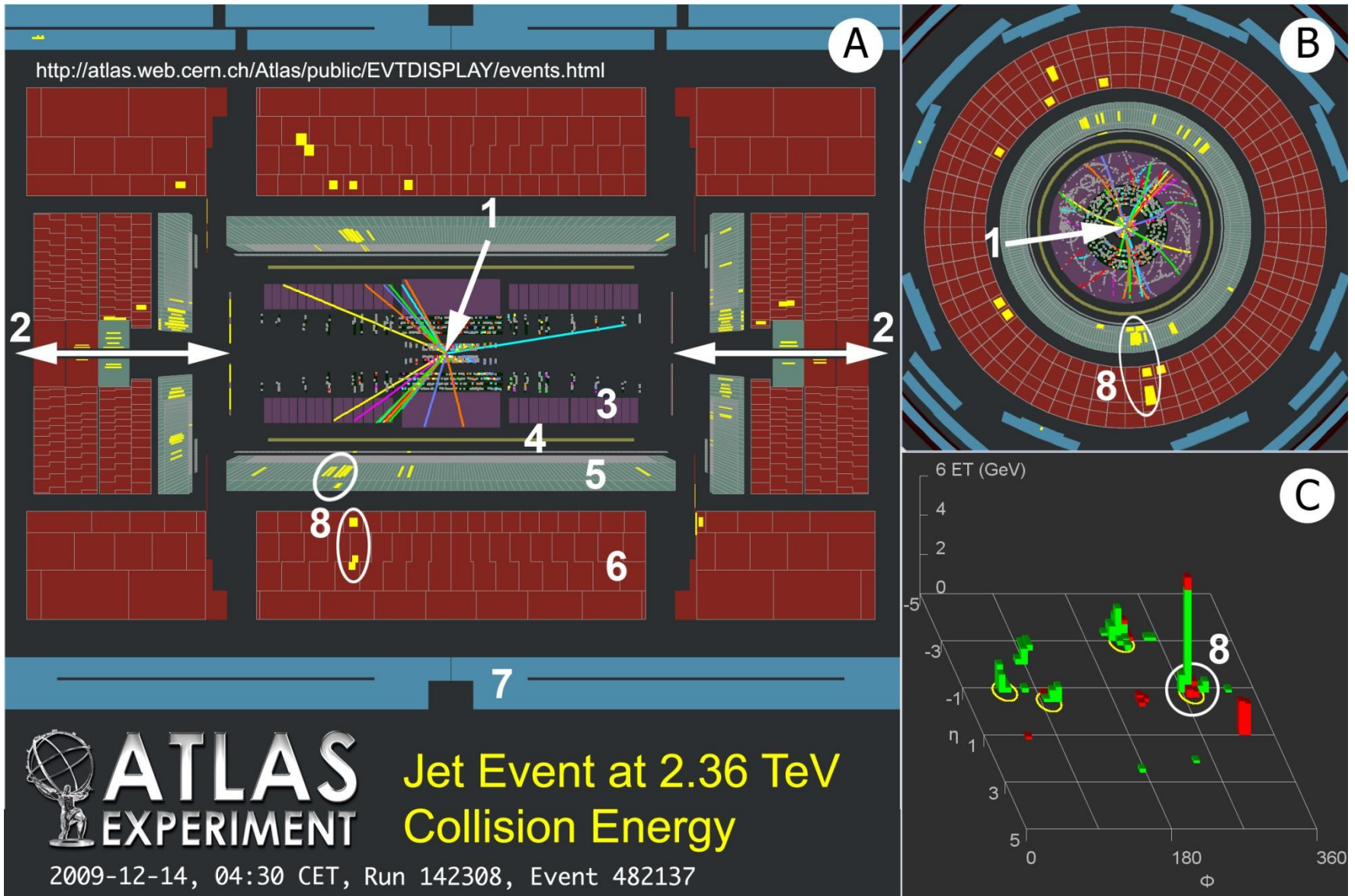
# What do events at the LHC look like?



Typical proton-proton collision



# What do events at the LHC look like?

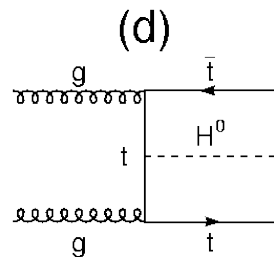
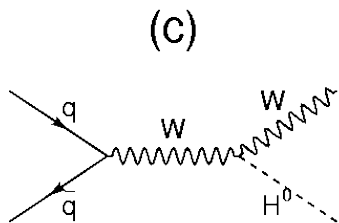
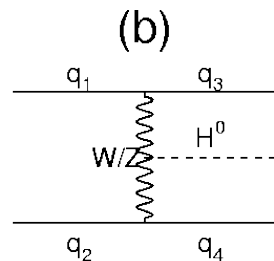
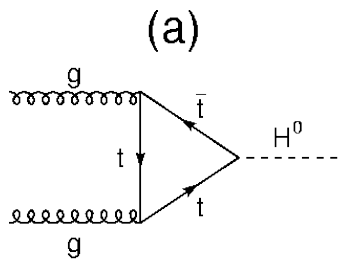
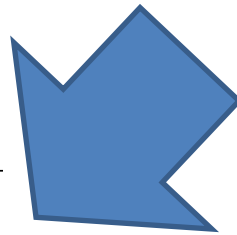


# Expected Rates for the Higgs boson

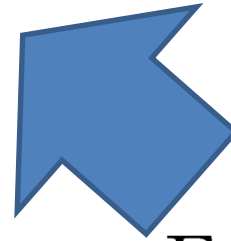
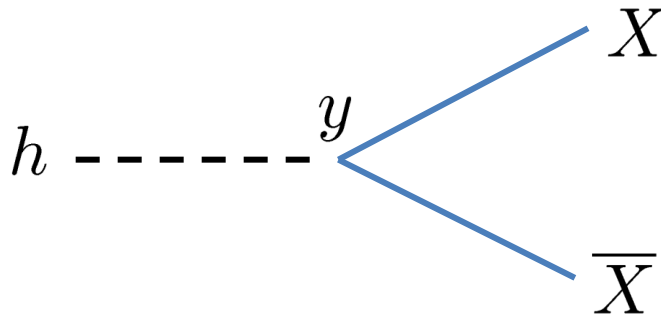
$$\text{Expected rate} = \sigma_{\text{prod}} \times \frac{\Gamma(H \rightarrow X\bar{X})}{\Gamma_H}$$

# Expected Rates for the Higgs boson

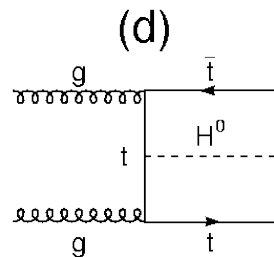
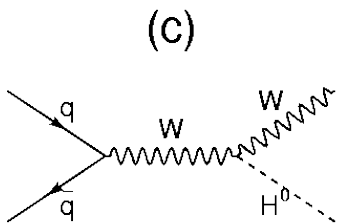
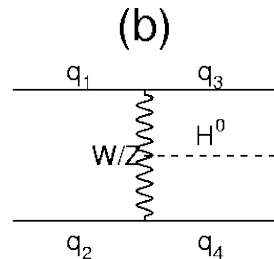
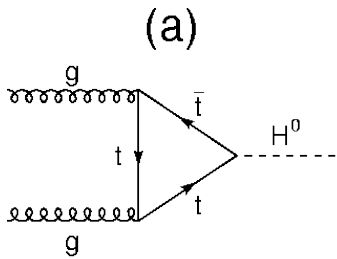
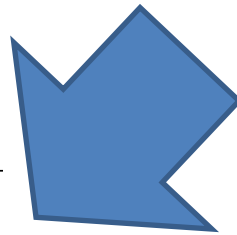
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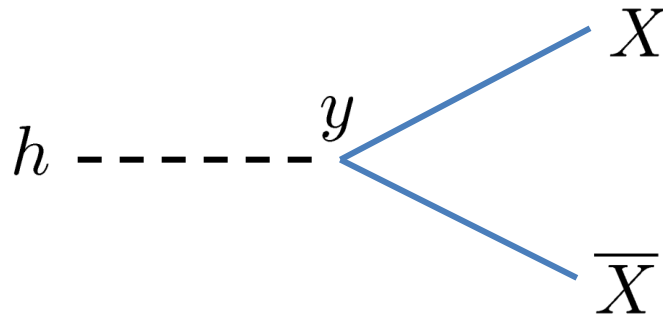
# Expected Rates for the Higgs boson



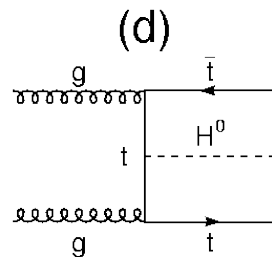
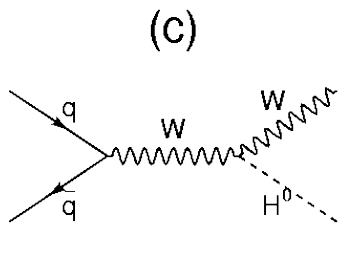
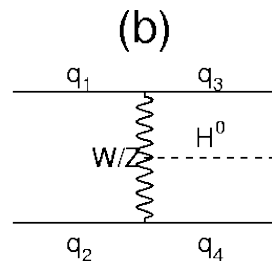
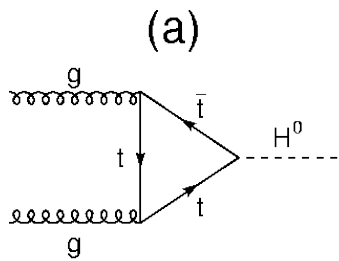
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# Expected Rates for the Higgs boson



Expected rate =  $\sigma_{\text{prod}} \times \frac{\Gamma(H \rightarrow X\bar{X})}{\Gamma_H}$

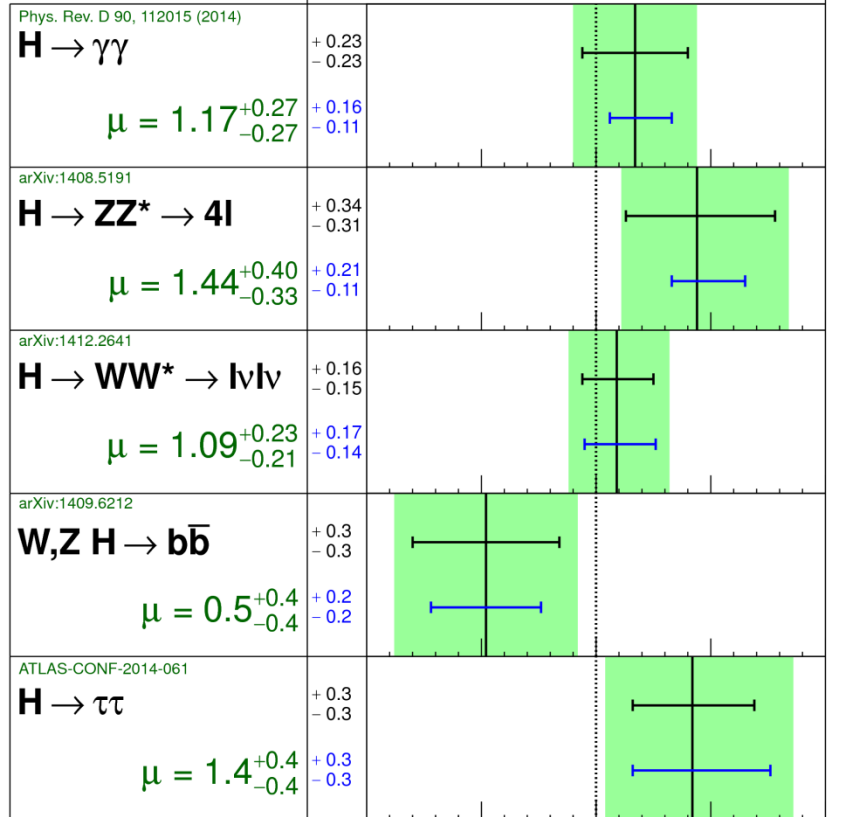


$$\Gamma_H = \sum_i \Gamma(H \rightarrow X_i \bar{X}_i)$$

# Experimental results

**ATLAS Prelim.**  
 $m_H = 125.36 \text{ GeV}$

—  $\sigma(\text{stat.})$  Total uncertainty  
 —  $\sigma(\text{sys inc.})$   $\pm 1\sigma$  on  $\mu$   
 —  $\sigma(\text{theory})$



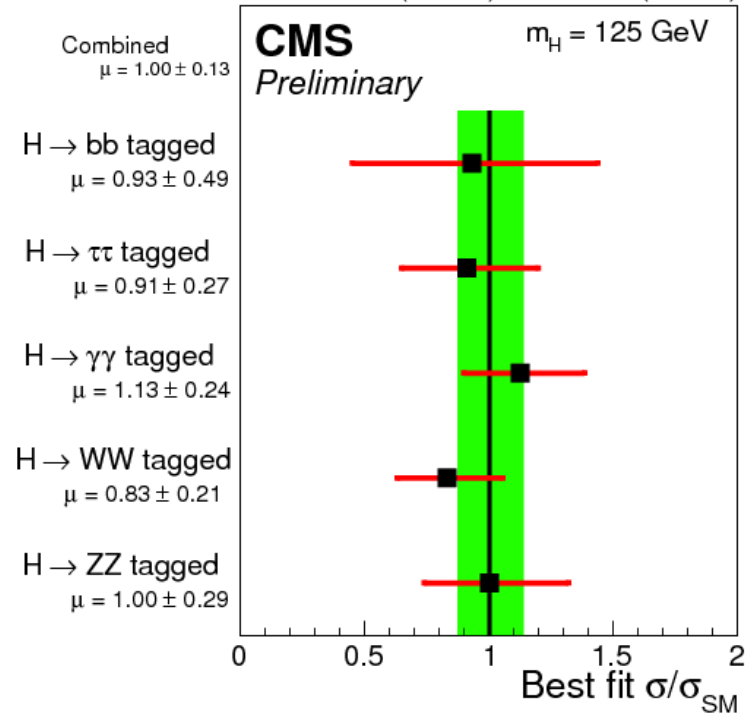
$\sqrt{s} = 7 \text{ TeV} \int L dt = 4.5\text{-}4.7 \text{ fb}^{-1}$

$\sqrt{s} = 8 \text{ TeV} \int L dt = 20.3 \text{ fb}^{-1}$

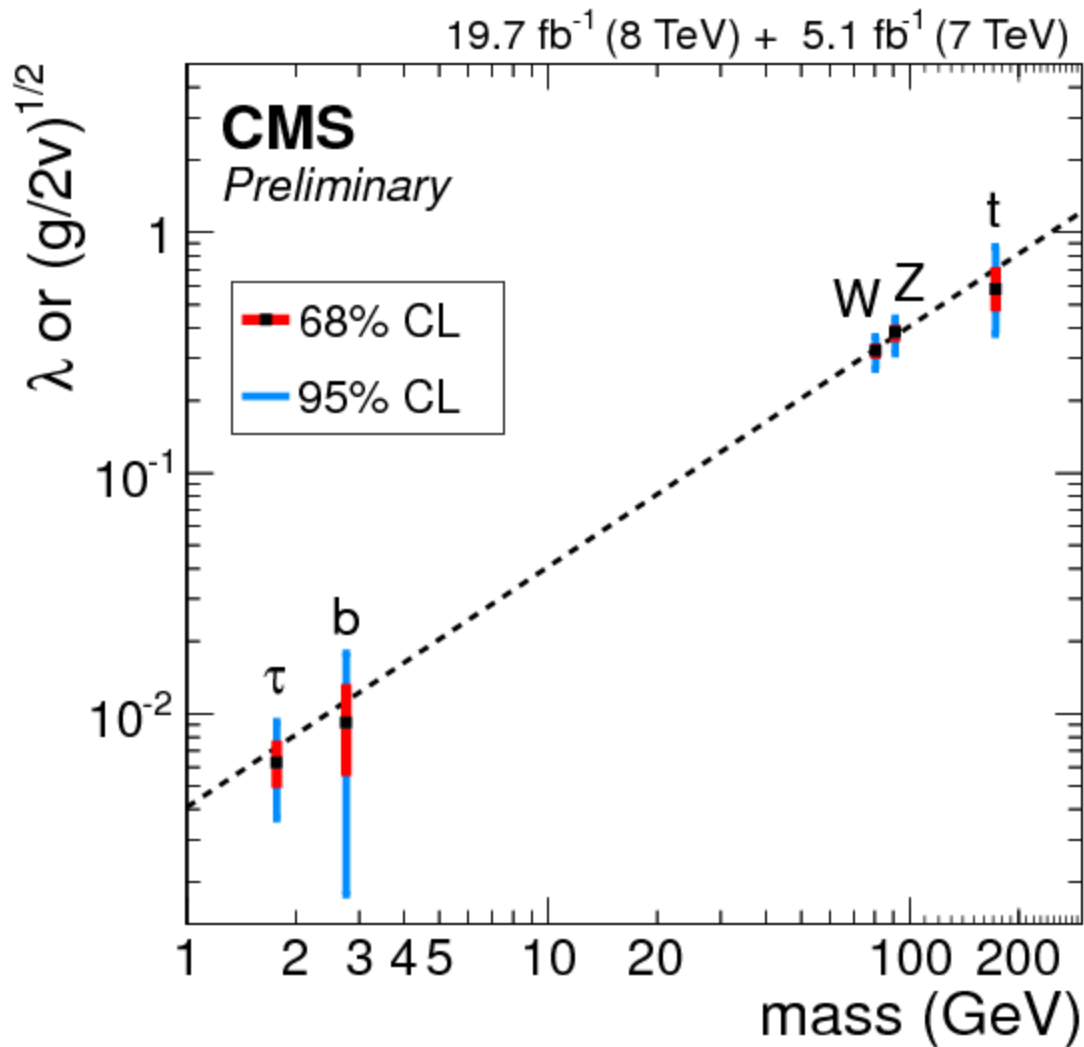
Signal strength ( $\mu$ )

released 12.01.2015

19.7  $\text{fb}^{-1}$  (8 TeV) + 5.1  $\text{fb}^{-1}$  (7 TeV)



# Turning observed rates into constraints on Higgs coupling



# Problems in extracting couplings

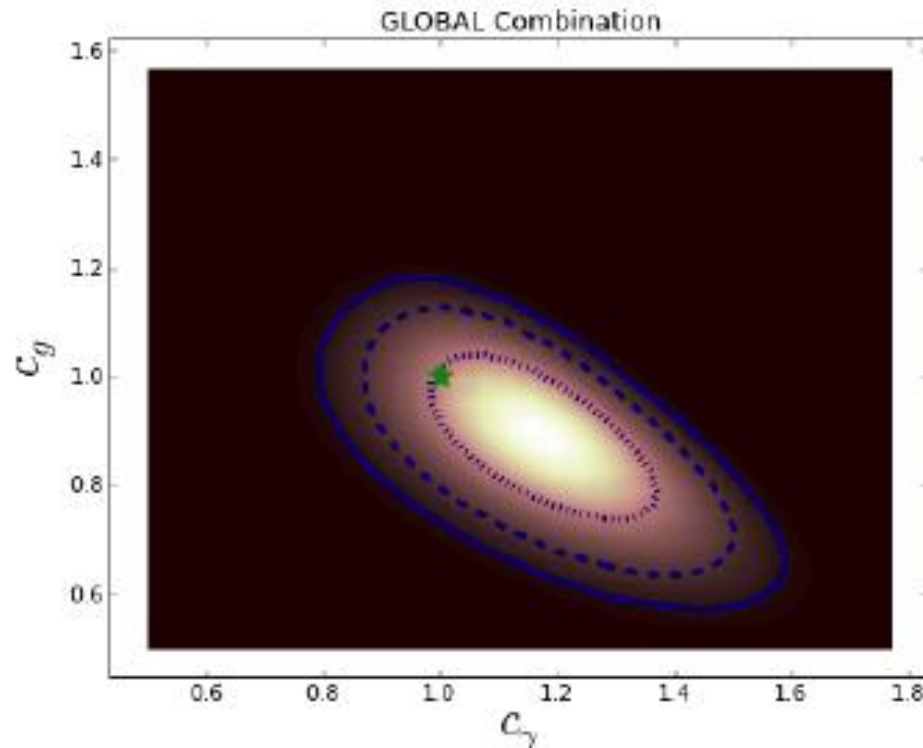
- A number of degeneracies (e.g. LHC flat direction).
- HWW coupling is important for consistency of the unitarization of the SM at high energies.
- However there are degeneracies in  $g_{HWW}$  measurements.
- Hgg coupling is sensitive to new colored particles that couple to the Higgs boson.
- Hgg coupling can not be directly measured because of the hadronic final state.



**Can we break some of these  
degeneracies by measuring the  
production modes?**

# Gluon couplings from global fit

$$\mathcal{L}_\Delta = - \left[ \frac{\alpha_s}{8\pi} c_g b_g G_{a\mu\nu} G_a^{\mu\nu} + \frac{\alpha_{em}}{8\pi} c_\gamma b_\gamma F_{\mu\nu} F^{\mu\nu} \right] \left( \frac{H}{V} \right)$$



$$c_\gamma = 1.18 \pm 0.12, \quad c_g = 0.88 \pm 0.11$$

**Global fit measurement of gluon coupling is indirect.**

**Can we get another handle on Higgs  
coupling to gluons and production  
mechanisms in general?**

# Separating Higgs production modes

## Naïve approach:

- a) Kinematic cuts on VBF/GF (forward jets)
- b) further kinematic cuts

# Kinematic separation: Rapidity gap

- Consider  $pp \rightarrow H + jj$  with  $H \rightarrow \gamma\gamma$
- Cuts: Large rapidity gap (CMS tight cuts)

$$\Delta\eta_{jj} > 3.5$$

$$M_{jj} > 500 \text{ GeV}$$

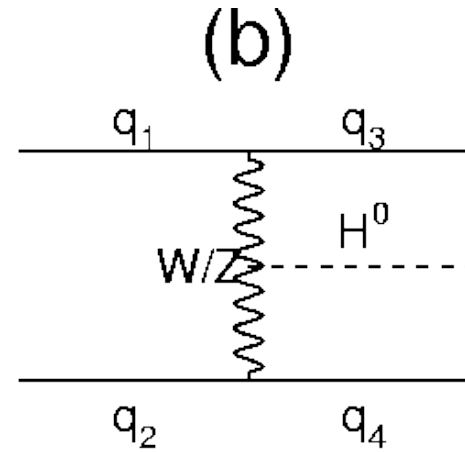
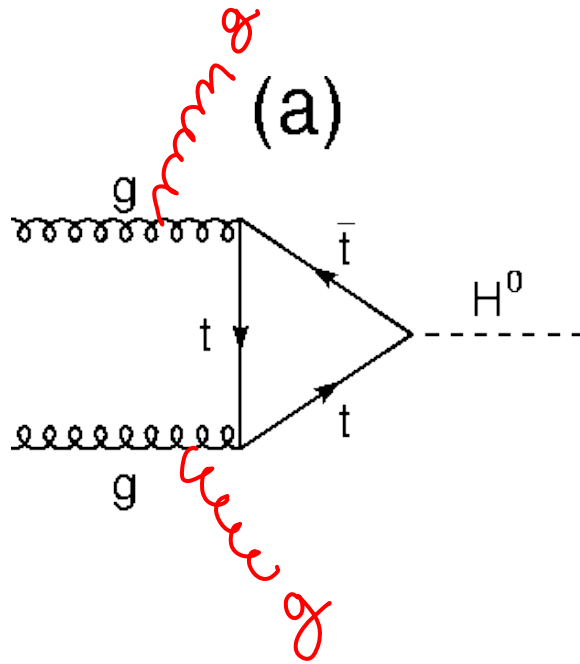
Tight

$$M_{jj} > 250 \text{ GeV}$$

loose

- Even after imposing these cuts sizeable GF contamination  $\sim 20\text{-}30\%$  and an  $O(1)$  background

# Contamination



# Kinematic Separation

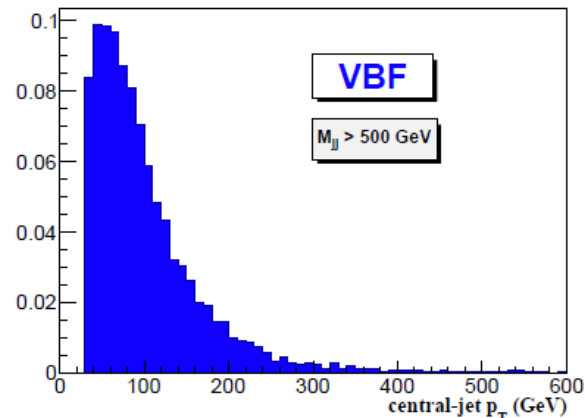
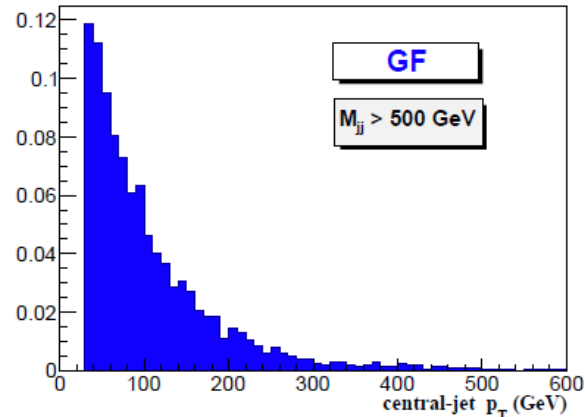


FIG. 1: Normalized  $p_T$  distribution of the central jet for GF (upper panel) and for VBF (lower panel) in  $H + 2$  jets events passing the tight selection cuts with  $M_{jj} > 500$  GeV.

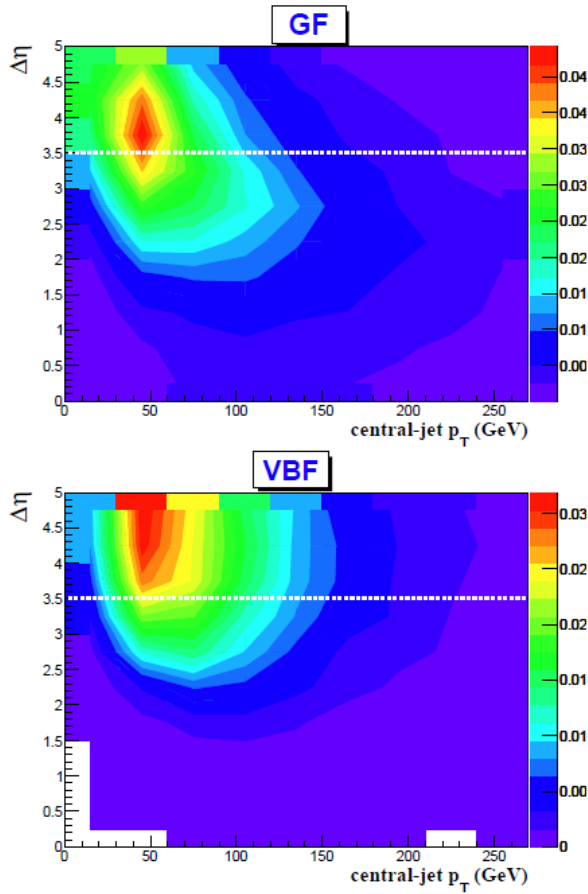


FIG. 2:  $p_T$  of the central jet vs  $\Delta\eta$  of the two jets for GF (upper panel) and for VBF (lower panel) in  $H + 2$  jets events, when only mild cuts on jets are applied. The dotted white line shows the value of the cut on  $\Delta\eta$  applied in the analysis.

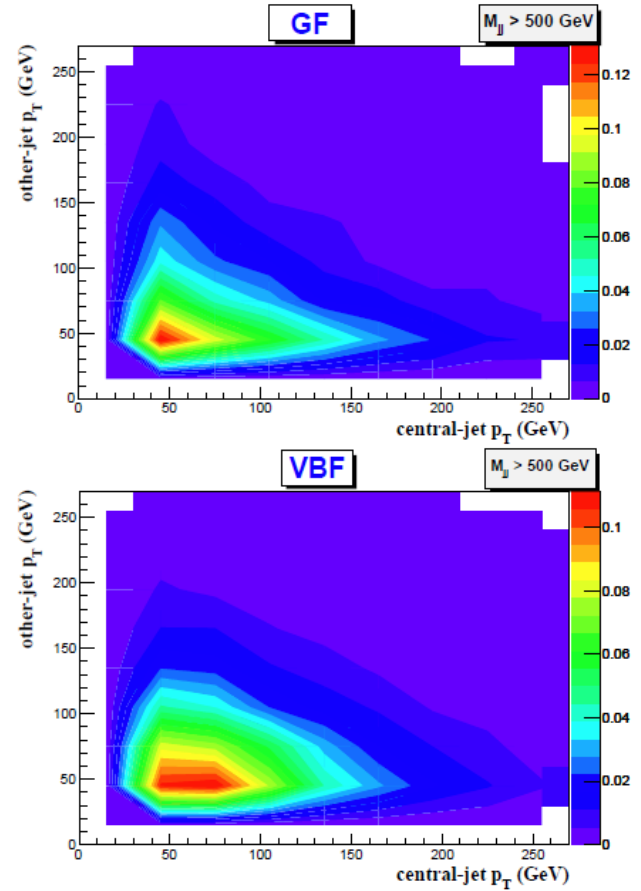


FIG. 3:  $p_T$  of the central jet vs  $p_T$  of the other jet for GF (upper panel) and for VBF (lower panel) in  $H + 2$  jets events passing the tight selection cuts with  $M_{jj} > 500$  GeV.



# Separating Higgs production modes

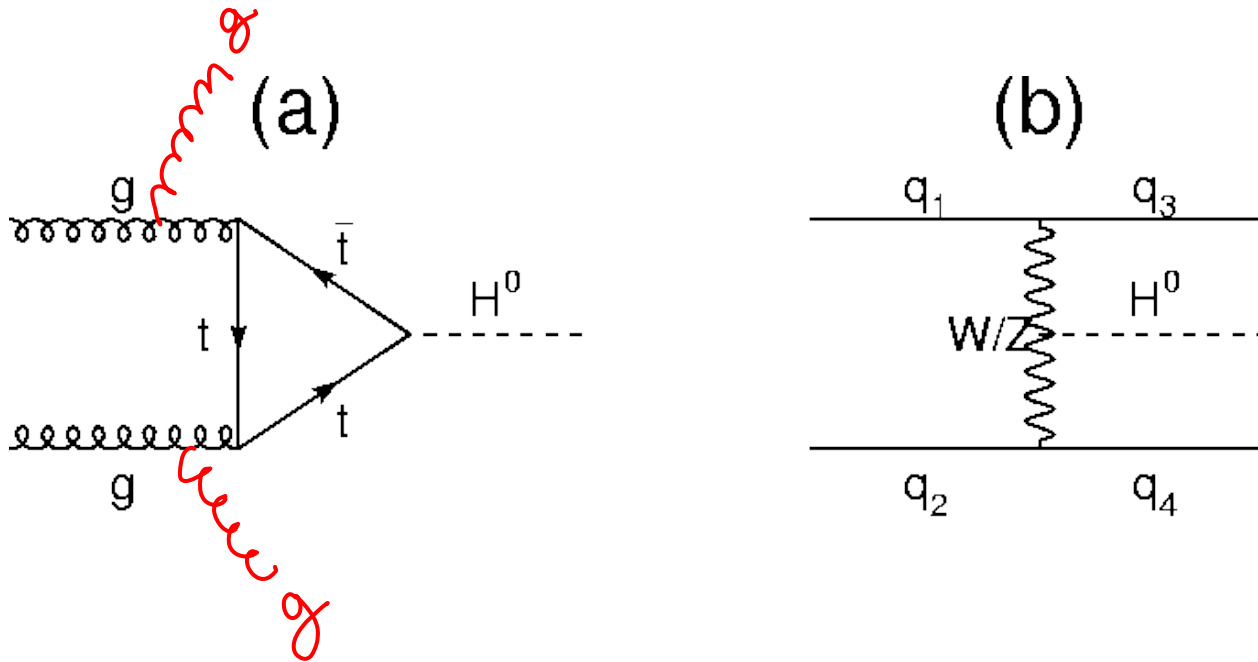
Naïve approach:

- a) Kinematic cuts on VBF/GF (forward jets)
- b) further kinematic cuts

Better handles:

- **Jet energy profiles: This talk**
- H + jet veto (T. Becher and M. Neubert)
- Hadronic event shapes (Englert, Spannowsky and Takeuchi)
- Matrix element method (Andersen, Englert and Spannowsky)
- Third jet veto (Cox, Forshaw, and Pilkington)

# An observation



- Jets associated with GF are mostly gluon like
- Jets associated with VBF are always quark like

**Any method to statistically measure ratio of quark and gluon jets efficiently could pin down the ratio of GF to VBF like events in a given Higgs sample.**

**We have proposed such a  
technique.**

# Advantages of this technique

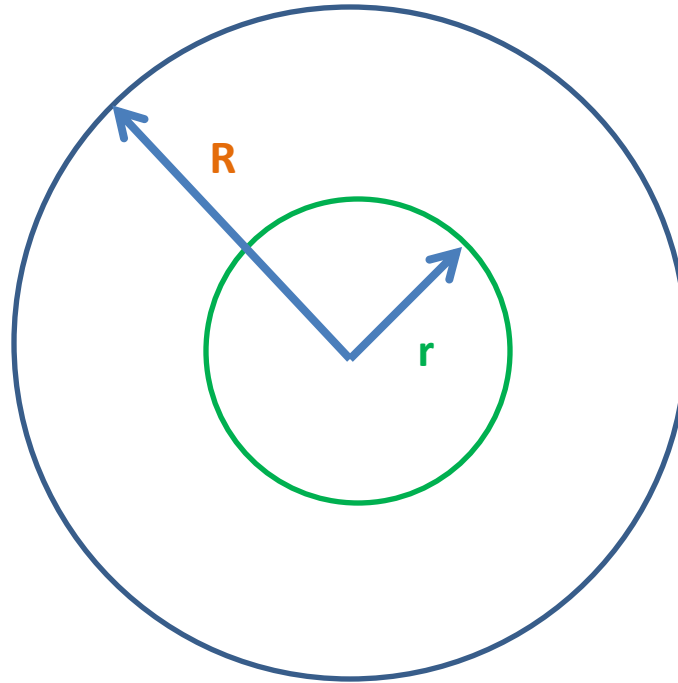
- Measurement independent of the branching fractions!

$$\text{Observed GF rate} = \sigma_{\text{GF}} \times \frac{\Gamma(H \rightarrow \gamma\gamma)}{\Gamma_H}$$

$$\text{Observed VBF rate} = \sigma_{\text{VBF}} \times \frac{\Gamma(H \rightarrow \gamma\gamma)}{\Gamma_H}$$

- Measuring ratio  $g_{Hgg} / g_{HWW}$  independently of the branching fractions
- Can be measured in many different kinematic regimes (not just with forward jets)

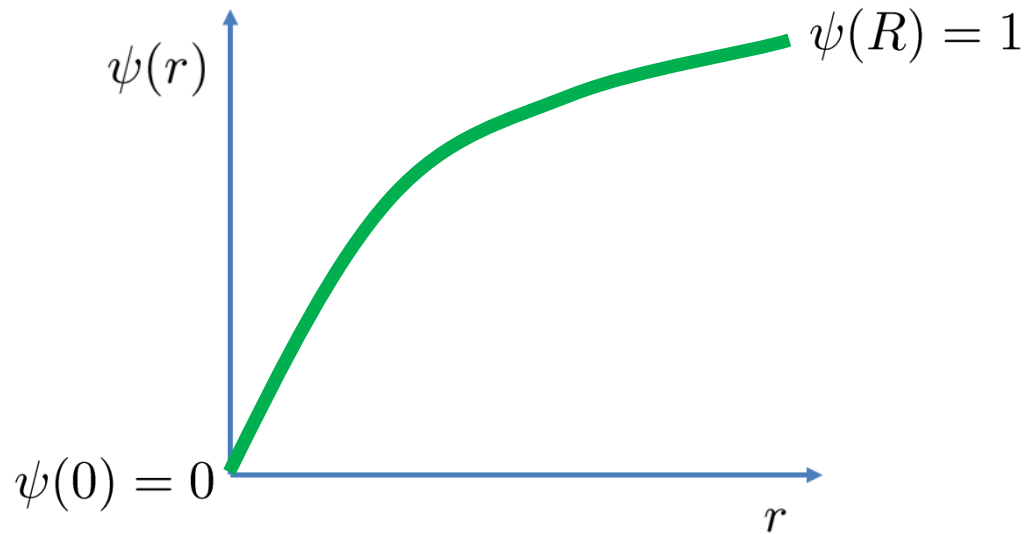
# How? Jet energy profiles



**Fraction of total jet  $p_T$  in a sub-cone of size  $r$ , inside a jet or size  $R$**

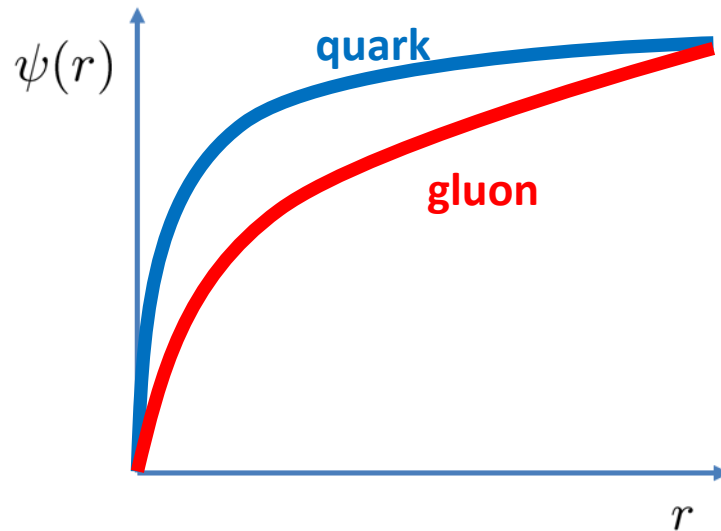
$$\psi(r) = \frac{\int_0^r \frac{dp_T}{dr'} dr'}{\int_0^R \frac{dp_T}{dr'} dr'}$$

# What to expect for the JEP



$R$  = jet cone size during clustering ( $\sim 0.7$ )

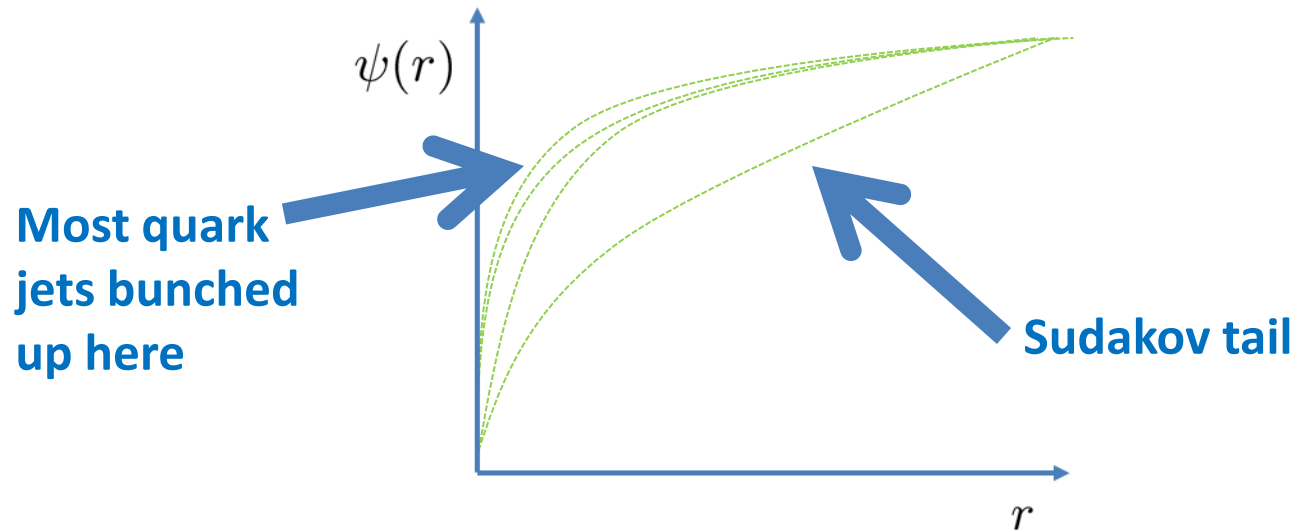
# Quark vs gluon jets



- Quark jets radiate relatively little and are narrower with a sharply rising JEP.
- Gluon jets radiate more and are broader so they have a slowly rising JEP.

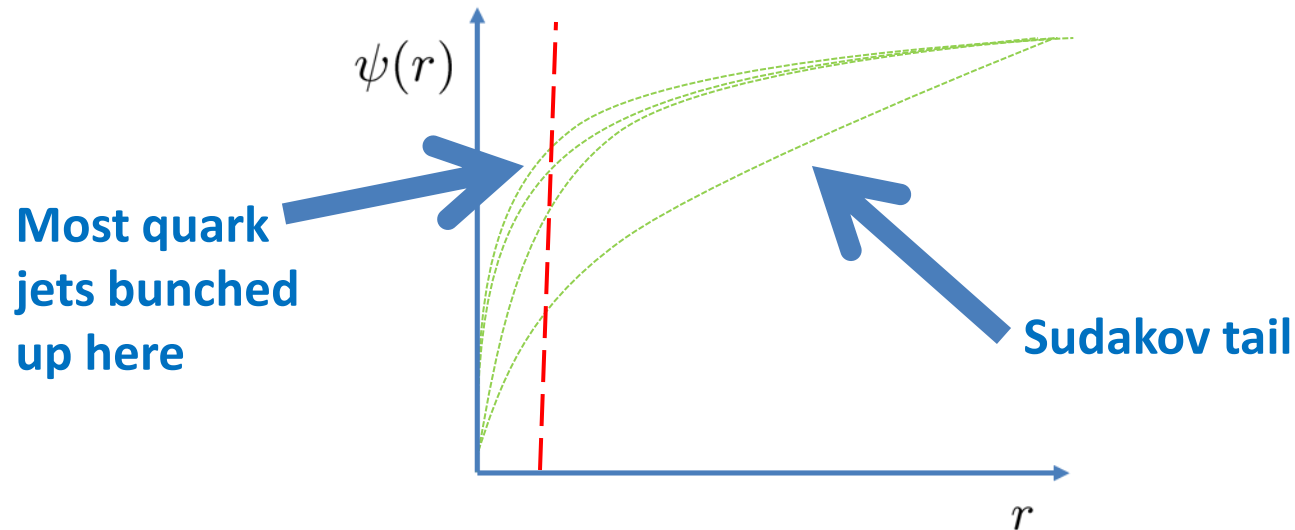


# Looking at a sample of (quark) jets



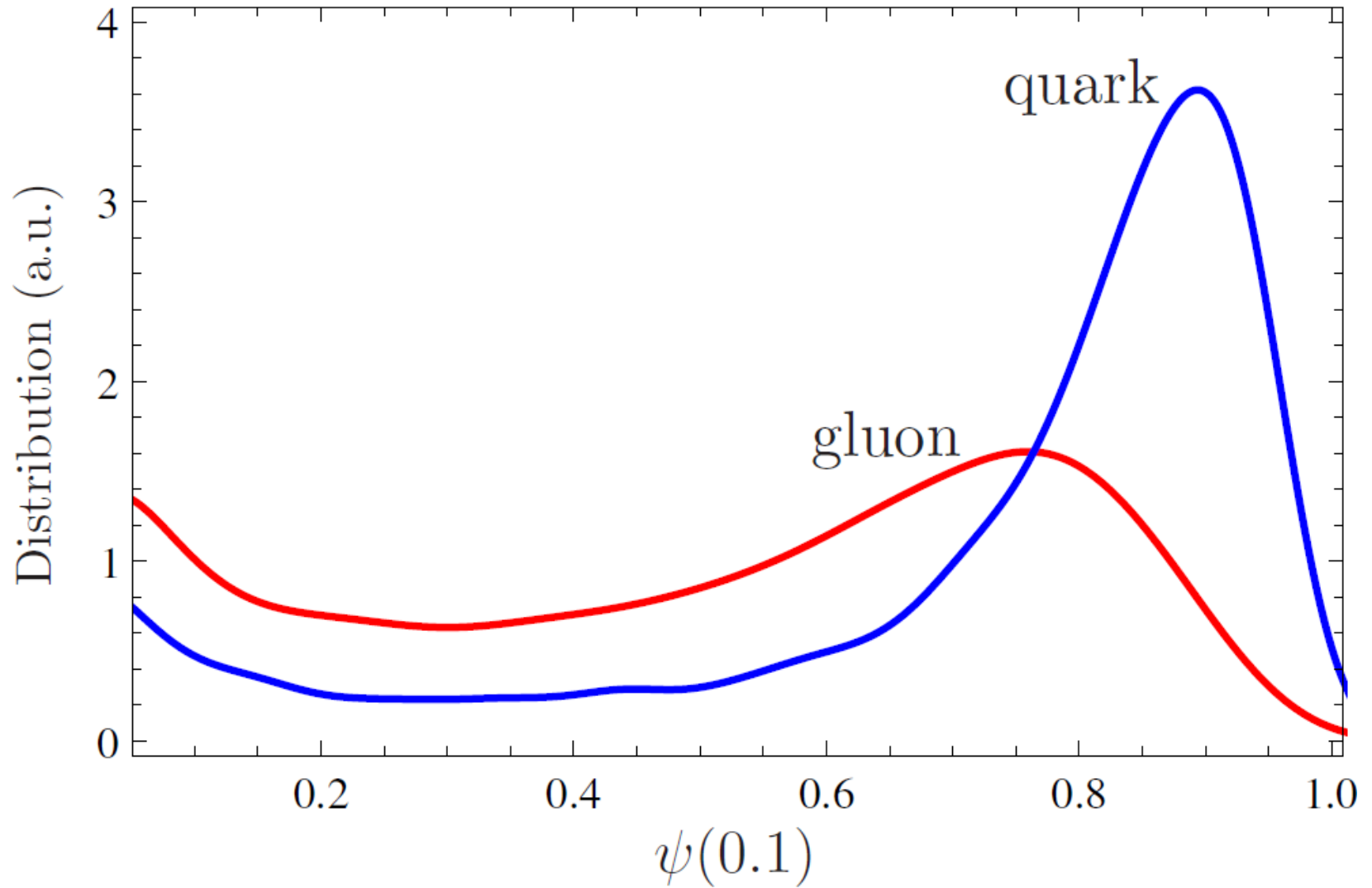
- For an individual quark/gluon jet the profile can fluctuate wildly.
- This fluctuation has an underlying distribution due to the underlying physics which is a Sudakov tail.
- The underlying distribution is not “gaussian” distributed about the average profile

# Looking at a sample of (quark) jets

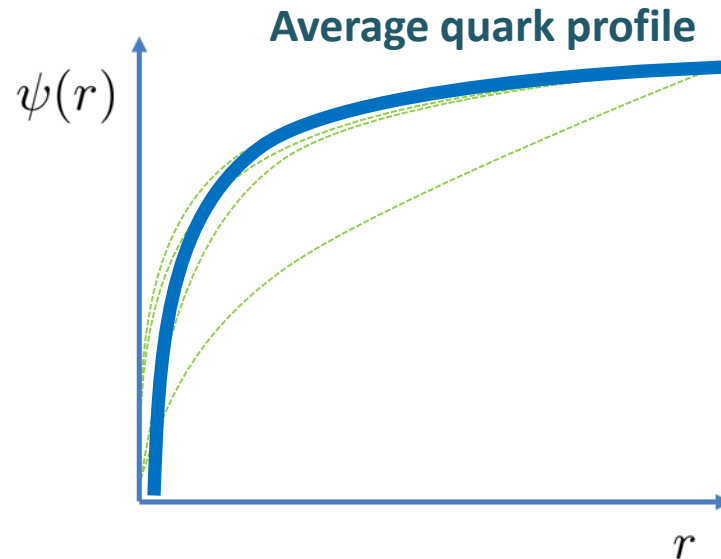


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# Slicing the JEP

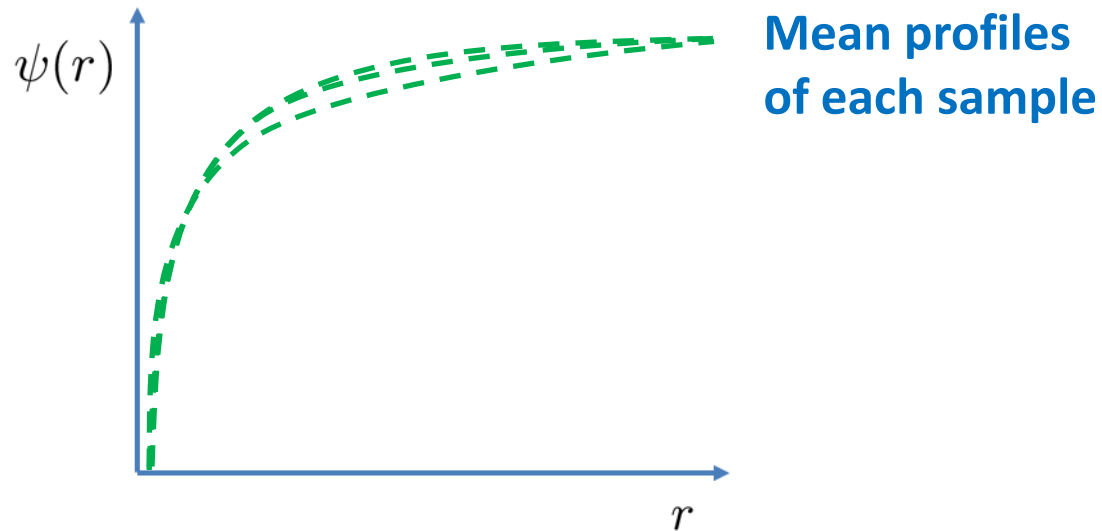


# Looking at a sample of (quark) jets



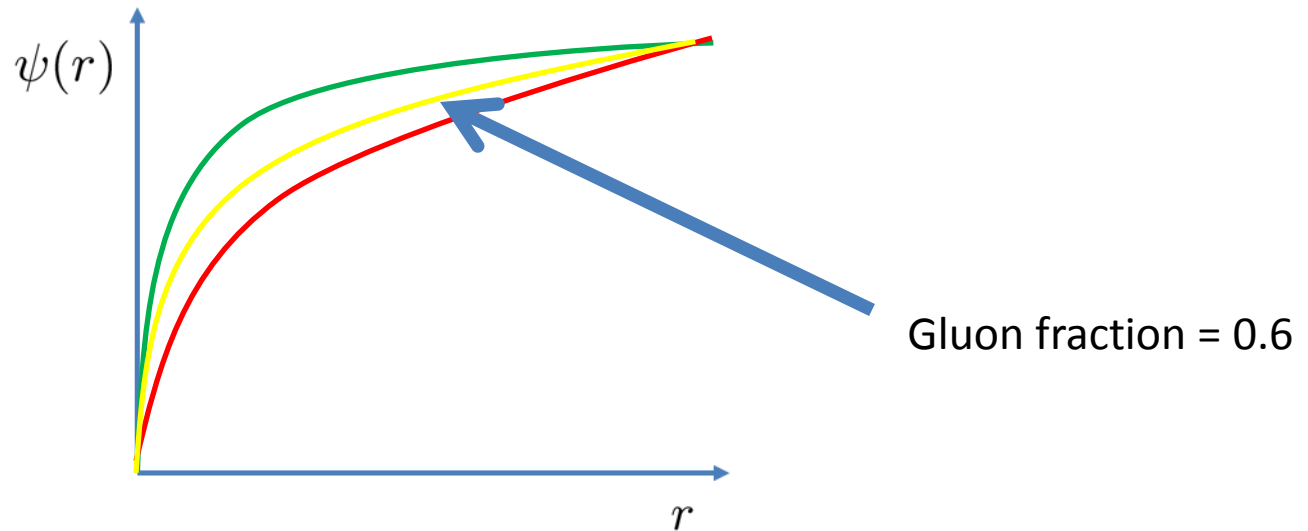
- For an individual quark/gluon jet the profile can fluctuate wildly.
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# Pseudo experiments of samples



- Consider many pseudo-experiments of  $N_{\text{exp}}$  quark jets.
- The average profile of this sample fluctuates less wildly.
- As a rule of thumb, for  $> 30$  events in the sample, the fluctuation in the average profile of the sample IS gaussian.

# From quarks and gluons to weighted samples



- Instead of talking about samples with pure quarks or pure gluons, we can talk about samples with a specific gluon fraction.
- The average profile is just a linear weighting of the average quark and gluon profiles.

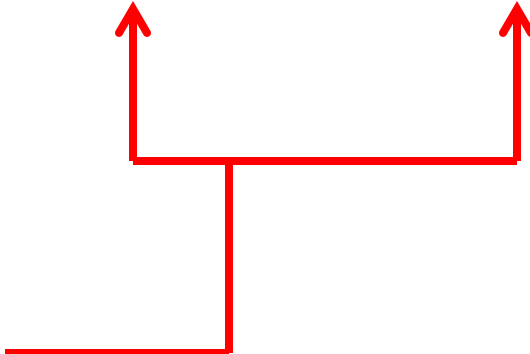
# Expected Average Profile

$$\psi(r)_{\text{EAP}} = \int \left( \frac{dN_q}{dp_T} \psi_q(r, p_T) + \frac{dN_g}{dp_T} \psi_g(r, p_T) \right) dp_T / (N_q + N_g)$$

# Expected Average Profile

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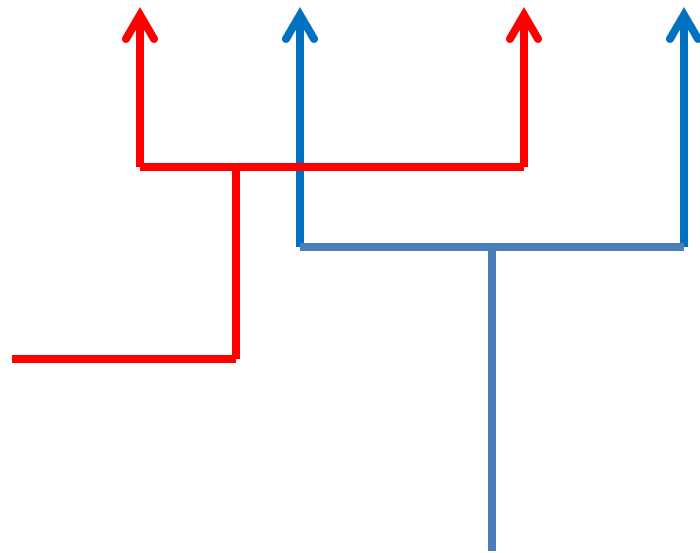
**Fluctuations of  
the gluon  
fraction of a  
given pseudo  
experiment  
(hard process)**





# Expected Average Profile

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**Fluctuations of the gluon fraction of a given pseudo experiment (hard process)**

**Fluctuations due to variation of individual quark/gluon jet energy profiles (soft process)**

# Pseudo experiments of samples

- **Pseudo Experiments:** For a given luminosity from Monte-Carlo we can generate samples of events with fluctuations in the number of total events and fraction of quark/gluon events (hard process) and fluctuations in the jet energy profiles (soft processes).
- For a given sample of  $N$  events we can study the average JEP.

# Strategy to separate VBF from GF

- Find the average profile for a SM like sample and the expected error.  
(Experimental measurement should lie within the error bars of this sample)

For comparison:

- Find the average profile for a pure VBF sample and the expected error.
- Find the average profile for a pure GF sample and the expected error.

# Three ways to determine the JEP

- Experimental data (control samples of pure quark or gluon jets or known gluon fraction)
- Theoretical calculations (NLO parton splitting or LL resummation)
- Pythia (tune dependent but allows statistical fluctuations of pseudo-experiments to be estimated)



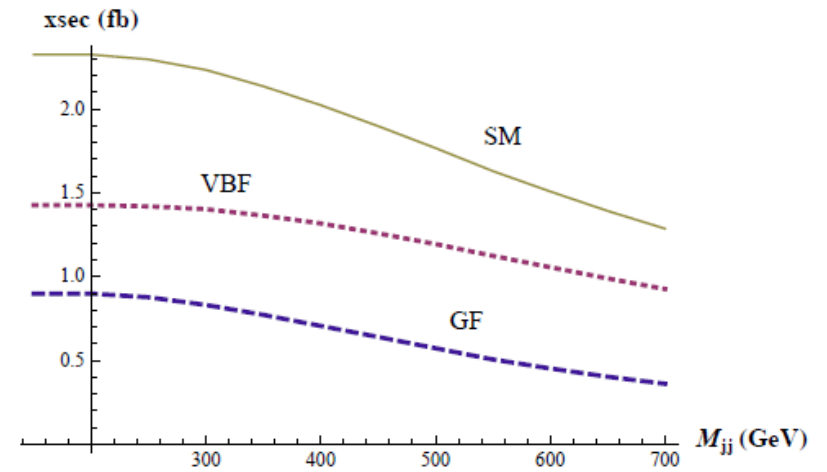
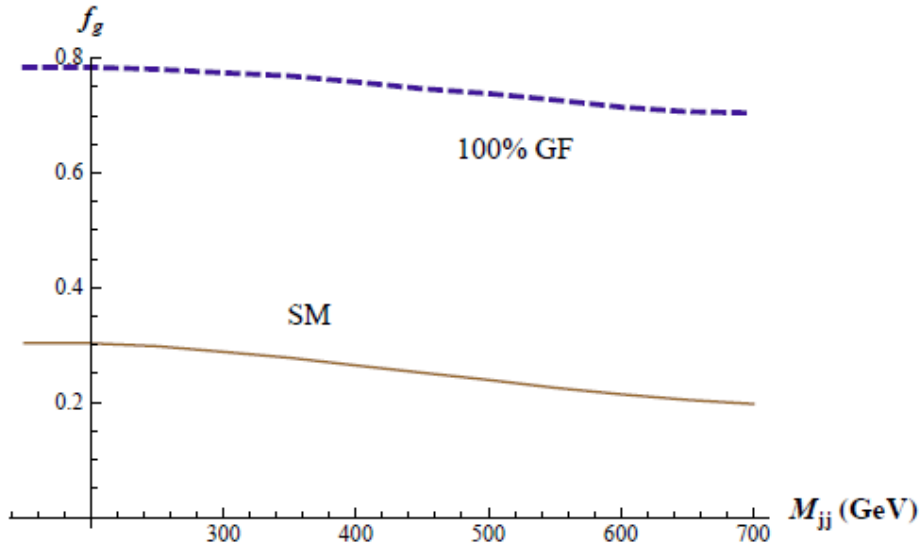
Tools available to theorists

# Separating VBF from GF

	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
14 TeV	GF	VBF	GF	VBF
$\text{MG} \times K_f^{\text{CMS}}$	32%	68%	38%	62%
	0.57 fb	1.2 fb	0.88 fb	1.4 fb

TABLE II: SM expected cross-sections at the 14 TeV LHC, using tight cuts with  $M_{jj} > 500 \text{ GeV}$  and with  $M_{jj} > 250 \text{ GeV}$ .

# Dijet invariant mass dependence



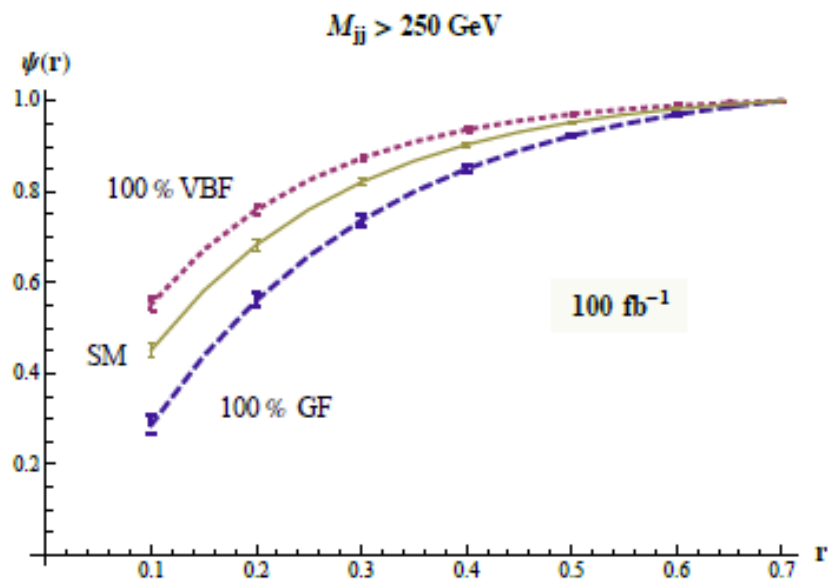
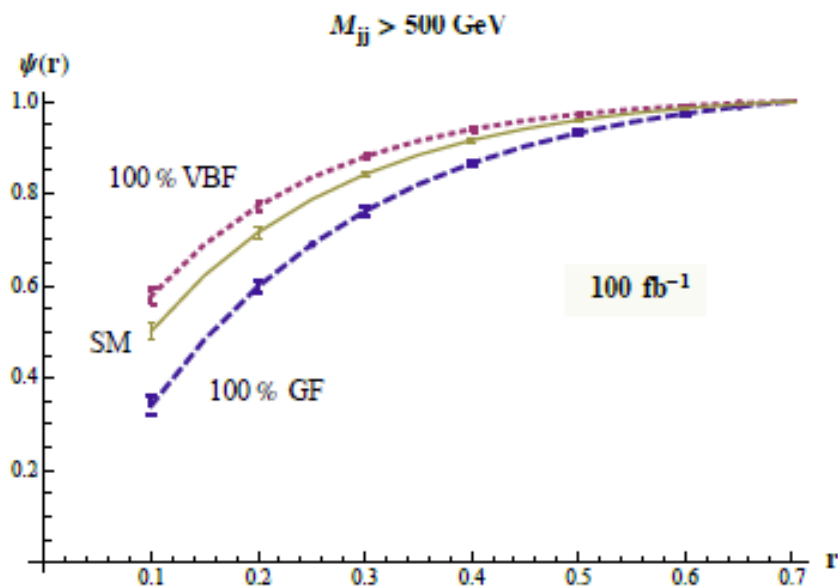
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- Find the average profile for a pure GF sample and the expected error.

# Jet energy profiles with error bars from Pythia



Caution: The error bar is the monte-carlo size of the error on the mean JEP. Individual jet profiles can fluctuate far more than the size of this error bar.



# Analytic approximation of JEPs

- We find the JEPs can be approximated by:

$$\psi(r) = \frac{1 - be^{-ar}}{1 - be^{-aR}}$$

- Define a one parameter linear interpolation between VBF and GF JEPs:

$$\psi_{f_V}(r) = f_V \psi_{\text{VBF}}(r) + (1 - f_V) \psi_{\text{GF}}(r)$$

- $f_V$  parameterizes the VBF fraction of the sample.
- The errors on the JEPs can be translated into errors on the fitted  $f_V$ .

# Measured value of $f_V$ with errors

$f_V$	$M_{jj} > 500 \text{ GeV}$	$M_{jj} > 250 \text{ GeV}$
SM	$0.68 \pm 0.05$	$0.62 \pm 0.04$
VBF	$1.00 \pm 0.04$	$1.00 \pm 0.03$
GF	$0.00 \pm 0.06$	$0.00 \pm 0.05$

Compare this to the simulated cross-section:

	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
14 TeV	GF	VBF	GF	VBF
$MG \times K_f^{\text{CMS}}$	32%	68%	38%	62%
	0.57 fb	1.2 fb	0.88 fb	1.4 fb

# Sensitivity and Reach

	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
	GF	VBF	GF	VBF
$\sigma$ -level	8.7	5.0	9.7	7.6

TABLE V: Expected  $\sigma$ -level distinction between SM and pure GF or VBF event samples using  $100 \text{ fb}^{-1}$  of luminosity at the 14 TeV LHC.

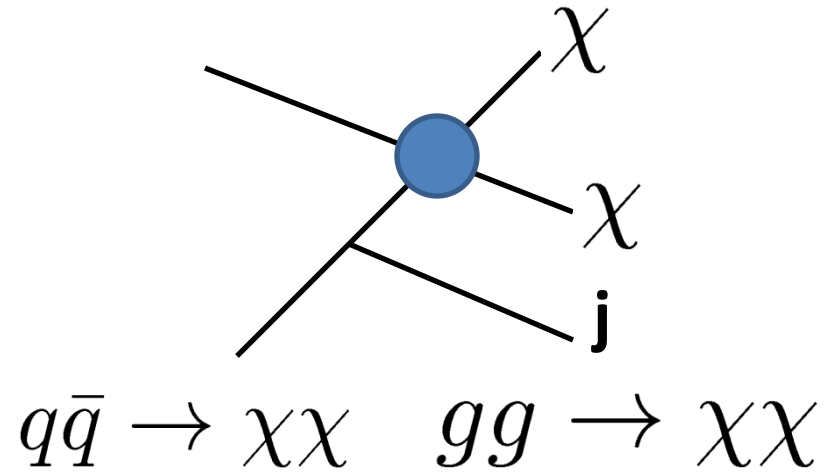
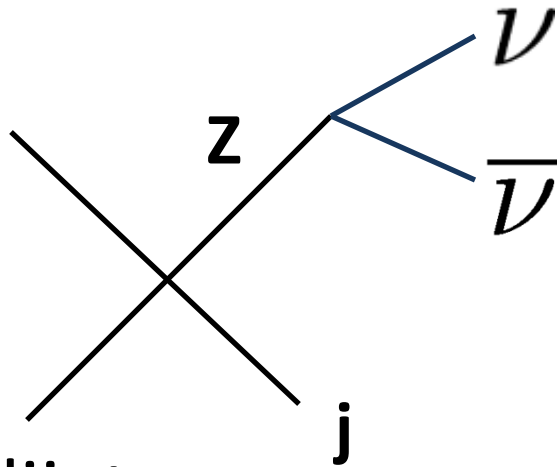
	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
	GF	VBF	GF	VBF
$5\sigma$				
Lum [ $\text{fb}^{-1}$ ]	33	100	27	43

TABLE VI: Integrated luminosity required to distinguish SM from pure GF or VBF event samples at the  $5\sigma$  level.

Lower invariant mass cut seems to be better but it also leads to increased background.

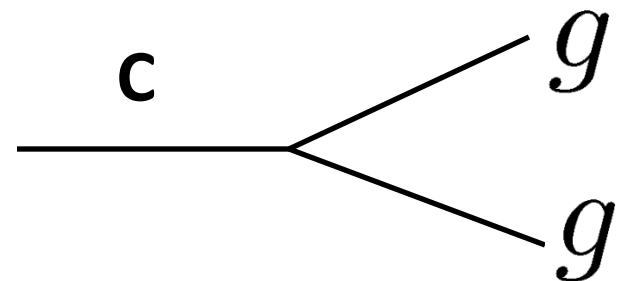
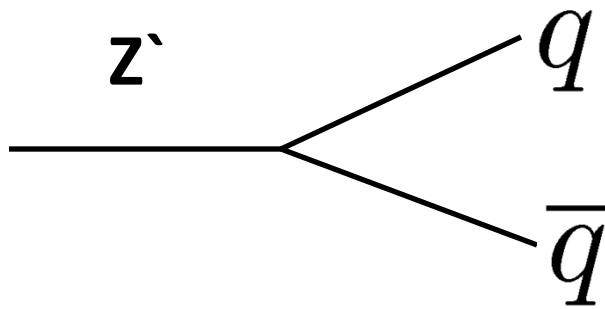
# Further applications of this technique

- Monojet searches (with P. Agrawal JHEP 1405 (2014) 098, hep-ph/1312.5325)



- New dijet resonances

R. S. Chivukula, E. H. Simmons, N. Vignaroli ,  
hep-ph/1412.3094



# Summary and Conclusions

- New Higgs observable  $f_V$  can break degeneracies in Higgs coupling extraction
- Allows identification of GF and VBF fractions to within 10% with  $100 \text{ fb}^{-1}$  of data
- Probe of Higgs coupling to gluons which is sensitive to new physics
- Independent of decay branching fractions
- Should be included in global fits
- Many possible applications of JEPs to separate quarks and gluons for new physics searches

QUESTIONS, COMMENTS, SUGGESTIONS?



# Three ways to determine the JEP

- Experimental data (control samples of pure quark or gluon jets or known gluon fraction)
- Theoretical calculations (NLO parton splitting or LL resummation)
- Pythia (tune dependent but allows statistical fluctuations of pseudo-experiments to be estimated)



Tools available to theorists

# Advantages and disadvantages of each approach

**All three should be used, each offers a different level of precision and each has its own limitations.**

- Experiments:
  1. Smallest error for low-moderate  $P_T$  jets  $\sim 200$  GeV.
  2. Suspect to systematics.
  3. No proof of factorizability (universality).
  4. Can not be extrapolated to regions where control samples are not available.
  5. Not available to theorists.
- Theory:

NLO prediction is not finite at  $r = 0$ . LL resummation provides a nice finite formula and shows factorizability but has two problems:

  1. Undetermined constants of integration.
  2. Can not generate statistical fluctuations.
- Pythia:
  1. Can generate pseudo experiments.
  2. Requires tuning.



# Estimating the effect of background

$$\psi_S(r) = \psi_{obs}(r) + \frac{B}{S} (\psi_{obs}(r) - \psi_B(r))$$

- Errors scale up by a factor  $\sqrt{1 + 2\frac{B}{S}}$

# Sensitivity including background

	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
$100 \text{ fb}^{-1}$	GF	VBF	GF	VBF
$\sigma$ level	6.4	3.6	6.4	5.0

	$M_{jj} > 500 \text{ GeV}$		$M_{jj} > 250 \text{ GeV}$	
$5\sigma$	GF	VBF	GF	VBF
Lum [ $\text{fb}^{-1}$ ]	61	190	61	100

TABLE VIII: Upper Table: Expected  $\sigma$ -level distinction between SM and pure GF/VBF event samples using  $100 \text{ fb}^{-1}$  of luminosity at the 14 TeV LHC including the estimated effect of background. Lower Table: Integrated luminosity required to distinguish SM from pure GF/VBF event samples at the  $5\sigma$  level after subtracting the background JEP.

Lower invariant mass cut is better even after including background.

# Comparison of resummed JEPs to the data

Li, Li, Yuan

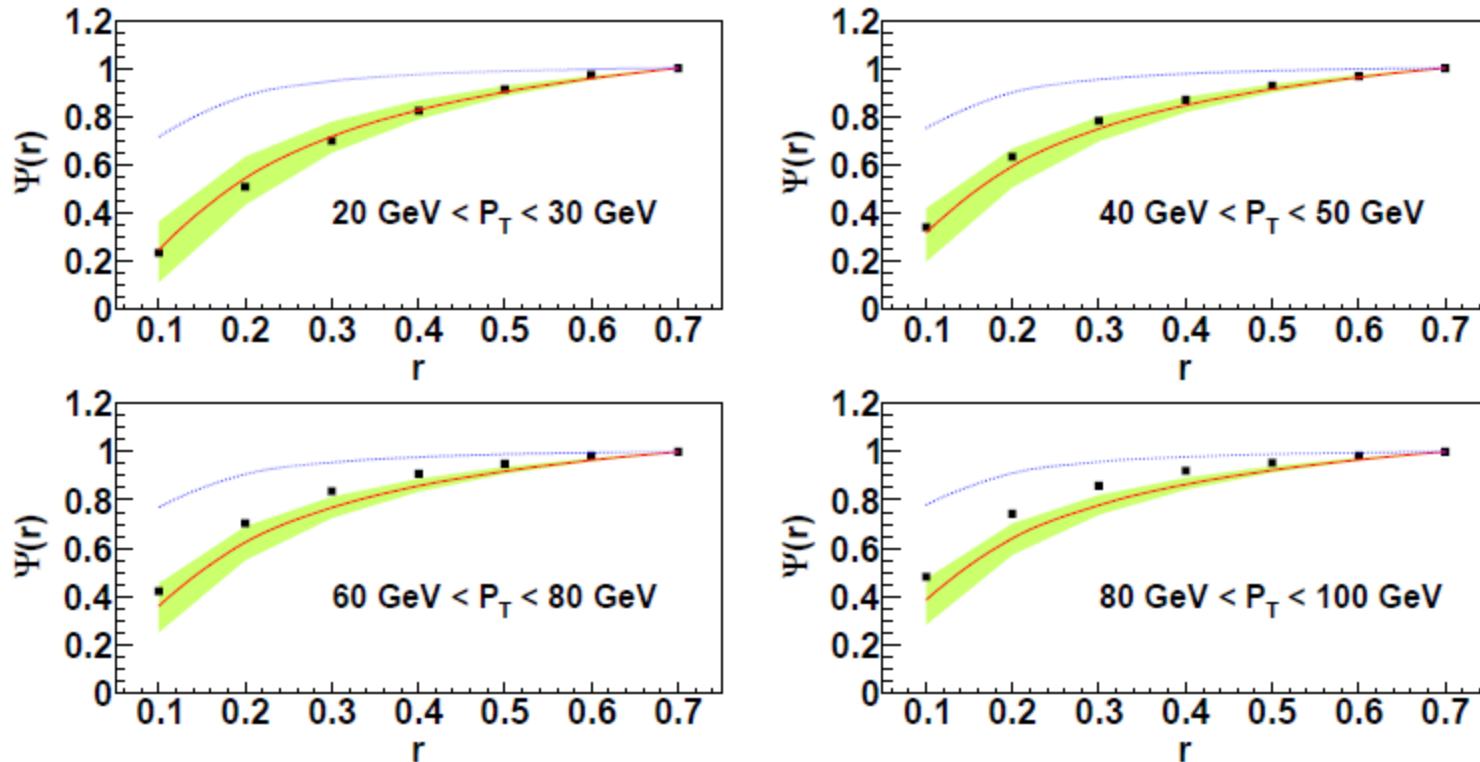


FIG. 10: Resummation predictions for the jet energy profiles with  $R = 0.7$  compared to LHC CMS data in various  $P_T$  intervals. The NLO predictions denoted by the dotted curves are also displayed.

**NLO: Blue line**

**LL: Red line**

**Black points: data WITH error bars**

The LL resummation calculation has a constant that parameterizes the NLL contribution. Varying the constant gives the green error band.

# Default Pythia tune cannot be relied upon to measure the jet profile

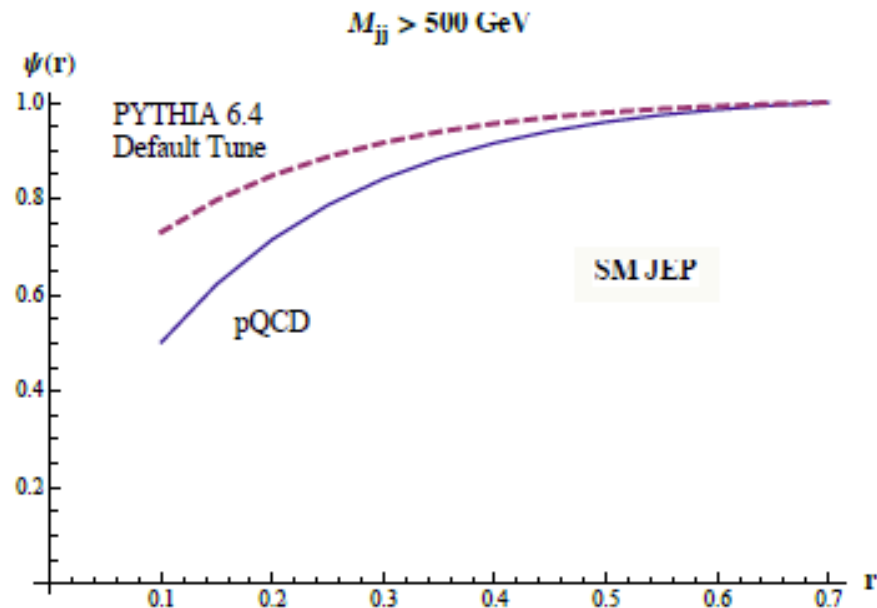
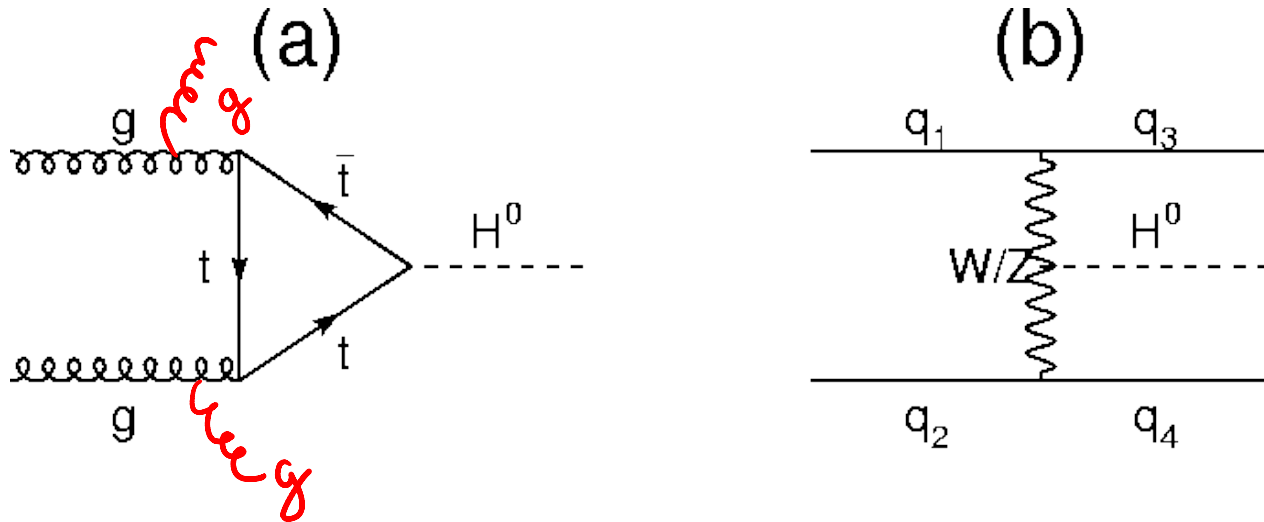


FIG. 7: Energy profile of the central jet for SM obtained by analyzing the jet substructure after Pythia v6.4 (default tune) showering, compared to the theoretical pQCD prediction using jet functions [11, 12].

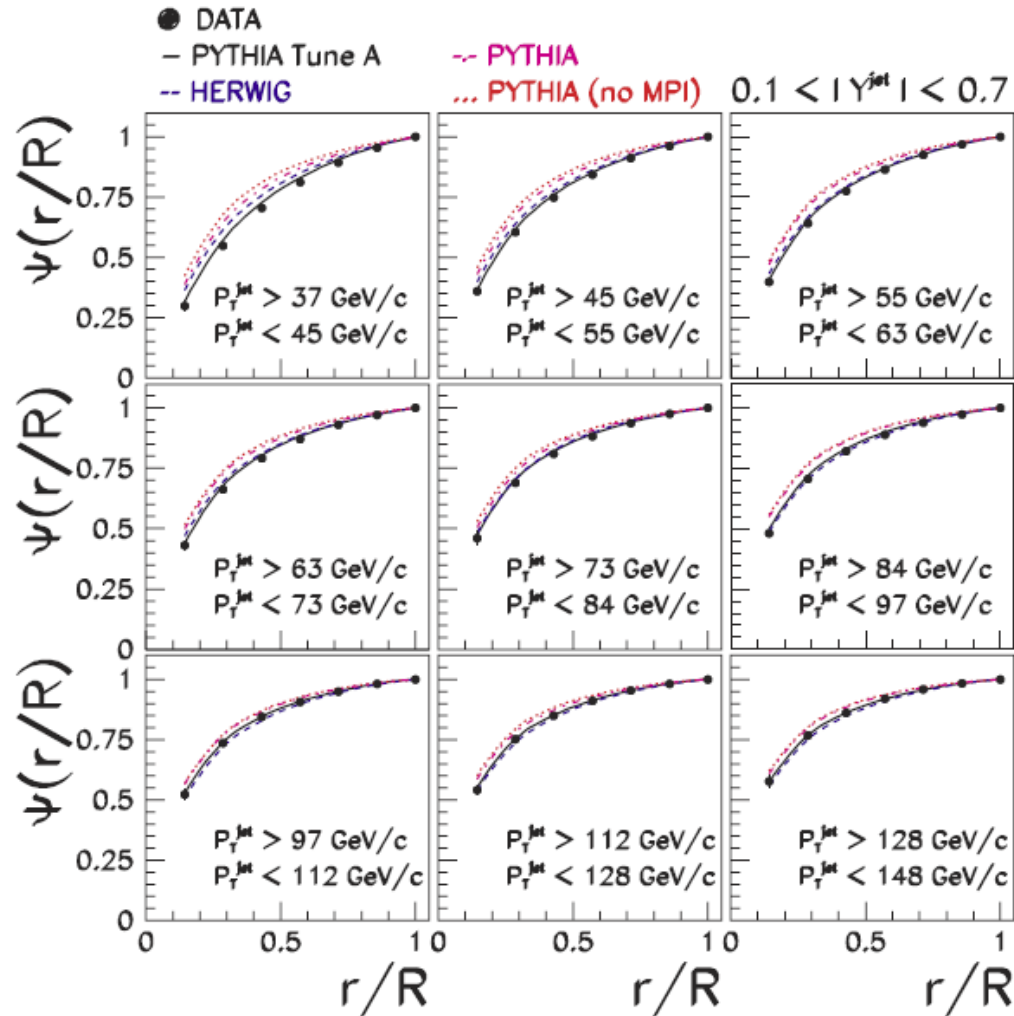
# Applying this to VBF vs GF separation



**The best approach is a hybrid approach combining all three strategies to measure JEPS.**

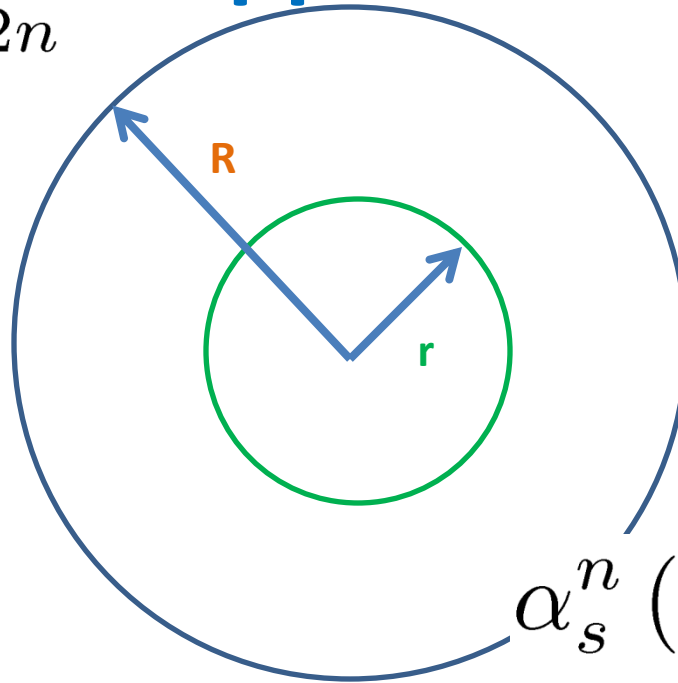
- Our choice is constrained because of lack of experimental data:
  1. We choose to use the average profile from the LL resummation calculation. The integration constants are fixed from Tevatron data and are mostly  $P_T$  independent.
  2. To estimate the error on the average profile, we conduct pseudo-experiments in (untuned) pythia and lift the error bars from the pythia JEPs and put them on the theoretical JEP.

# Experimental JEPs and Pythia (CDF)



# Jet energy profile: theoretical approach

$$\alpha_s^n (\log(R/r))^{2n}$$

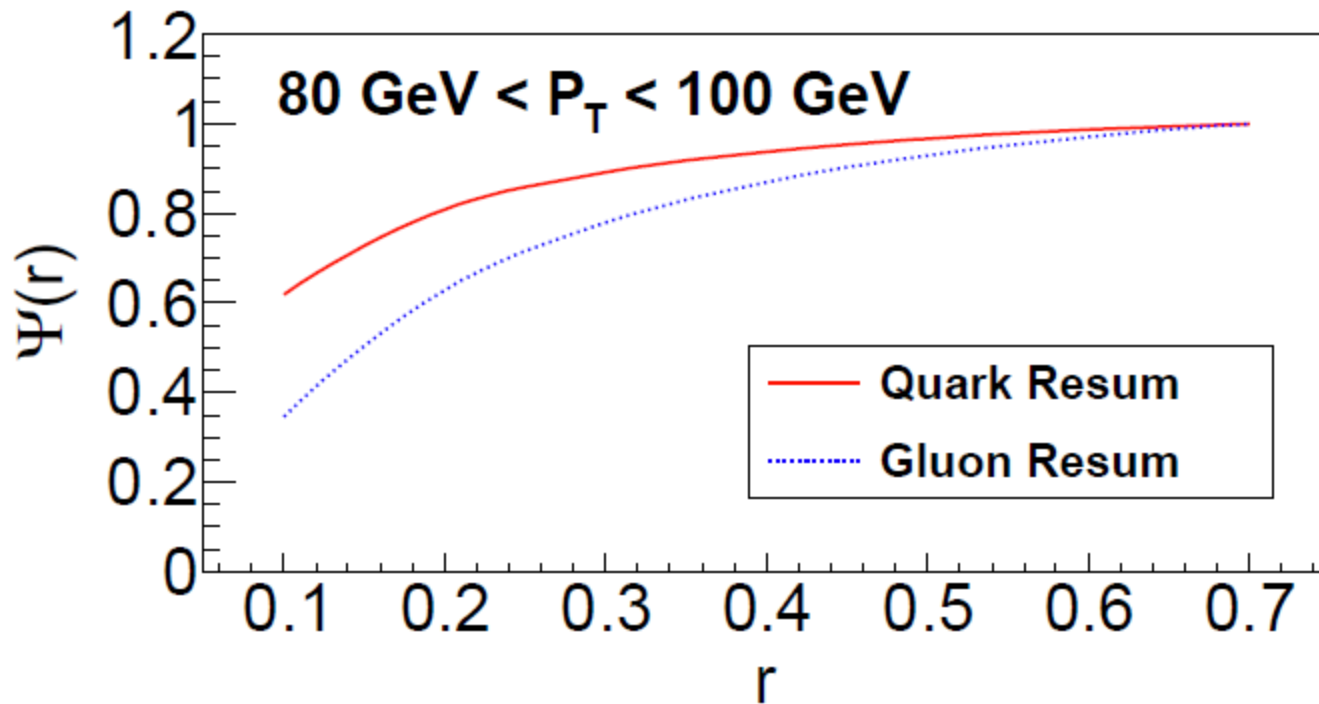


$$\alpha_s^n (\log(R/r))^{2n-1}$$

**Fraction of total jet pT in a sub-cone of size r, inside a jet or size R**

$$\Psi(r) = \frac{1}{N_J} \sum_J \frac{\sum_{r_i < r, i \in J} P_{Ti}}{\sum_{r_i < R, i \in J} P_{Ti}}$$

# Resummed jet energy profile for quark vs gluon jets





# Comparison of resummed JEPs to the data

Li, Li, Yuan

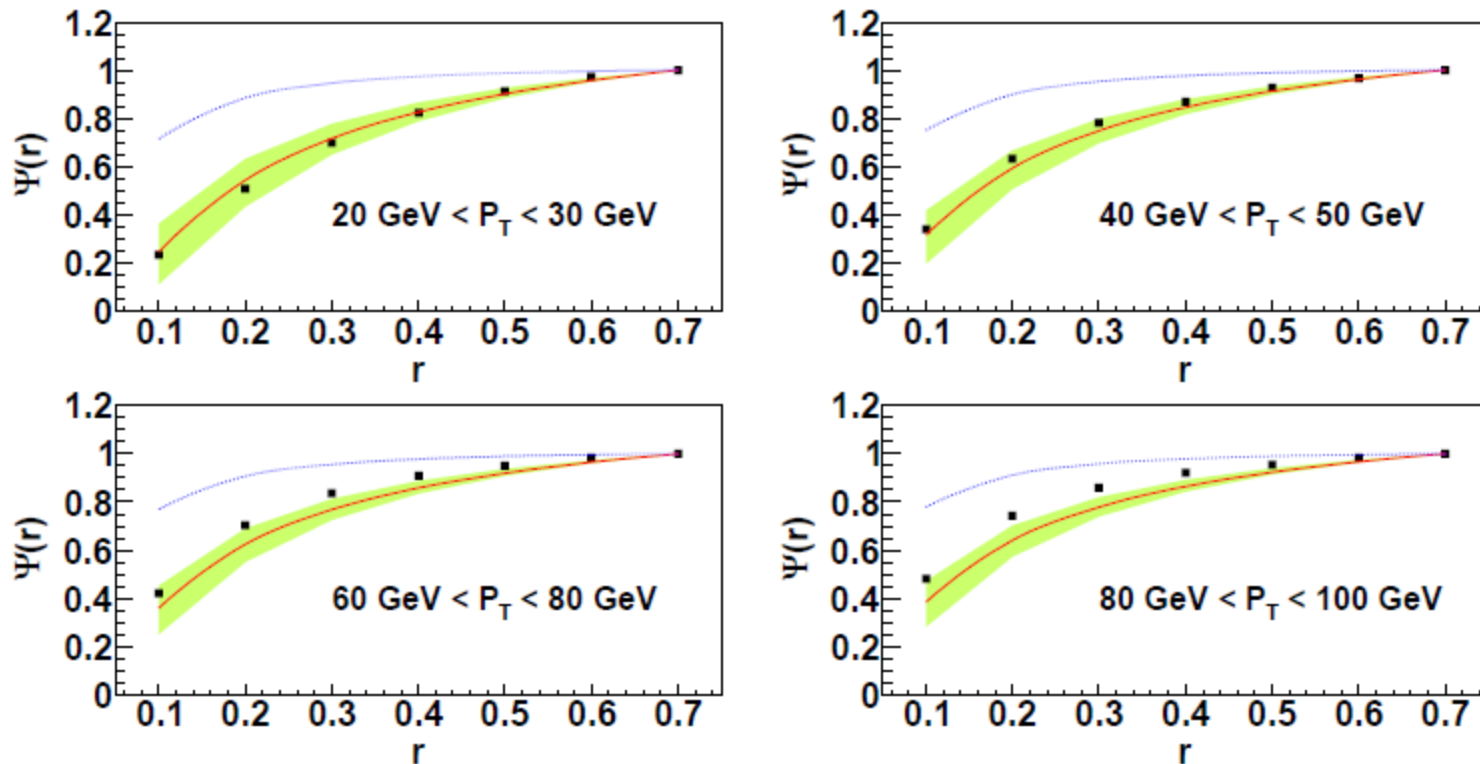


FIG. 10: Resummation predictions for the jet energy profiles with  $R = 0.7$  compared to LHC CMS data in various  $P_T$  intervals. The NLO predictions denoted by the dotted curves are also displayed.

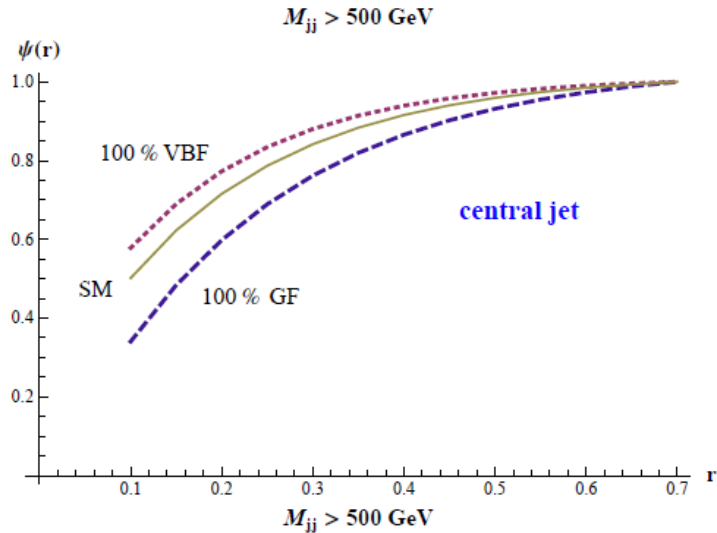
**NLO: Blue line**

**LL: Red line**

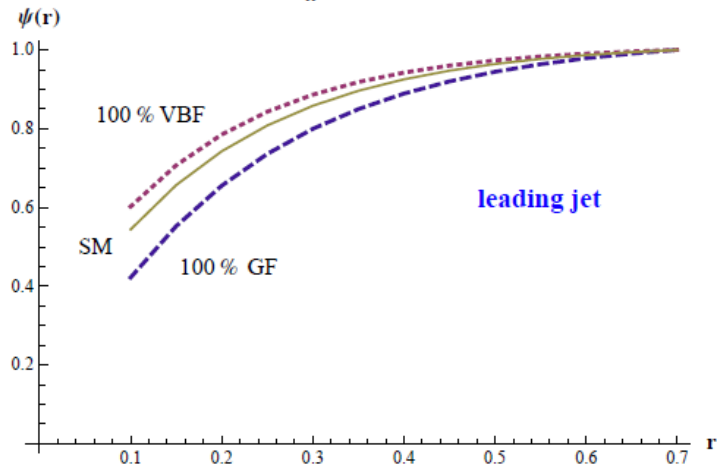
**Black points: data WITH error bars**

The LL resummation calculation has a constant that parameterizes the NLL contribution. Varying the constant gives the green error band.

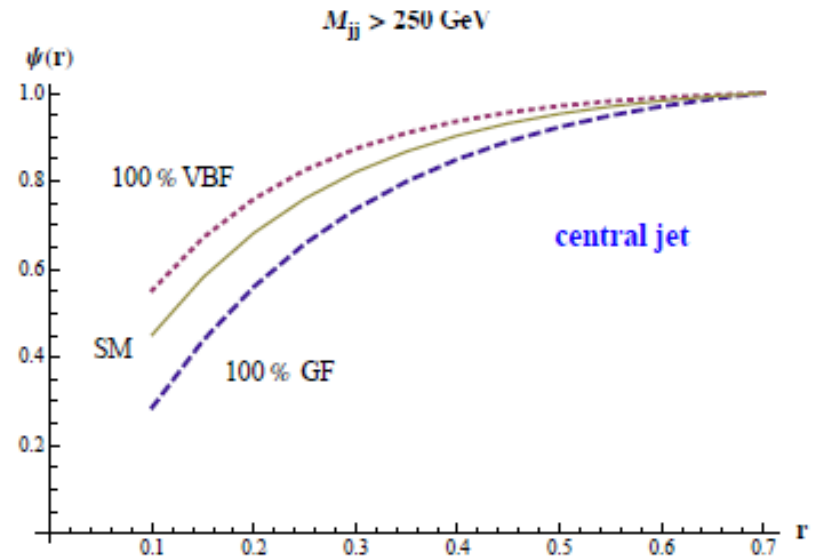
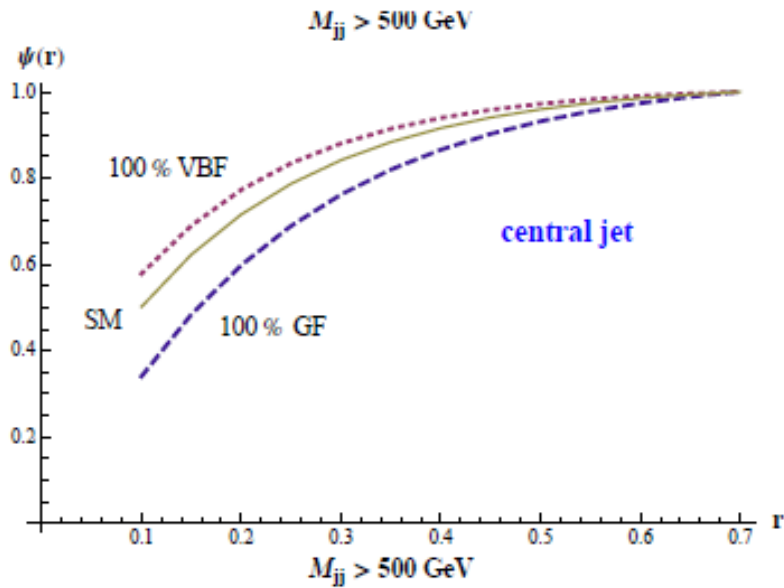
# We use the central jet



- Better reconstruction
- Better separation of JEPs



# Separation of profiles for different cuts



# Default Pythia tune cannot be relied upon to measure the jet profile

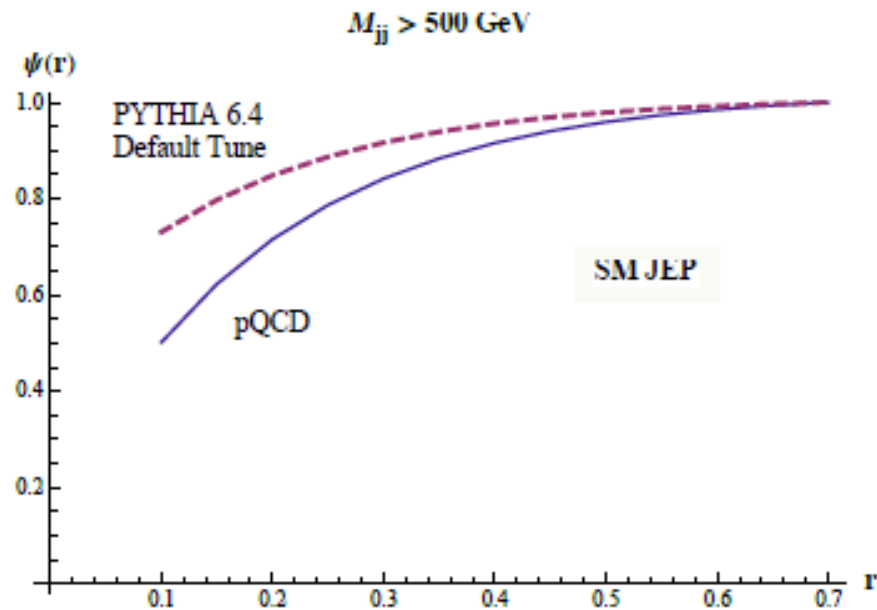


FIG. 7: Energy profile of the central jet for SM obtained by analyzing the jet substructure after Pythia v6.4 (default tune) showering, compared to the theoretical pQCD prediction using jet functions [11, 12].