

# Lattice developments for $\Delta M_{d,s}$

Elvira Gámiz (on behalf of Fermilab Lattice-MILC)

(with C. Bouchard and E. Freeland)



Universidad de Granada / CAFPE

- **9th International Workshop on the CKM Unitarity Triangle,**  
TIFR, Mumbai, Nov 28-Dec 2 2016 •

# Neutral $B$ mixing

- \* Particularly interesting process for indirect NP searches and constraining BSM theories.
- \*\* Tension between  $\Delta M_{s,d}$  and  $\varepsilon_K$ . See M. Blanke talk
- \*\* Main physical observables  $\Delta M_{s,d}$  measured at the subpercent level.

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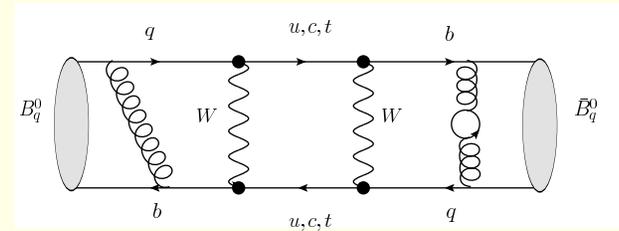
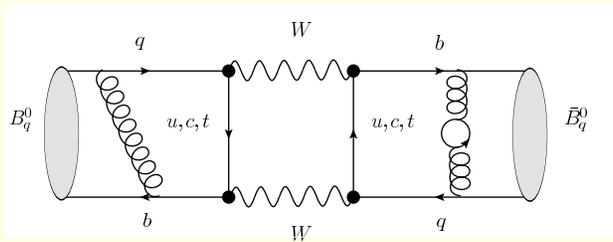
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- \* Using experimental measurements of  $\Delta M_{s,d}$  and theoretical determinations of the relevant hadronic matrix elements  
→ extract the CKM matrix elements  $|V_{ts}|$ ,  $|V_{td}|$ .

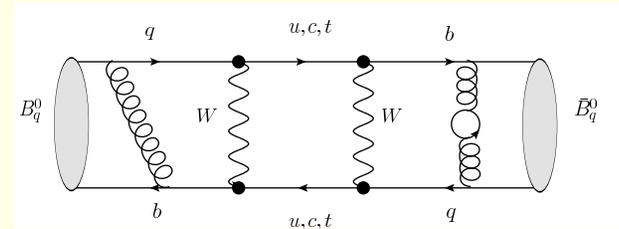
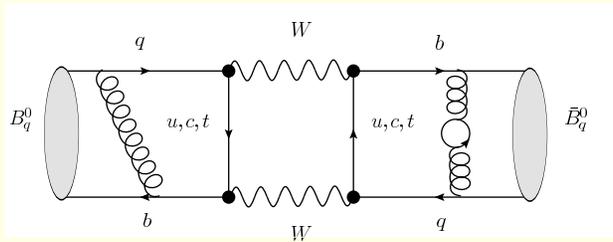
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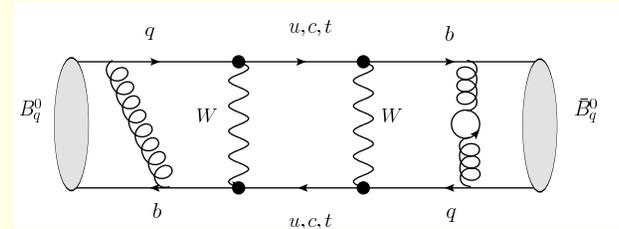
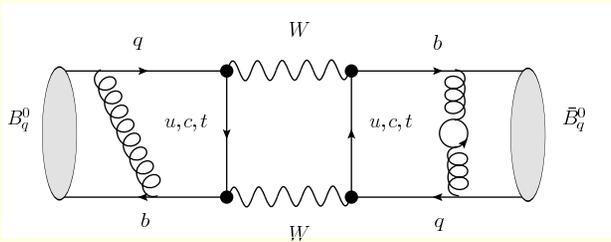


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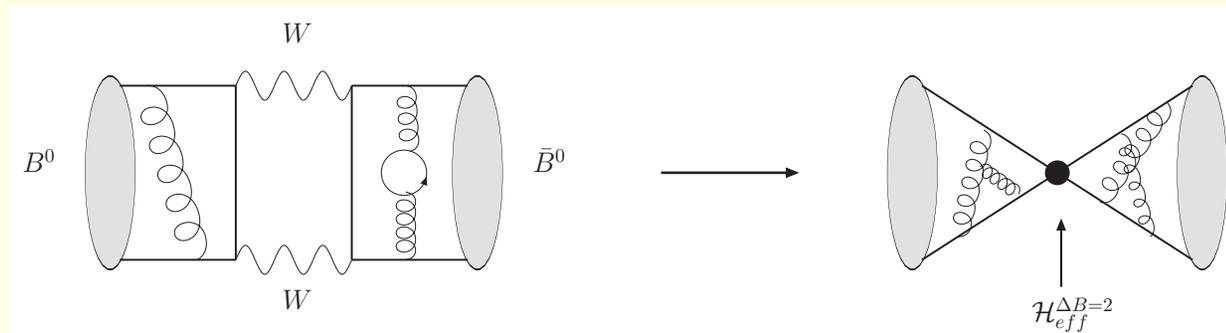
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- \* At tree level (flavour changing neutral currents)

Through a combination of **GIM mechanism** and **Cabibbo suppression**, the **top dominates** quark loop contributions



# Neutral $B$ mixing

And the mixing is described to a good approximation by the effective hamiltonian

$$\mathcal{H}_{eff}^{\Delta B=2} = \sum_{i=1}^5 C_i \mathcal{O}_i + \sum_{i=1}^3 \tilde{C}_i \tilde{\mathcal{O}}_i \quad \text{with}$$

$$\mathcal{O}_1^q = \left( \bar{b}^i \gamma^\nu (1 - \gamma_5) q^i \right) \left( \bar{b}^j \gamma^\nu (1 - \gamma_5) q^j \right) \quad \text{SM}$$

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$$\tilde{\mathcal{O}}_{1,2,3}^q = \mathcal{O}_{1,2,3}^q \text{ with the replacement } (1 \pm \gamma_5) \rightarrow (1 \mp \gamma_5)$$

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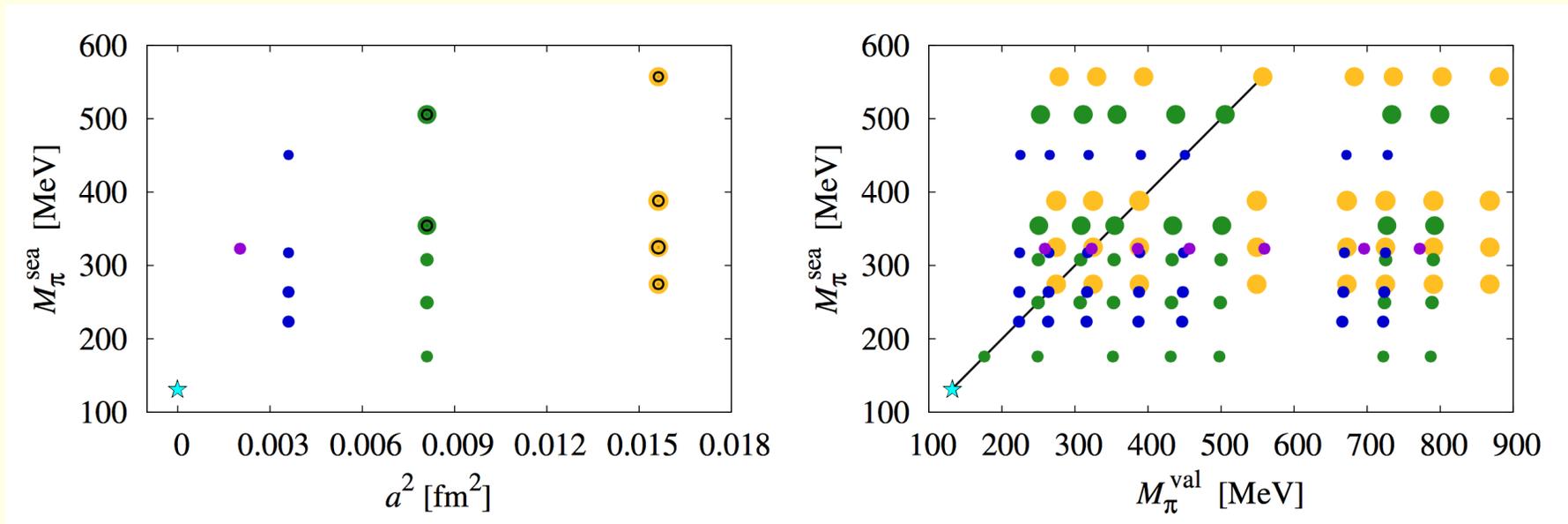
In **this talk**:

Calculation of the five hadronic matrix elements (and combinations of them) using three-flavour lattice QCD **FNAL-MILC 1602.03560** (SM prediction of  $\Delta M_{d,s}$  and  $\xi$ )

# 1.1 Simulation details

MILC  $N_f = 2 + 1$  asqtad ensembles

- \* 600-2000 gauge fields per ensemble
- \* pions as light as 177 MeV



## 1.2 Matching and renormalization

\* Mostly non-perturbative renormalization (mNPR).

$$\mathcal{O}_i = Z_{V_{bb}^4} Z_{V_{dd}^4} \rho_{ij} \mathcal{O}_j + \mathcal{O}(\alpha_s a, a^2)$$

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( $O_{1,2,3}$  mix under renormalization, as well as  $O_{4,5}$ )

## 1.3 Chiral-Continuum extrapolation

Extrapolate the lattice data to the continuum and infinite volume limits, and physical light quark masses in the Heavy Meson (HM)ChPT framework:

- \* Including dominant light quark discretization effects (NLO Staggered HMChPT) and NNLO ChPT analytic terms
- \* Gluon and light-quark discretization effects a la Symanzik
- \* Heavy-quark discretization effects (derived in HQET)
- \* Fine tuning  $m_b$ .
- \* Include higher order renormalization effects,  $\mathcal{O}(\alpha_s^2)$  in the fit.

$$F_i = F_i^{\text{logs}} + F_i^{\text{analytic}} + F_i^{\alpha_s a^2} + F_i^{\text{HQ disc.}} + F_i^{m_b \text{ tune}} + F_i^{\text{renor.}}$$

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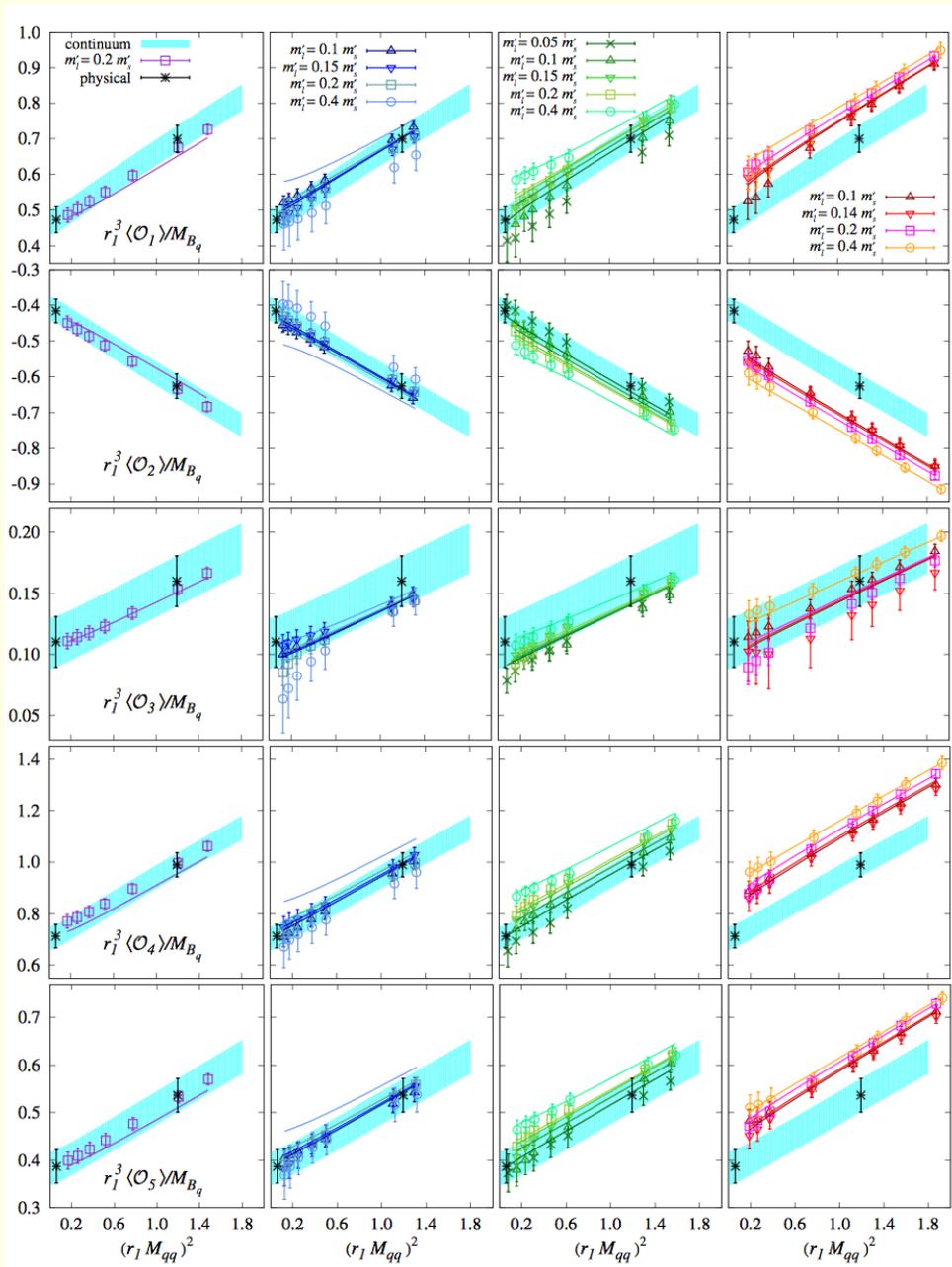
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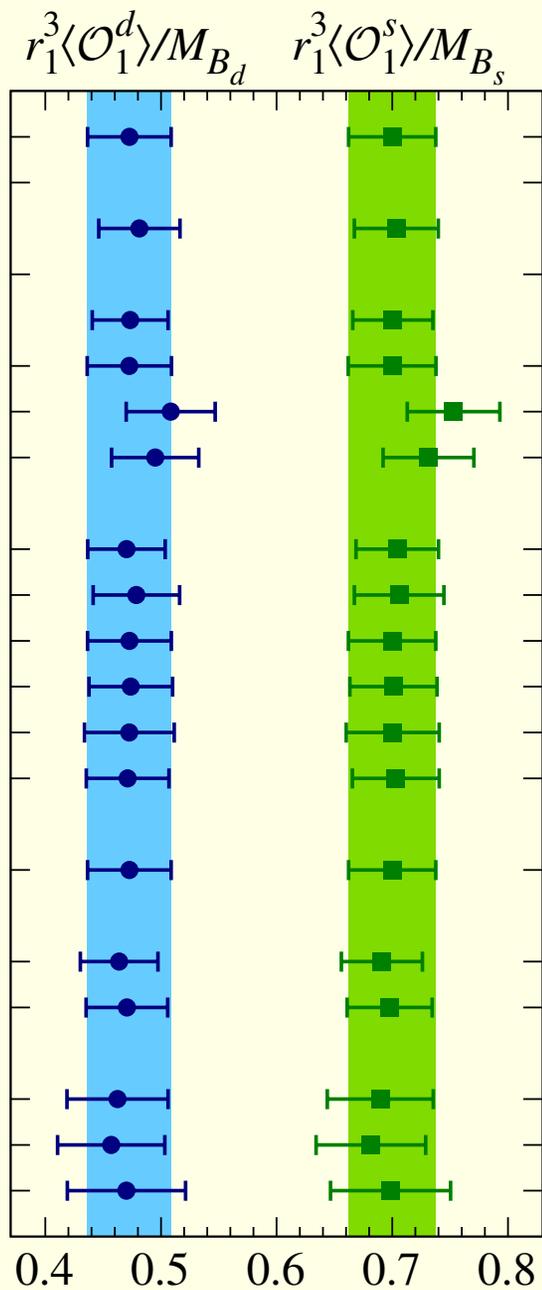


\*  $O_{1,2,3}$  and  $O_{4,5}$  also mix within ChPT.

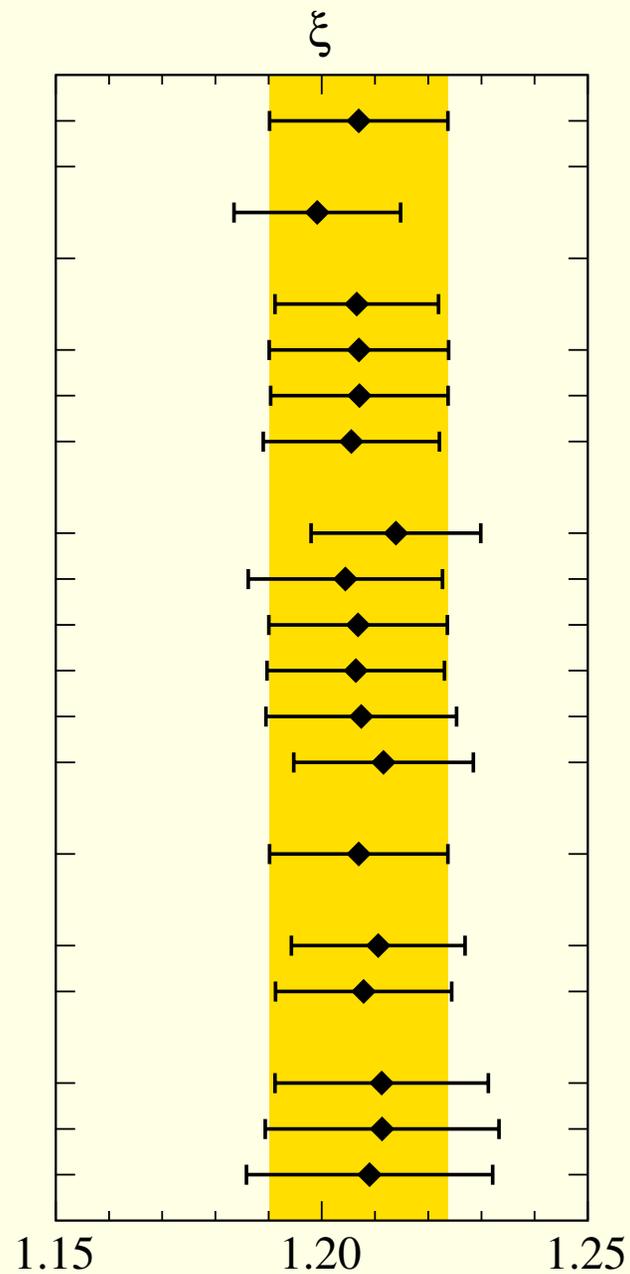
\* All operators are correlated via common gauge fields and valence quarks.

\* Perform a simultaneous (Bayesian) fit to all five operators.

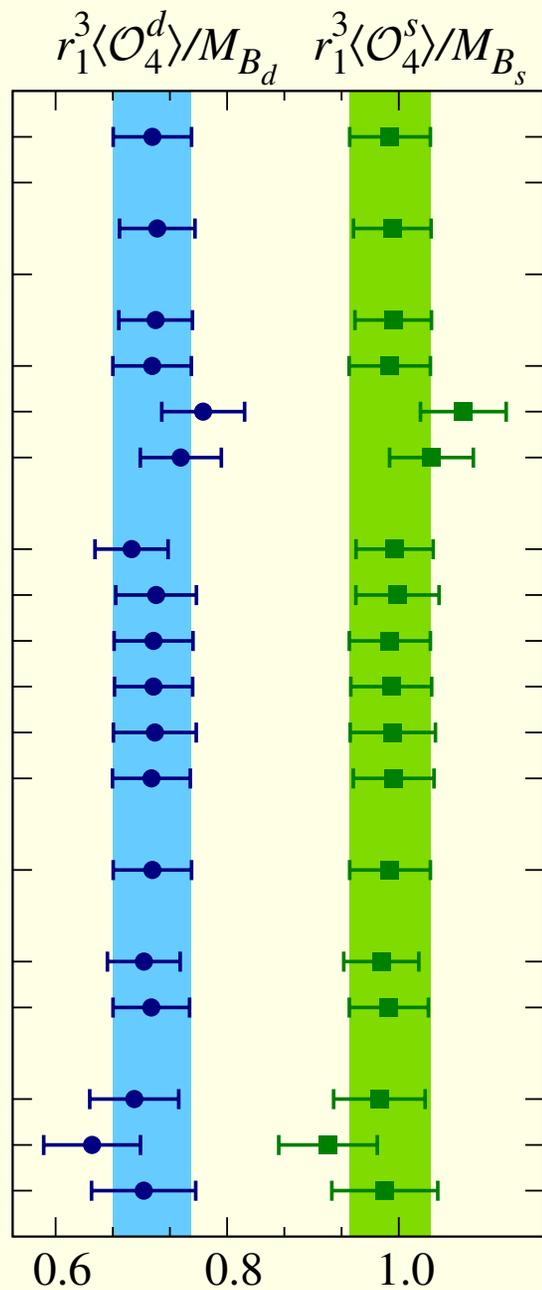
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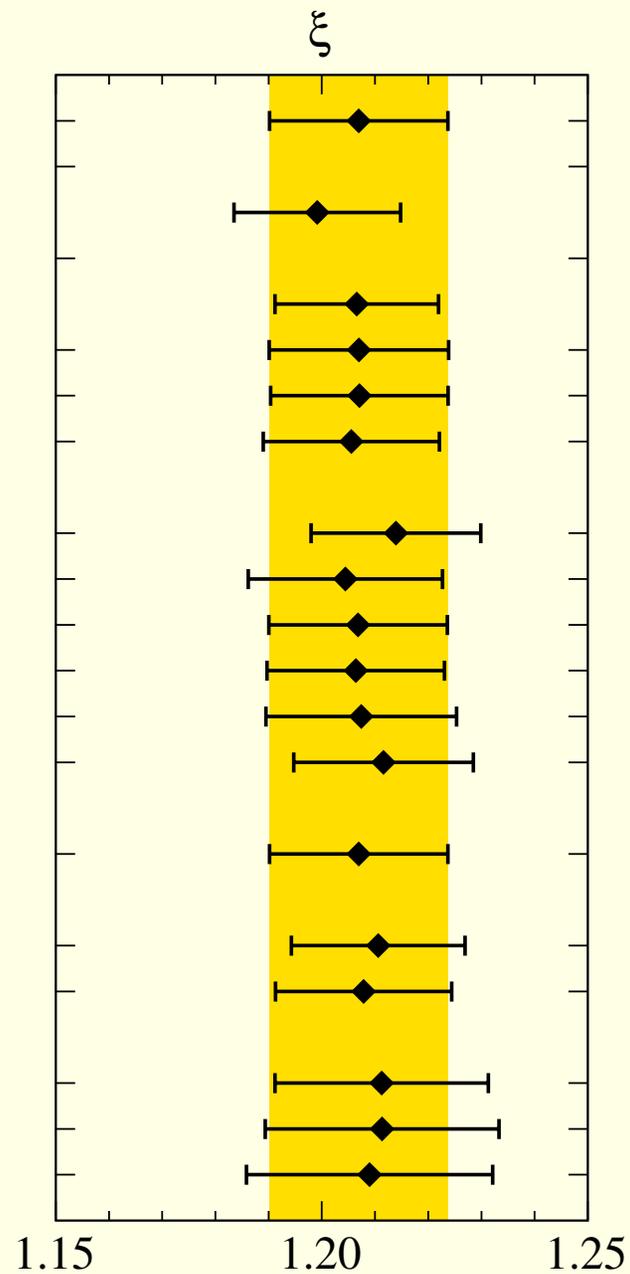
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- mNPR
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- PT<sub>P</sub> +  $\alpha_s^2$
- PT<sub>L</sub> +  $\alpha_s^2$
- NLO ( $m_q < 0.65m_s$ )
- N<sup>3</sup>LO
- LO × 2
- NLO × 2
- NNLO × 2
- no splitting
- generic  $O(\alpha_s a)$
- HQ  $O(\alpha_s a)$  only
- HQ  $O(\alpha_s a, a^2)$  only
- no  $a \approx 0.12$  fm
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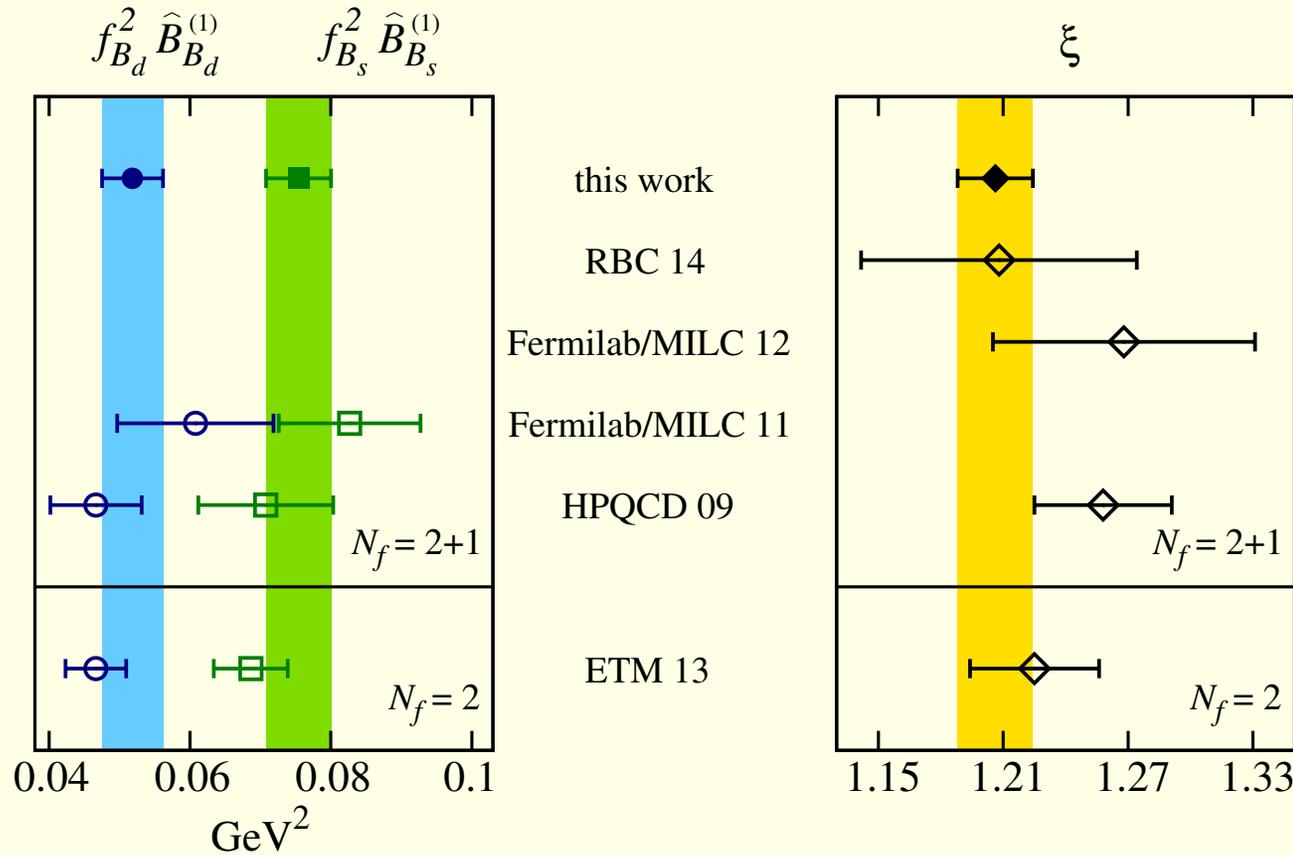


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## 2.1. Matrix elements relevant for SM $\Delta M_{s,d}$

In the SM,  $\Delta M_q \propto |V_{tq}^* V_{tb}|^2 f_{B_q}^2 \hat{B}_{B_q}^{(1)}$ , where  $\frac{8}{3} f_{B_q}^2 B_{B_q}^{(1)}(\mu) M_{B_q}^2 = \langle \mathcal{O}_1^q \rangle(\mu)$



**This work:** 1602.03560,  
**RBC 14:** 1406.6192,  
**Fermilab/MILC 12:**  
 1205.7013, **Fermilab/MILC**  
**11:** 1112.5642  
 (proceedings), **HPQCD 09:**  
 0902.1815, **ETM 13:**  
 1308.1851

In the  $SU(3)$ -breaking ratio  $\xi = \sqrt{\frac{f_{B_s}^2 \hat{B}_{B_s}^{(1)}}{f_{B_d}^2 \hat{B}_{B_d}^{(1)}}}$ , statistical and systematic uncertainties largely cancel (1.5% error dominated by statistics and HQ disc.)

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In the SM, using tree-level inputs for the CKM matrix elements **CKMfitter** and **Fermilab-MILC 1602.03560** results

$$f_{B_d} \sqrt{\hat{B}_{B_d}^{(1)}} = 227.7(9.5)(2.3) \text{ MeV} , f_{B_s} \sqrt{\hat{B}_{B_s}^{(1)}} = 274.6(8.4)(2.7) \text{ MeV} \quad ,$$
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we get

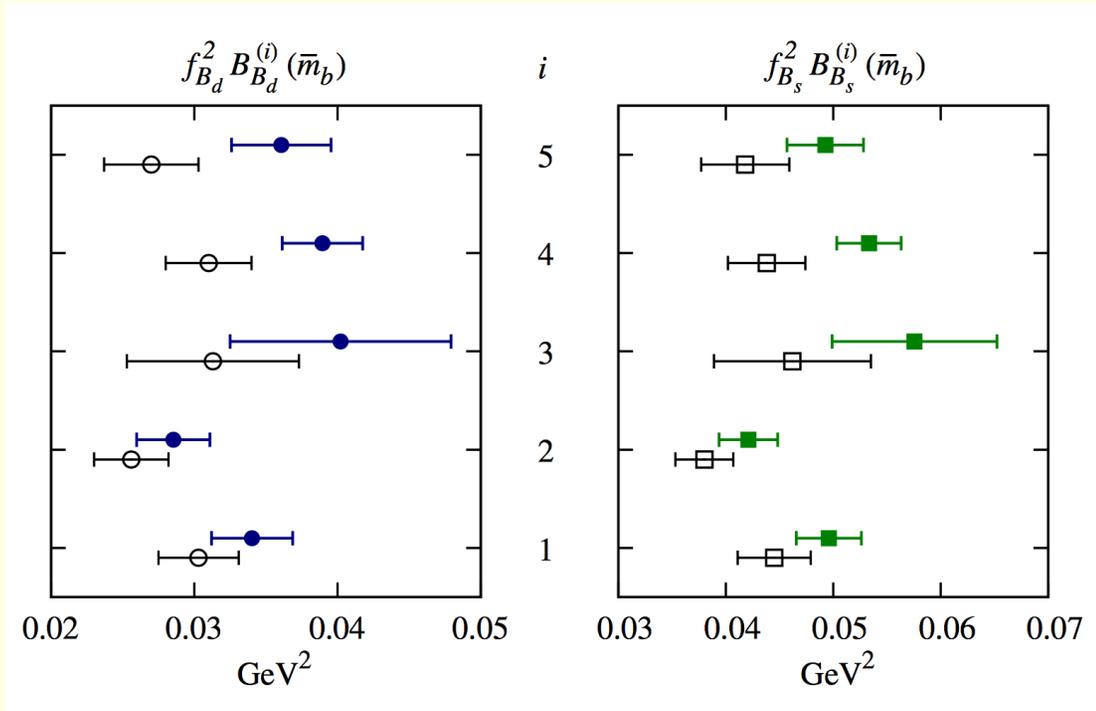
$$\begin{aligned} \Delta M_d^{SM} &= 0.630(53)(42)(5)(13) \text{ ps}^{-1} & \Delta M_d^{expt,HFAG} &= 0.5064(19) \text{ ps}^{-1} \\ \Delta M_s^{SM} &= 19.6(1.2)(1.0)(0.2)(0.4) \text{ ps}^{-1} & \Delta M_s^{expt,HFAG} &= 17.757(21) \text{ ps}^{-1} \\ (\Delta M_d / \Delta M_s)^{SM} &= 0.0321(10)(15)(0)(3) \text{ ps}^{-1} \end{aligned}$$

(where the errors are from lattice, CKM matrix elements, other inputs in SM expression, omission of charm quark on the sea, respectively)

\* These amount to tensions of  $2.1\sigma$ ,  $1.3\sigma$  and  $2.9\sigma$ , respectively.

## 2.2 Matrix elements relevant for BSM physics

Comparison with  $N_f = 2$  ETM collaboration 1308.1851 results



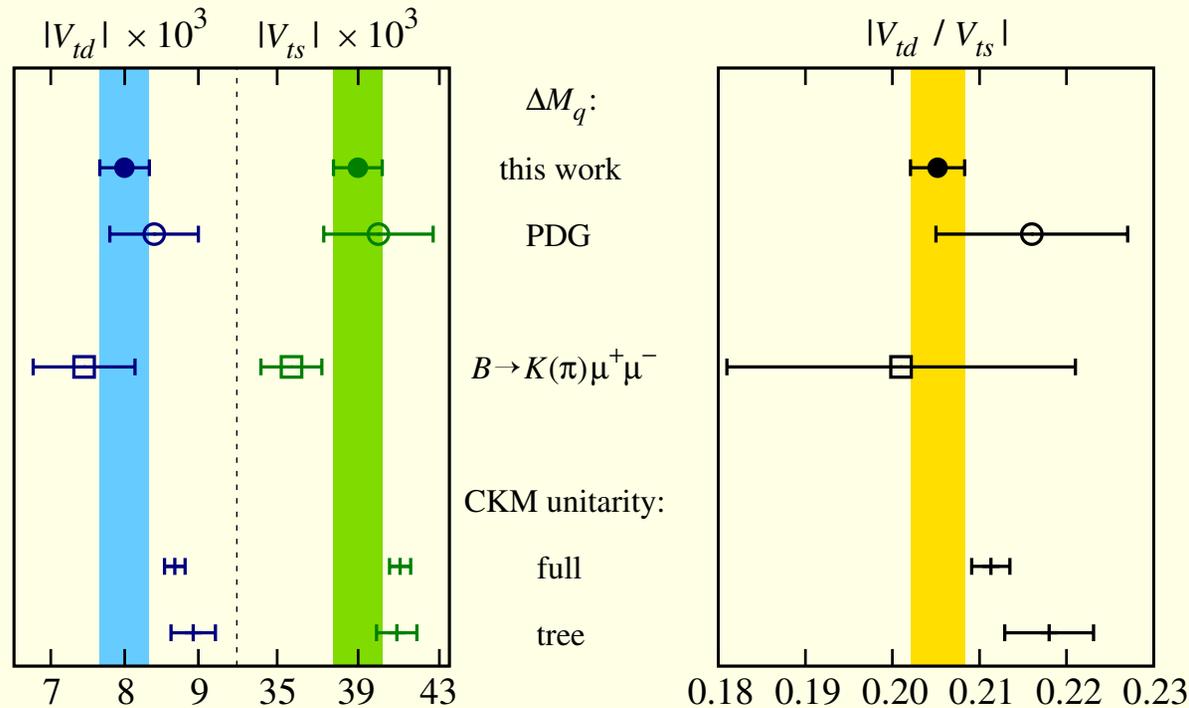
Open symbols: ETM

Full symbols: our results

\* Errors range from  $\sim 5 - 15\%$ , larger for  $B_d$  matrix elements.

## 2.3 Extraction of CKM matrix elements

Alternatively, use  $\Delta M_q^{expt}$  **HFAG 2014** and determine CKM factors

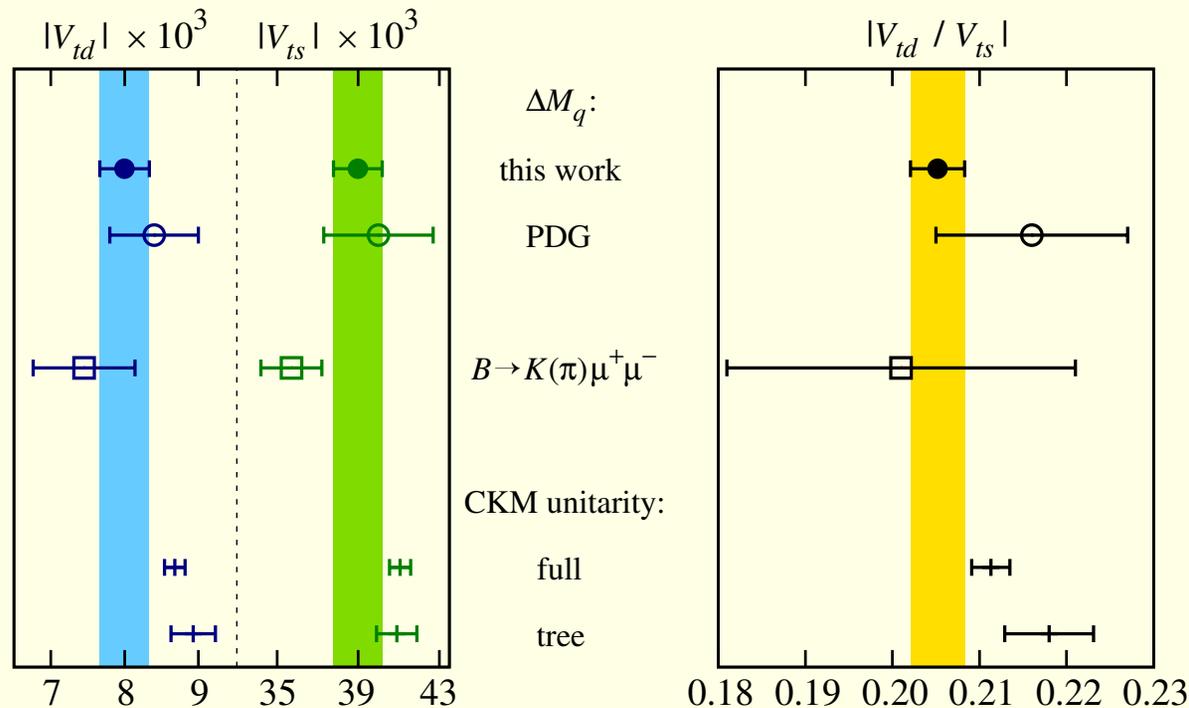


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Our results for  $|V_{td}|$ ,  $|V_{ts}|$  are  $2\sigma$ ,  $2.9\sigma$  below the CKM tree-fit results

\* Errors dominated by lattice mixing matrix elements

## 2.4 Bag parameters

Matrix elements of four fermion operators are often recast in terms of bag parameters:  $\langle \bar{B}_q | \mathcal{O}_i | B_q \rangle \propto f_{B_q}^2 B_{B_q}^{(i)}$

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\* Using  $f_B = 193.6(4.2)$  MeV,  $f_{B_s} = 228.6(3.8)$  MeV  $f_{B_s}/f_B = 1.187(15)$  from **Rosner, Stone, Van de Water**, PDG review, 1509.02220 and our results  $\rightarrow$  full set of bag parameters (in the SM and beyond) and correlations

\*\* For the SM RGI bag parameters we get

$$\hat{B}_{B_d}^{(1)} = 1.38(12)(6), \quad \hat{B}_{B_s}^{(1)} = 1.443(88)(48), \quad \frac{\hat{B}_{B_d}^{(1)}}{\hat{B}_{B_s}^{(1)}} = 1.033(31)(26)$$

(errors from matrix elements and decay constants respectively)

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**In progress:** Correlated calculation of decay constants  $\rightarrow$  decrease bag parameters errors.

## 2.5 Rare decays $B \rightarrow \mu^+ \mu^-$

Bag parameters  $\hat{B}_{B_{d,s}}$  describing B-meson mixing in the SM can be used for (indirect) theoretical predictions of  $\mathcal{B}(B \rightarrow \mu^+ \mu^-)$  **Buras** hep-ph/0303060, **Bobeth et al** 1311.0903

$$\left( \frac{\Gamma(B_q \rightarrow \mu^+ \mu^-)}{\Delta M_q} \right)^{\text{SM}} = \frac{3}{\pi^3} \frac{(G_F M_W m_\mu)^2}{\eta_{2B} S_0(x_t)} \frac{C_A^2(\mu_b)}{\hat{B}_{B_q}^{(1)}} \sqrt{1 - \frac{4m_\mu^2}{M_{B_q}^2}}$$

(with  $C_A(\mu_b)$  including NLO EW and NNLO QCD corrections)

**Herman, Misiak, Steinhauser** 1311.1347, **Bobeth, Gorbahn, Stamou**, 1311.1348

## 2.5 Rare decays $B \rightarrow \mu^+ \mu^-$

Bag parameters  $\hat{B}_{B_{d,s}}$  describing **B-meson** mixing in the **SM** can be used for (indirect) theoretical predictions of  $\mathcal{B}(B \rightarrow \mu^+ \mu^-)$  **Buras** hep-ph/0303060, **Bobeth et al** 1311.0903

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- \* Using our  $N_f = 2 + 1 \hat{B}_{B_{d,s}}$ , including the effects of a non-vanishing  $\Delta\Gamma_s$  to compute the **time-averaged branching fractions**  $\bar{\mathcal{B}}$  measured in experiment  
 $(\bar{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)^{\text{SM}} = \tau_{H_s} \Gamma(B_s \rightarrow \mu^+ \mu^-)^{\text{SM}}, \bar{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-) = \mathcal{B}(B_d \rightarrow \mu^+ \mu^-))$   
 and the experimental  $\Delta M_q$  **HFAG 2014**

$$\bar{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-)^{\text{SM}} = 9.06(85)(4)(16) \cdot 10^{-11}$$

$$\bar{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)^{\text{SM}} = 3.22(22)(0)(6) \cdot 10^{-9}$$

$$\left( \frac{\bar{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-)}{\bar{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)} \right)^{\text{SM}} = 0.02786(109)(12)(19)$$

(with errors coming from bag parameters, experimental  $\Delta M_q$  and others, respectively)

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\* SM predictions using

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To be compared with the experimental averages from **LHCb** and **CMS 1411.4413**

$$\overline{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-)^{\text{exp}} = 3.9^{(+1.6)}_{(-1.4)} \times 10^{-10}$$

$$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)^{\text{exp}} = 2.8^{(+0.7)}_{(-0.6)} \times 10^{-9}$$

$$\left( \frac{\overline{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-)}{\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)} \right)^{\text{exp}} = 0.14^{(+0.08)}_{(-0.06)}$$

$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)$  agrees with experiment,  $\overline{\mathcal{B}}(B_d \rightarrow \mu^+ \mu^-)$  is  $2\sigma$  above (symmetrizing exp. errors), and the ratio  $1.6\sigma$  below.

### 3. Matrix elements contributing to $\Delta\Gamma_{d,s}$

At NLO in the heavy quark expansion  $\Delta\Gamma_q^{\text{SM}}$  depends on

$$\langle \mathcal{O}_1 \rangle, \langle \mathcal{O}_3 \rangle, \langle R_0 \rangle, \langle R_{1,2,3} \rangle$$

- \* With **FNAL/MILC 1602.03560**:  $\langle \mathcal{O}_1 \rangle$  and  $\langle \mathcal{O}_3 \rangle$  known with 6% and 13% error.
- \*  $R_0 = \mathcal{O}_1 + \alpha_1 \mathcal{O}_2 + \frac{\alpha_2}{2} \mathcal{O}_3$  and  $R_1 = \frac{m_q}{m_b} \mathcal{O}_4$  calculated in **FNAL/MILC 1602.03560**
- \* VSA estimates for dimension-7 operators  $R_2$  and  $R_3$  (50% error)

$$\langle R_2 \rangle = \frac{1}{m_b^2} \left( \bar{b}^i \overleftarrow{D}_\alpha \gamma^\nu (1 - \gamma_5) D^\alpha q^i \right) \left( \bar{b}^j \gamma^\nu (1 - \gamma_5) q^j \right)$$

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Dominant  $\Delta\Gamma_s^{\text{SM}}$  uncertainties:  $\langle\mathcal{O}_1\rangle$  (14%  $\rightarrow$  6%),  $\langle R_2\rangle$  (15%) and renormalization scale (8%) **Artuso, Borissov, Lenz 1511.09466**

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**On-going:**  $N_f = 2 + 1 + 1$  **HPQCD** calculation of dimension-7 operators  $R_2$  (and  $R_3$ )  
(see **M. Wingate** talk at Lattice 2016)

- \* **Goal:** Reduce error in  $\langle R_2 \rangle$  to 25%  $\implies \Delta\Gamma_s$  error 19%  $\rightarrow$  14%

## 4. Conclusions and outlook

- \* First three-flavor results for full set of  $B_{s,d}$  mixing matrix elements
  - \*\* All source of systematic uncertainty controlled.
  - \*\* Most precise determination (1.6% error) of  $\xi$  and  $\langle \mathcal{O}_1^{d,s} \rangle$ .

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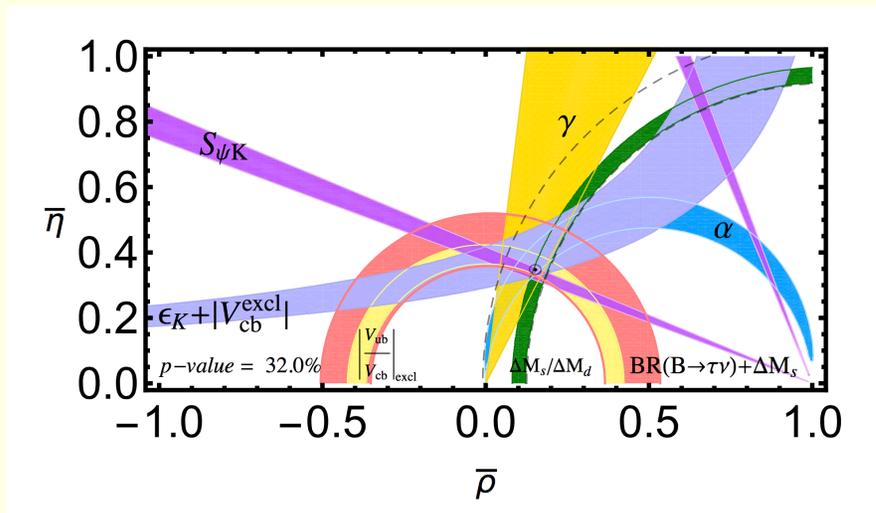
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- \* Using Fermilab-MILC results for  $B_{s,d}$ -meson mixing parameters 1602.03560,  $V_{cb}$  1403.0635, 1503.07237 and  $V_{ub}$  1503.07839



Compatible with SM at  $p = 0.32$ ,  
but still ample room for BSM  
flavor-changing neutral currents

## 4. Conclusions and outlook

**On-going** Fermlab-MILC

Combined analysis of matrix elements and decay constants →  
correlations → reduction of errors for bag parameters

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Use MILC  $N_f = 2 + 1 + 1$  HISQ configurations and HISQ valence quarks

- \* Physical light quark masses → reduce (eliminate) chiral extr. error
- \* Eliminate charm quark sea error
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**On-going:** another Lattice collaborations

- \*  $N_f = 2 + 1 + 1$  HPQCD with HISQ light quarks and non-relativistic  $b$ .  
Preliminary results 1411.6989

