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## Higher dimensional HQET parameters

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[Ayesh Gunawardana, GP *to appear*]

# Outline

- Motivation
- A little bit of history
- HQET (and NRQCD) operators at dimension 8 and above
- Conclusions

# Motivation

## Motivation

- Inclusive semileptonic  $B$  decays and  $Q_{7\gamma} - \bar{Q}_{7\gamma}$  contribution to  $B \rightarrow X_s \gamma$  can be described by local OPE

$$\Gamma = \sum_{n=0}^{\infty} \frac{1}{m_b^n} \sum_k c_{k,n} \langle O_{k,n} \rangle$$

- $c_{k,n}$  are Wilson coefficients: perturbative
- $\langle O_{k,n} \rangle$  are matrix elements of HQET operators: non-perturbative  
Often called HQET parameters  
 $O_{k,n} \sim \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h \cdot T_{\mu_1 \dots \mu_n \lambda}$ ,  $T$  is a Lorentz tensor
- $|V_{cb}|$  extraction from inclusive  $B$  decays  
uses dimension 7 and 8 HQET operators  
[Gambino, Healey, Turczyk PLB **763**, 60 (2016)]

# Motivation

- $|V_{cb}|$  extraction from inclusive  $B$  decays  
uses dimension 7 and 8 HQET operators  
[Gambino, Healey, Turczyk PLB **763**, 60 (2016)]
- Question:  
Are these *all* possible HQET operators at dimension 7,8?
- Question: Can we go higher?
- Question:  
Since HQET and NRQCD are related,  
[Manohar PRD **56**, 230 (1997), Gunawardana, GP *to appear*]  
What are the NRQCD operators?
- Question:  
What can we learn about the structure of effective field theories?  
[“Higher dimension operators in the SM EFT”  
Henning, Lu, Melia, Murayama, arXiv:1512.03433]

# A little bit of history

## Prehistory

$$D_t = \frac{\partial}{\partial t} + ieA^0, \quad \mathbf{D} = \nabla - ie\mathbf{A}$$

- Schrödinger equation:  $iD_t + \frac{\mathbf{D}^2}{2M}$
- Hydrogen Fine Structure:
  - Spin-Orbit:  $\boldsymbol{\sigma} \cdot \mathbf{B}$
  - Relativistic correction:  $\mathbf{D}^4$
  - Darwin term:  $\boldsymbol{\partial} \cdot \mathbf{E}$
- Organize operators in Lagrangian form
- The  $\text{dim}=5,6$  were given in [Caswell, Lepage PLB **167**, 437 (1986)]

$$\mathcal{L}_{\text{NRQED}}^{\text{dim}=5,6} = \psi^\dagger \left\{ iD_t + \frac{\mathbf{D}^2}{2M} + \frac{\mathbf{D}^4}{8M^3} + c_F g \frac{\boldsymbol{\sigma} \cdot \mathbf{B}}{2M} + c_D g \frac{[\boldsymbol{\partial} \cdot \mathbf{E}]}{8M^2} \right. \\ \left. + i c_S g \frac{\boldsymbol{\sigma} \cdot (\mathbf{D} \times \mathbf{E} - \mathbf{E} \times \mathbf{D})}{8M^2} + c_W 1 g \frac{\{\mathbf{D}^2, \boldsymbol{\sigma} \cdot \mathbf{B}\}}{8M^3} \right\} \psi$$

- First systematic discussion of HQET parameters  
[Mannel, PRD 50, 428 (1994)]
- Between HQET fields  $\bar{h} \dots h$  the Dirac basis reduces to  $\{1, \sigma\} = \{1, s^\lambda\}$  with  $v \cdot s = 0$
- $\bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h$  is the general operator  
Since  $iv \cdot Dh = 0$ 
  - $v_{\mu_1} \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h = 0$
  - $v_{\mu_n} \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h = 0$
  - $v_\lambda \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h = 0$
- Dimension 3:  $\bar{h} h$
- Dimension 4:  $\bar{h} iD^\mu h = 0$
- Dimension 5: Two operators  $\bar{h} iD^{\mu_1} iD^{\mu_2} h, \bar{h} iD^{\mu_1} iD^{\mu_2} s^\lambda h$
- Dimension 6: Two operators  $\bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} h, \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} s^\lambda h$
- Dimension 7: Wrong decomposition



- First systematic discussion of HQET parameters

[Mannel, PRD 50, 428 (1994)]

- Dimension 5:

$$\langle B(v) | \bar{h}(iD_\alpha)(iD_\beta)h | B(v) \rangle = 2M_H [g_{\alpha\beta} - v_\alpha v_\beta] \frac{1}{3} \lambda_1$$

$$\langle B(v) | \bar{h}(iD_\alpha)(iD_\beta) s_\lambda h | B(v) \rangle = 2M_H d_H i \epsilon_{\nu\alpha\beta\lambda} v^\nu \frac{1}{6} \lambda_2$$

- Dimension 6:

$$\langle B(v) | \bar{h}(iD_\alpha)(iD_\mu)(iD_\beta) h_\nu | B(v) \rangle = 2M_H [g_{\alpha\beta} - v_\alpha v_\beta] v_\mu \frac{1}{3} \rho_1$$

$$\langle B(v) | \bar{h}(iD_\alpha)(iD_\mu)(iD_\beta) s_\lambda h | B(v) \rangle = 2M_H d_H i \epsilon_{\nu\alpha\beta\lambda} v^\nu v_\mu \frac{1}{6} \rho_2.$$

- Dimension 7: Wrong decomposition

# HQET-NRQCD (NRQED) correspondence

- HQET: Heavy Quark Effective Theory and NRQCD: Non Relativistic Quantum Chromo Dynamics differ in the kinetic term

$$\mathcal{L}_{HQET}^{kinetic} = \bar{h} i v \cdot D h$$

$$\mathcal{L}_{NRQCD}^{kinetic} = \psi^\dagger i D_t \psi + \frac{\mathbf{D}^2}{2M} \psi$$

and power counting

- Correspondence between HQET and NRQCD (NRQED) operators [Manohar PRD **56**, 230 (1997); Gunawardana, GP *to appear*]

	NRQED (1920's-1980's)	HQET(1990's)
Dimension 5	$\mathbf{D}^2$	$(iD_\perp)^2$
	$\boldsymbol{\sigma} \cdot \mathbf{B}$	$(iD_\perp^\mu)(iD_\perp^\nu)(-i\sigma^{\mu\nu})$
Dimension 6	$[\partial \cdot \mathbf{E}]$	$(iD_{\perp\mu})(iv \cdot D)(iD_\perp^\mu)$
	$\boldsymbol{\sigma} \cdot (\mathbf{D} \times \mathbf{E} - \mathbf{E} \times \mathbf{D})$	$(iD_\perp^\mu)(iv \cdot D)(iD_\perp^\nu)(-i\sigma^{\mu\nu})$

where  $D_\perp^\mu = D^\mu - v \cdot D v^\mu$

- The  $\text{dim}=7$  NRQCD Lagrangian was given in [Manohar PRD **56**, 230 (1997)]

$$\mathcal{L}_{\text{NRQCD}}^{\text{dim}=7} = \psi^\dagger \left\{ \frac{\mathbf{D}^4}{8M^3} + ic_M g \frac{\{\mathbf{D}^i, [\partial \times \mathbf{B}]^i\}}{8M^3} + c_{A1} g^2 \frac{\mathbf{B}^2 - \mathbf{E}^2}{8M^3} - c_{A2} g^2 \frac{\mathbf{E}^2}{16M^3} + \right. \\ \left. + c_{p'p} g \frac{\boldsymbol{\sigma} \cdot \mathbf{D} \mathbf{B} \cdot \mathbf{D} + \mathbf{D} \cdot \mathbf{B} \boldsymbol{\sigma} \cdot \mathbf{D}}{8M^3} \right. \\ \left. c_{W1} g \frac{\{\mathbf{D}^2, \boldsymbol{\sigma} \cdot \mathbf{B}\}}{8M^3} - c_{W2} g \frac{D^i \boldsymbol{\sigma} \cdot \mathbf{B} D^i}{4M^3} + \right. \\ \left. + ic_{B1} g^2 \frac{\boldsymbol{\sigma} \cdot (\mathbf{B} \times \mathbf{B} - \mathbf{E} \times \mathbf{E})}{8M^3} - ic_{B2} g^2 \frac{\boldsymbol{\sigma} \cdot (\mathbf{E} \times \mathbf{E})}{8M^3} \right\} \psi$$

- Last line is non-zero only for NRQCD
- The operators are listed, but not derived
- NRQCD dimension 7:
  - 4 SI (spin-independent) operators
  - 5 SD (spin-dependent) operators

## Dimension 7 HQET operators

- How many HQET operators do we have at dimension 7?
  - 1) [Mannel, PRD 50, 428 (1994)]      2 SI operators      5 SD operators
  - 2) [Manohar PRD **56**, 230 (1997)]      4 SI operators      5 SD operators
  - 3) [Dassinger, Mannel, Turczyk  
JHEP **0703**, 087 (2007)]      3 SI operators      2 SD operators
  - 4) [Mannel, Turczyk, Uraltsev  
JHEP **1011**, 109 (2010)]      4 SI operators      5 SD operators
- HQET-NRQCD correspondance implies 4 SI and 5 SD operators
- No systematic derivation in either source

# Dimension 8 HQET operators

- What about dimension 8?  
[Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)]  
7 SI operators and 11 SD operators  
No systematic derivation
- The **dim=8** NRQED Lagrangian was given in  
[Hill, Lee, GP, Solon, PRD **87** 053017 (2013)]

$$\mathcal{L}_{\text{NRQED}}^{\text{dim}=8} = \psi^\dagger \left\{ c_{X1} g^2 \frac{[\mathbf{D}^2, \mathbf{D} \cdot \mathbf{E} + \mathbf{E} \cdot \mathbf{D}]}{M^4} + c_{X2} g^2 \frac{\{\mathbf{D}^2, [\partial \cdot \mathbf{E}]\}}{M^4} \right. \\
+ c_{X3} g^2 \frac{[\partial^2 \partial \cdot \mathbf{E}]}{M^4} + i c_{X4} g^2 \frac{\{\mathbf{D}^i, [\mathbf{E} \times \mathbf{B}]^i\}}{M^4} \\
+ i c_{X5} g^2 \frac{D^i \boldsymbol{\sigma} \cdot (\mathbf{D} \times \mathbf{E} - \mathbf{E} \times \mathbf{D}) D^i}{M^4} + i c_{X6} g^2 \frac{\epsilon^{ijk} \sigma^i D^j [\partial \cdot \mathbf{E}] D^k}{M^4} \\
+ c_{X7} g^2 \frac{\boldsymbol{\sigma} \cdot \mathbf{B} [\partial \cdot \mathbf{E}]}{M^4} + c_{X8} g^2 \frac{[\mathbf{E} \cdot \partial \boldsymbol{\sigma} \cdot \mathbf{B}]}{M^4} + c_{X9} g^2 \frac{[\mathbf{B} \cdot \partial \boldsymbol{\sigma} \cdot \mathbf{E}]}{M^4} \\
\left. + c_{X10} g^2 \frac{[E^i \boldsymbol{\sigma} \cdot \partial B^i]}{M^4} + c_{X11} g^2 \frac{[B^i \boldsymbol{\sigma} \cdot \partial E^i]}{M^4} + c_{X12} g^2 \frac{\boldsymbol{\sigma} \cdot \mathbf{E} \times [\partial_t \mathbf{E} - \partial \times \mathbf{B}]}{M^4} \right\} \psi$$

4 SI operators and 8 SD operators

Missing operators are presumably NRQCD operators

No systematic derivation

## Lesson from History

- Finding all of the HQET operators at a given dimension is not easy
- Is there a systematic way to do that?

# HQET (and NRQCD) operators at dimension 8 and above

[Ayesh Gunawardana, GP *to appear*]

## General method

- We consider matrix elements of the form

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle$$

- We use
  - Parity and Time reversal symmetry
  - Hermitian conjugation
  - Four dimensions

To decompose them in terms of tensors

- $v^{\mu_i}$
- $g^{\mu_i \mu_j}$
- $\epsilon^{\alpha \beta \rho \sigma}$



## General method: PT symmetry

- We consider matrix elements of the form

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle$$

- Parity and Time reversal are symmetries of HQET

In particular under PT:

$$- p = (p^0, \vec{p}) \xrightarrow{PT} (p^0, \vec{p}) = p \Rightarrow v = p/m \xrightarrow{PT} v$$

$$- iD^\mu \xrightarrow{PT} iD^\mu$$

$$- \bar{h} h \xrightarrow{PT} \bar{h} h$$

$$- \bar{h} s^\lambda h \xrightarrow{PT} -\bar{h} s^\lambda h$$

- Since  $T$  is anti-linear

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle \stackrel{PT}{=} \langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle^*$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle \stackrel{PT}{=} -\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle^*$$

- SI matrix elements are real, SD matrix elements are complex

## General method: Hermitian conjugation

- We consider matrix elements of the form

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle$$

- $\bar{h}h$ ,  $\bar{h}s^\lambda h$ ,  $iD^\mu$  are hermitian using Hermitian conjugation

$$\begin{aligned} \langle B | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h | B \rangle &= \langle B | \left( \bar{h} iD^{\mu_1} \dots iD^{\mu_n} (s^\lambda) h \right)^\dagger | B \rangle^* \\ &= \langle B | \bar{h} iD^{\mu_n} \dots iD^{\mu_1} (s^\lambda) h | B \rangle^* \end{aligned}$$

- Combining with the  $PT$  constraints  
Under inversion of the indices:
  - SI matrix elements are symmetric
  - SD matrix elements are anti-symmetric

## General method: Tensor decomposition

- We consider matrix elements of the form

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle$$

- $B$  is a pseudo-scalar  $\Rightarrow$  matrix element can only depend on  $v^{\mu_i}$ ,  $g^{\mu_i\mu_j}$ , and  $\epsilon^{\alpha\beta\rho\sigma}$

- Alternatively following

[Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)]

Define  $\Pi^{\mu\nu} = g^{\mu\nu} - v^\mu v^\nu$

For the standard choice of  $v = (1, 0, 0, 0)$ :  $\Pi^{00} = 0$  and  $\Pi^{ij} = -\delta^{ij}$

- Since the indices in  $\epsilon^{\alpha\beta\rho\sigma}$  cannot be independent of  $v^\mu$  replace  $\epsilon^{\alpha\beta\rho\sigma}$  by  $\epsilon^{\alpha\beta\rho\sigma} v_\sigma$

- Matrix element depend on  $v^{\mu_i}$ ,  $\Pi^{\mu_i\mu_j}$ , and  $\epsilon^{\alpha\beta\rho\sigma} v_\sigma$

## General method: Four dimensions

- We consider matrix elements of the form

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} h | B(v) \rangle$$

$$\langle B(v) | \bar{h} iD^{\mu_1} \dots iD^{\mu_n} s^\lambda h | B(v) \rangle$$

- Four dimensions  $\Rightarrow$  only four independent directions
- Not all tensors with more than four indices are independent
- Example: for dimension 7 SD HQET operators  
need  $\Pi^{\mu\nu} \epsilon^{\alpha\beta\rho\sigma} v_\sigma$ : three indices are the same  
Tensors obtained by permuting indices are not linearly independent
- Example: for dimension 11 SI HQET operators  
need  $\Pi^{\mu_1\mu_2} \Pi^{\mu_3\mu_4} \Pi^{\mu_5\mu_6} \Pi^{\mu_7\mu_8}$ : four indices are the same

## Results: Dimension 7 HQET operators

- Using the general method

$$\frac{1}{2M_B} \langle B | \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} iD^{\mu_4} h | B \rangle = a_{12}^{(7)} \Pi^{\mu_1 \mu_2} \Pi^{\mu_3 \mu_4} + a_{13}^{(7)} \Pi^{\mu_1 \mu_3} \Pi^{\mu_2 \mu_4} +$$

$$+ a_{14}^{(7)} \Pi^{\mu_1 \mu_4} \Pi^{\mu_2 \mu_3} + b^{(7)} \Pi^{\mu_1 \mu_4} v^{\mu_2} v^{\mu_3}$$

$$\frac{1}{2M_B} \langle B | \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} iD^{\mu_4} s^\lambda h | B \rangle = i\tilde{b}^{(7)} v^{\mu_2} v^{\mu_3} \epsilon^{\mu_1 \mu_4 \lambda \rho} v_\rho +$$

$$+ i\tilde{a}_{12}^{(7)} \left( \Pi^{\mu_1 \mu_2} \epsilon^{\mu_3 \mu_4 \lambda \rho} v_\rho - \Pi^{\mu_4 \mu_3} \epsilon^{\mu_2 \mu_1 \lambda \rho} v_\rho \right) +$$

$$+ i\tilde{a}_{13}^{(7)} \left( \Pi^{\mu_1 \mu_3} \epsilon^{\mu_2 \mu_4 \lambda \rho} v_\rho - \Pi^{\mu_4 \mu_2} \epsilon^{\mu_3 \mu_1 \lambda \rho} v_\rho \right) +$$

$$+ i\tilde{a}_{14}^{(7)} \Pi^{\mu_1 \mu_4} \epsilon^{\mu_2 \mu_3 \lambda \rho} v_\rho + i\tilde{a}_{23}^{(7)} \Pi^{\mu_2 \mu_3} \epsilon^{\mu_1 \mu_4 \lambda \rho} v_\rho$$

- How many HQET operators do we have at dimension 7?  
4 SI operators and 5 SD operators confirms

[Manohar PRD **56**, 230 (1997)]

[Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)]

## Results: Dimension 8 HQET operators

- Using the general method

$$\begin{aligned} \frac{1}{2M_B} \langle B | \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} iD^{\mu_4} iD^{\mu_5} h | B \rangle &= a_{12}^{(8)} (\Pi^{\mu_1 \mu_2} \Pi^{\mu_3 \mu_5} v^{\mu_4} + \Pi^{\mu_1 \mu_3} \Pi^{\mu_5 \mu_4} v^{\mu_2}) + \\ a_{13}^{(8)} (\Pi^{\mu_1 \mu_3} \Pi^{\mu_2 \mu_5} v^{\mu_4} + \Pi^{\mu_5 \mu_3} \Pi^{\mu_4 \mu_1} v^{\mu_2}) &+ a_{15}^{(8)} (\Pi^{\mu_1 \mu_5} \Pi^{\mu_3 \mu_4} v^{\mu_2} + \Pi^{\mu_1 \mu_5} \Pi^{\mu_3 \mu_2} v^{\mu_4}) + \\ b_{12}^{(8)} \Pi^{\mu_1 \mu_2} \Pi^{\mu_4 \mu_5} v^{\mu_3} + b_{14}^{(8)} \Pi^{\mu_1 \mu_4} \Pi^{\mu_2 \mu_5} v^{\mu_3} &+ b_{15}^{(8)} \Pi^{\mu_1 \mu_5} \Pi^{\mu_2 \mu_4} v^{\mu_3} + \\ c^{(8)} \Pi^{\mu_1 \mu_5} v^{\mu_2} v^{\mu_3} v^{\mu_4} \end{aligned}$$

$$\begin{aligned} \frac{1}{2M_B} \langle B | \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} iD^{\mu_4} iD^{\mu_5} s^\lambda h | B \rangle &= \\ i\tilde{a}_{12}^{(8)} (v^{\mu_3} \Pi^{\mu_1 \mu_2} \epsilon^{\mu_4 \mu_5 \lambda \rho} v_\rho - v^{\mu_3} \Pi^{\mu_4 \mu_5} \epsilon^{\mu_2 \mu_1 \lambda \rho} v_\rho) &+ i\tilde{a}_{14}^{(8)} (v^{\mu_3} \Pi^{\mu_1 \mu_4} \epsilon^{\mu_2 \mu_5 \lambda \rho} v_\rho - v^{\mu_3} \Pi^{\mu_5 \mu_2} \epsilon^{\mu_4 \mu_1 \lambda \rho} v_\rho) \\ + i\tilde{a}_{15}^{(8)} v^{\mu_3} \Pi^{\mu_1 \mu_5} \epsilon^{\mu_2 \mu_4 \lambda \rho} v_\rho + i\tilde{a}_{24}^{(8)} v^{\mu_3} \Pi^{\mu_2 \mu_4} \epsilon^{\mu_1 \mu_5 \lambda \rho} v_\rho &+ \\ + i\tilde{b}_{13}^{(8)} (v^{\mu_2} \Pi^{\mu_1 \mu_3} \epsilon^{\mu_4 \mu_5 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_5 \mu_3} \epsilon^{\mu_2 \mu_1 \lambda \rho} v_\rho) &+ i\tilde{b}_{14}^{(8)} (v^{\mu_2} \Pi^{\mu_1 \mu_4} \epsilon^{\mu_3 \mu_5 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_5 \mu_2} \epsilon^{\mu_3 \mu_1 \lambda \rho} v_\rho) \\ + i\tilde{b}_{15}^{(8)} (v^{\mu_2} \Pi^{\mu_1 \mu_5} \epsilon^{\mu_3 \mu_4 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_1 \mu_5} \epsilon^{\mu_3 \mu_2 \lambda \rho} v_\rho) &+ i\tilde{b}_{34}^{(8)} (v^{\mu_2} \Pi^{\mu_3 \mu_4} \epsilon^{\mu_1 \mu_5 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_3 \mu_2} \epsilon^{\mu_1 \mu_5 \lambda \rho} v_\rho) \\ + i\tilde{b}_{35}^{(8)} (v^{\mu_2} \Pi^{\mu_3 \mu_5} \epsilon^{\mu_1 \mu_4 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_3 \mu_1} \epsilon^{\mu_5 \mu_2 \lambda \rho} v_\rho) &+ i\tilde{b}_{45}^{(8)} (v^{\mu_2} \Pi^{\mu_4 \mu_5} \epsilon^{\mu_1 \mu_3 \lambda \rho} v_\rho - v^{\mu_4} \Pi^{\mu_2 \mu_1} \epsilon^{\mu_5 \mu_3 \lambda \rho} v_\rho) \\ + i\tilde{c}^{(8)} v^{\mu_2} v^{\mu_3} v^{\mu_4} \epsilon^{\mu_1 \mu_5 \lambda \rho} v_\rho. \end{aligned}$$

- How many HQET operators do we have at dimension 8?

7 SI operators and 11 SD operators, confirms

[Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)]

## Results: Dimension 9 HQET operators

- Using the general method:

24 possible SI Dimension 9 HQET operators (New!)

$$\begin{aligned}
 \langle B | \bar{h} iD^{\mu_1} iD^{\mu_2} iD^{\mu_3} iD^{\mu_4} iD^{\mu_5} iD^{\mu_6} h | B \rangle = & \\
 & \Pi^{\mu_1\mu_2} \Pi^{\mu_3\mu_4} \Pi^{\mu_5\mu_6} + (\Pi^{\mu_1\mu_2} \Pi^{\mu_3\mu_5} \Pi^{\mu_4\mu_6} + \Pi^{\mu_1\mu_3} \Pi^{\mu_2\mu_4} \Pi^{\mu_5\mu_6}) + \\
 & (\Pi^{\mu_1\mu_2} \Pi^{\mu_3\mu_6} \Pi^{\mu_4\mu_5} + \Pi^{\mu_1\mu_4} \Pi^{\mu_2\mu_3} \Pi^{\mu_5\mu_6}) + \Pi^{\mu_1\mu_3} \Pi^{\mu_2\mu_5} \Pi^{\mu_4\mu_6} + \\
 & (\Pi^{\mu_1\mu_3} \Pi^{\mu_2\mu_6} \Pi^{\mu_4\mu_5} + \Pi^{\mu_1\mu_5} \Pi^{\mu_2\mu_3} \Pi^{\mu_4\mu_6}) + \Pi^{\mu_1\mu_4} \Pi^{\mu_2\mu_5} \Pi^{\mu_3\mu_6} + \\
 & (\Pi^{\mu_1\mu_4} \Pi^{\mu_2\mu_6} \Pi^{\mu_3\mu_5} + \Pi^{\mu_1\mu_5} \Pi^{\mu_2\mu_4} \Pi^{\mu_3\mu_6}) + \Pi^{\mu_1\mu_5} \Pi^{\mu_2\mu_6} \Pi^{\mu_3\mu_4} + \\
 & \Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_3} \Pi^{\mu_4\mu_5} + \Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_4} \Pi^{\mu_3\mu_5} + \\
 & \Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_5} \Pi^{\mu_3\mu_4} + (\Pi^{\mu_1\mu_2} \Pi^{\mu_4\mu_6} v^{\mu_3} v^{\mu_5} + \Pi^{\mu_1\mu_3} \Pi^{\mu_5\mu_6} v^{\mu_2} v^{\mu_4}) + \\
 & (\Pi^{\mu_1\mu_2} \Pi^{\mu_4\mu_6} v^{\mu_3} v^{\mu_5} + \Pi^{\mu_1\mu_3} \Pi^{\mu_5\mu_6} v^{\mu_2} v^{\mu_4}) + \Pi^{\mu_1\mu_2} \Pi^{\mu_5\mu_6} v^{\mu_3} v^{\mu_4} + \\
 & (\Pi^{\mu_1\mu_2} \Pi^{\mu_4\mu_6} v^{\mu_3} v^{\mu_5} + \Pi^{\mu_1\mu_3} \Pi^{\mu_5\mu_6} v^{\mu_2} v^{\mu_4}) + \Pi^{\mu_1\mu_3} \Pi^{\mu_4\mu_6} v^{\mu_2} v^{\mu_5} + \\
 & \Pi^{\mu_1\mu_4} \Pi^{\mu_3\mu_6} v^{\mu_2} v^{\mu_5} + (\Pi^{\mu_1\mu_4} \Pi^{\mu_2\mu_6} v^{\mu_3} v^{\mu_5} + \Pi^{\mu_1\mu_5} \Pi^{\mu_3\mu_6} v^{\mu_2} v^{\mu_4}) + \\
 & \Pi^{\mu_1\mu_5} \Pi^{\mu_2\mu_6} v^{\mu_3} v^{\mu_4} + (\Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_3} v^{\mu_4} v^{\mu_5} + \Pi^{\mu_1\mu_6} \Pi^{\mu_4\mu_5} v^{\mu_2} v^{\mu_3}) + \\
 & (\Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_4} v^{\mu_3} v^{\mu_5} + \Pi^{\mu_1\mu_6} \Pi^{\mu_3\mu_5} v^{\mu_2} v^{\mu_4}) + \Pi^{\mu_1\mu_6} \Pi^{\mu_2\mu_5} v^{\mu_3} v^{\mu_4} + \\
 & \Pi^{\mu_1\mu_6} \Pi^{\mu_3\mu_4} v^{\mu_2} v^{\mu_5} + \Pi^{\mu_1\mu_6} v^{\mu_2} v^{\mu_3} v^{\mu_4} v^{\mu_5}
 \end{aligned}$$

# Conclusions



# Conclusions

- We presented a general method to construct HQET operators:  
using tensor decomposition of HQET matrix elements  
Several applications:
  - Easily relate different bases
    - Dimension 7: Manohar '97 to Mannel-Turczyk-Uraltsev '10
    - Dimension 8: Mannel-Turczyk-Uraltsev '10 to Hill, Lee, GP, Solon '12
  - SI dimension 9 HQET operators
  - Moments of the leading order shape function  
up to and including dimension 9 HQET operators
- We will present the dimension 8 NRQCD Lagrangian

# Conclusions

- [Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)] lists dimension 7 and 8 HQET operators
- Question: Are these *all* possible HQET operators at dimension 7,8?
  - Answer: Yes
- Question: Can we go higher?
  - Answer: Yes, we presented the general method for that
- Question: Since HQET and NRQCD are related, what are the NRQCD operators?
  - Answer will be given in [Gunawardana, GP *to appear*]
- Question:  
What can we learn about the structure of effective field theories?  
[Henning, Lu, Melia, Murayama, arXiv:1512.03433]
  - Answer: Simpler than we think...