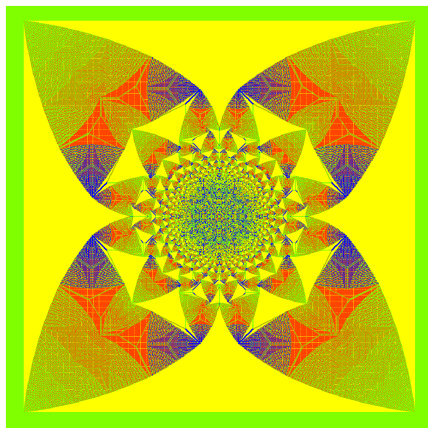
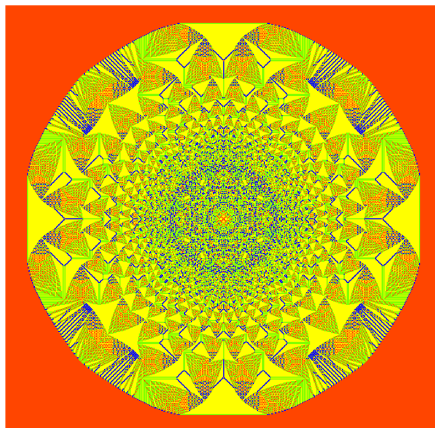


# Pattern formation in sandpile models of self-organized criticality

Tridib Sadhu

28 Dec 2010

# The patterns



# Generating the patterns

- ▶ Non-negative integer height  $z_i$  on an infinite square lattice
- ▶ If  $z_i \geq 4$ , **Unstable**.

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$$z_i \rightarrow z_i - 4$$

- ▶ Iterate until stable .

The relaxed height distribution forms deterministic complex patterns.

# The patterns

$N = 4 \times 10^4$ . The color code: Red=0, Blue=1, Green=2, Yellow=3.

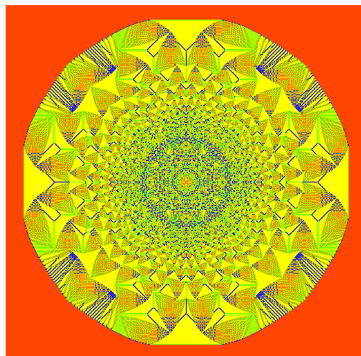


Figure: Background all  $z_i = 0$

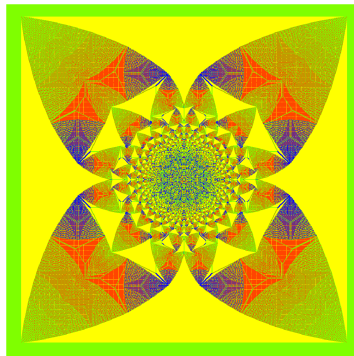


Figure: Background all  $z_i = 2$



- ▶ Motivation
- ▶ Characterization of the patterns
- ▶ Robustness to external noise
- ▶ Possible connection to some interesting mathematics

# Motivation

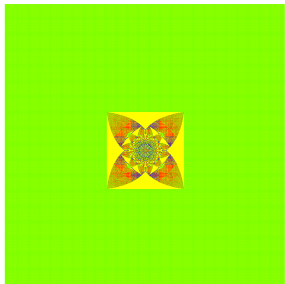


Figure:  $N = 4 \times 10^4$

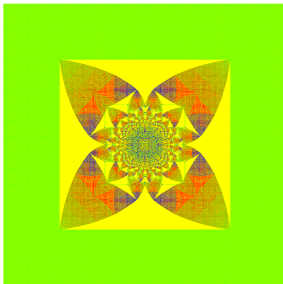


Figure:  $N = 2 \times 10^5$

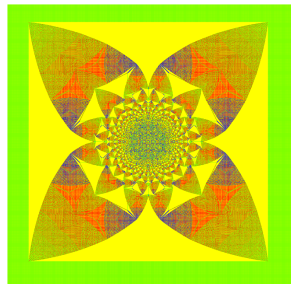


Figure:  $N = 4 \times 10^5$

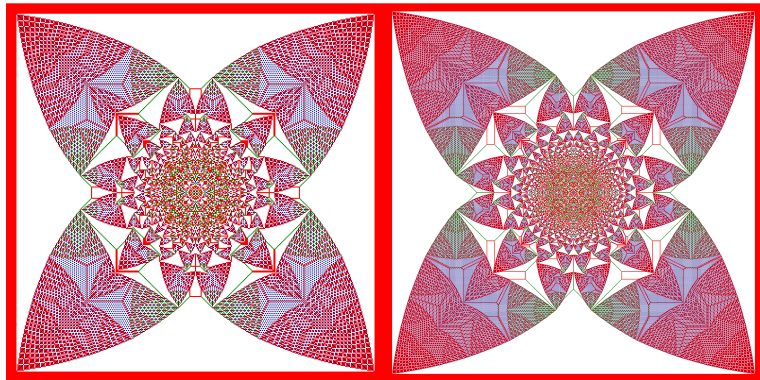
Diameter  $\sim \sqrt{N}$ .

# Proportionate growth

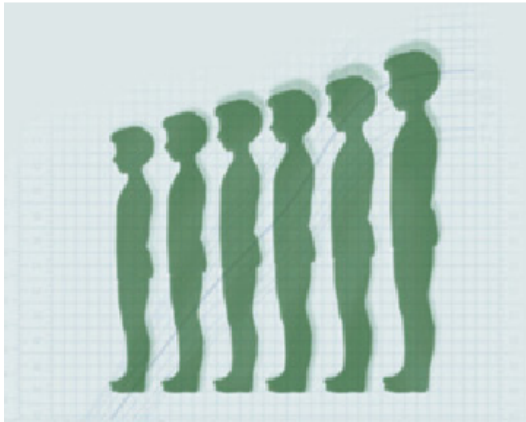
Color Code: 0 1 2 3

N=60k Color Code: 0 1 2 3

N=250k



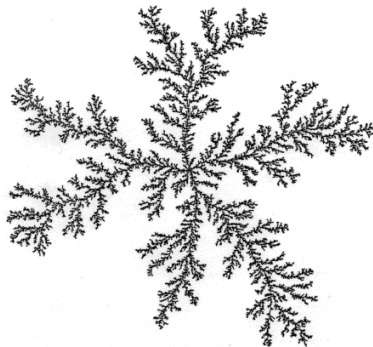
## Proportionate growth



**Figure:** Different body parts in animals grow roughly at the same rate.

# Motivation . . . contd.

- ▶ Proportionate growth requires regulation, and/or communication between different parts.
- ▶ Most existing models of growth in physics literature **DLA**, **KPZ growth**, **Invasion percolation**, etc does not have this property.
- ▶ Extra symmetries and robustness



- ▶ Emergence of **complex structures from simple local rules**, e.g. Fractals
- ▶ Sandpile patterns are complex, yet simpler to analytically characterize.
- ▶ Exact characterization of the pattern involves some interesting mathematics
  - ▶ Discrete analytic functions
  - ▶ Tropical polynomials

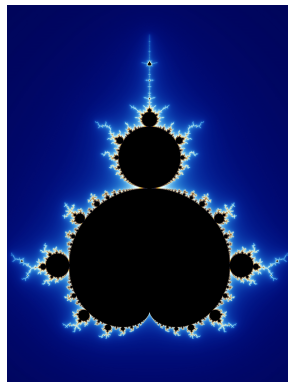


Figure: Mandelbrot set

# Characterizing the pattern

- ▶ Proportionate growth + Diameter  $\sim \sqrt{N}$

⇒ Describe in reduced coordinates

$$\xi = x/\sqrt{N}, \quad \eta = y/\sqrt{N}$$

- ▶ Characterize pattern in terms of density of heights  $\rho(\xi, \eta)$   
= the height averaged over an area  $\delta\xi\delta\eta$  around  $(\xi, \eta)$ , with  $1/\sqrt{N} \ll \delta\xi \ll 1$  and  $1/\sqrt{N} \ll \delta\eta \ll 1$ .

# Characterizing the pattern

Let  $T(x, y) = \#$  of toppling at  $(x, y)$ .

$$\sum' T(x', y') - 4T(x, y) = \Delta z(x, y) - N\delta_{x,0}\delta_{y,0} \quad (1)$$

Define

$$\phi(\xi, \eta) = \lim_{N \rightarrow \infty} \frac{T(x, y)}{N}$$

Then

$$\nabla^2 \phi(\xi, \eta) = \Delta \rho(\xi, \eta) - \delta(\xi)\delta(\eta),$$

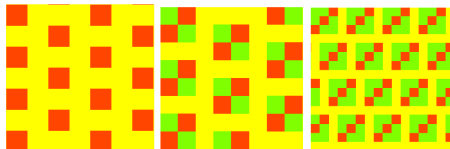
where  $\Delta \rho$  is the change in density.

The complete specification of  $\phi$  determines the patterns.



# Determining $\phi$

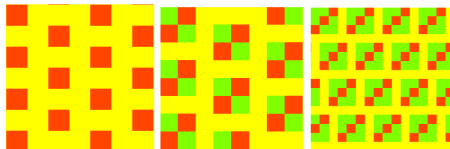
- ▶ Patches with periodic heights.



▶ pattern

# Determining $\phi$

- ▶ Patches with periodic heights.

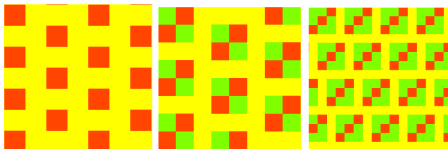


▶ pattern

- ▶  $\rho(\xi, \eta)$  is constant within a patch.

# Determining $\phi$

- ▶ Patches with periodic heights.



▶ pattern

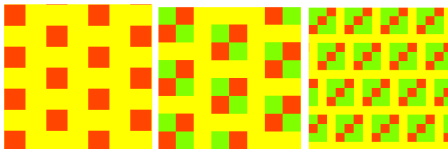
- ▶  $\rho(\xi, \eta)$  is constant within a patch.
- ▶ Lemma:

$\phi$  is a quadratic function of  $\xi, \eta$  in each patch.

▶ Proof

# Determining $\phi$

- ▶ Patches with periodic heights.



▶ pattern

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- ▶ Lemma:

$\phi$  is a quadratic function of  $\xi, \eta$  in each patch.

▶ Proof

- ▶ Continuity of  $\phi$  and its first derivatives along the patch boundaries imposes constraints.

Solve the constraints and determine  $\phi$ .

# Simpler pattern

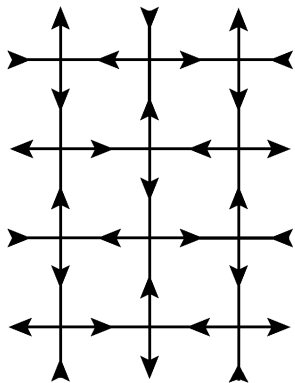


Figure: F-lattice with  $z_c = 2$

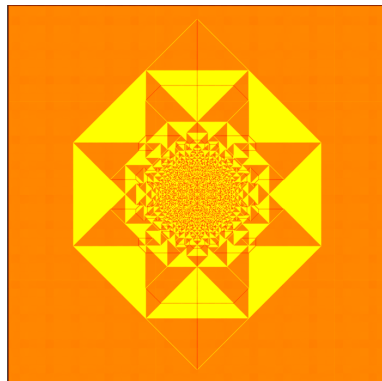


Figure:  $N = 2 \times 10^5$  on checkerboard background of 1 and 0 heights

# Adjacency graph

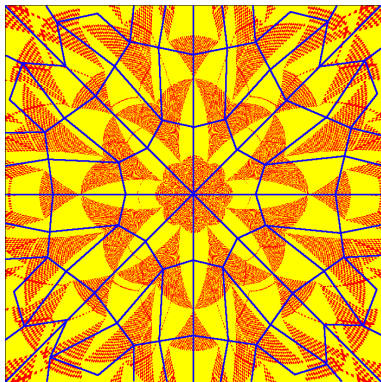


Figure: Adjacency graph

▶ 1/z2 picture

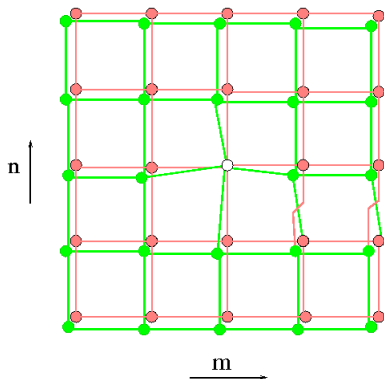


Figure: Representation as a square lattice on two sheeted Riemann surface

# Quantitative characterization

- ▶ Label patches using  $(m, n)$ .

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- ▶ Label patches using  $(m, n)$ .
- ▶ Potential in a dense patch  $(m, n)$ ,

$$\phi(\xi, \eta) = \frac{1}{8}(m+1)\xi^2 + \frac{1}{4}n\xi\eta + \frac{1}{8}(1-m)\eta^2 + d\xi + e\eta + f$$

In a light patch

$$\phi(\xi, \eta) = \frac{1}{8}m\xi^2 + \frac{1}{4}n\xi\eta - \frac{1}{8}m\eta^2 + d_{m,n}\xi + e_{m,n}\eta + f_{m,n}$$



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In a light patch

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- ▶ Continuity of  $\phi$  and its first derivatives along the patch boundaries imply,  $d_{m,n}$  and  $e_{m,n}$  follows

$$\psi_{m+1,n+1} + \psi_{m+1,n-1} + \psi_{m-1,n+1} + \psi_{m-1,n-1} - 4\psi_{m,n} = 0,$$

Boundary condition:

- ▶  $d(0, 0) + ie(0, 0) = 0$
- ▶ For large  $|m + in|$ ,

$$d(m, n) + ie(m, n) \simeq \pm \frac{1}{\sqrt{2\pi}} \sqrt{m + in}$$

Solve this set of linear equations numerically on a large grid.

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Solve this set of linear equations numerically on a large grid.

The pattern has exact eight fold rotational symmetry. ▶ Aside

# Multiple sites of addition

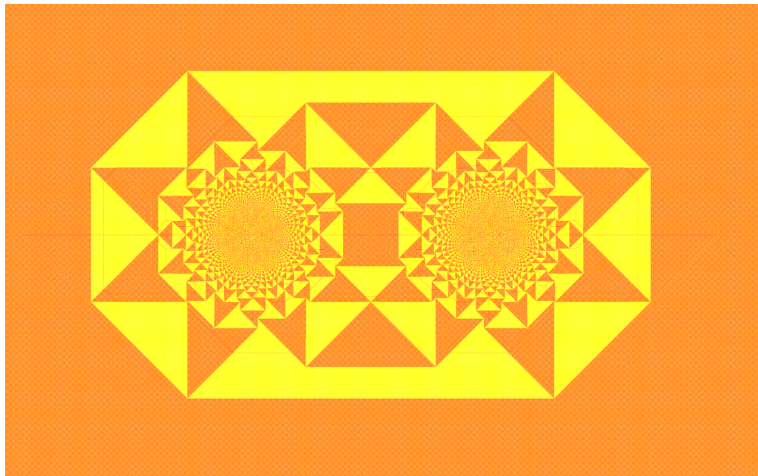


Figure:  $N = 10^5$  added each at sites  $(-400, 0)$  and  $(400, 0)$

# Adjacency graph

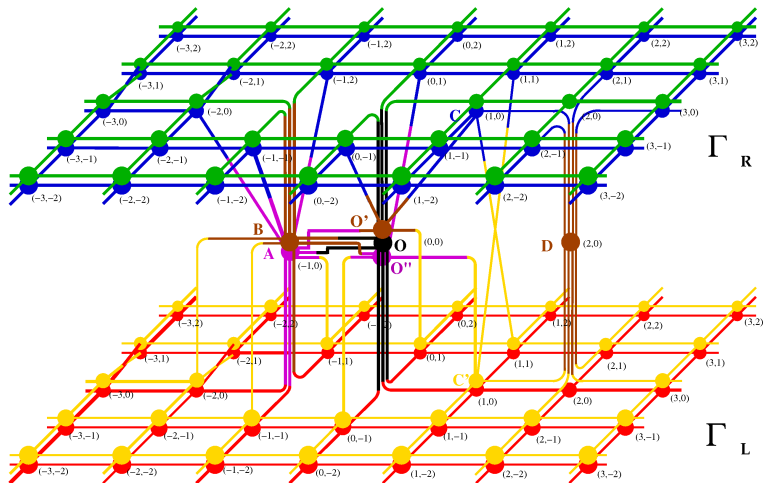


Figure: Adjacency graph

# Line of sink sites

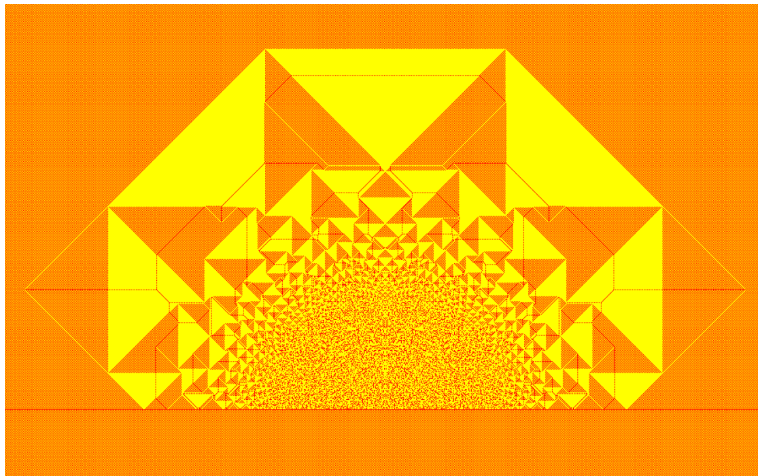


Figure:  $N = 10^5$  added at site  $(0, 1)$  with sink sites along the x-axis.

# Adjacency graph

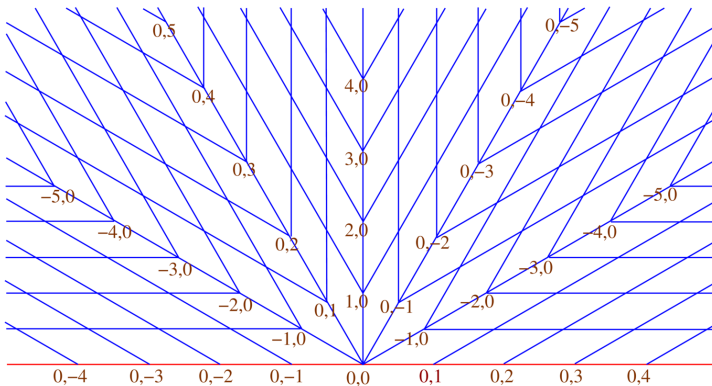


Figure: Adjacency graph

"Diameter" of the pattern

$$\Lambda \sim N^\alpha$$

In absence of sink sites  $\alpha = 1/2$ .

For the line sink:

- ▶ # un-absorbed particles  $N_r \sim \Lambda^2$
- ▶ # absorbed particles  $N_a \sim \Lambda^2 \int_{1/\Lambda}^1 d\xi \frac{\partial}{\partial \eta} \phi \Big|_{\eta=0}$
- ▶ Close to the sink line

$$\phi \sim \frac{\cos(\theta)}{r} \Rightarrow N_a \sim \Lambda^3$$

- ▶ Hence,

$$\boxed{C_1 \Lambda^3 + C_2 \Lambda^2 = N} \Rightarrow \Lambda \sim N^{1/3} \text{ For large } N.$$



- ▶ The equation gives correction to scaling.
- ▶ **Unexpected accuracy:** For  $C_1 = 0.1853$ , and  $C_2 = 0.528$  the solution of this equation differs from the actual  $\Lambda(N)$  by at most 1, for  $100 < N < 3 \times 10^6$

## Other sink geometries

- ▶ For a wedge of  $\theta$ ,  $\Lambda \sim N^\alpha$ , with  $\alpha = 1/(2 + \pi/\theta)$ .
- ▶ For a point sink adjacent to the site of addition  $\Lambda \sim \sqrt{N/\log N}$ .
- ▶ Generalizable to higher dimensions

# Fast growing sandpiles

- ▶ If the initial background density is low enough everywhere,

$$\Lambda \sim N^{1/2}$$

- ▶ If many sites have large height

$$\Lambda = \infty \quad \text{for finite } N$$

- ▶ For an in-between set of periodic backgrounds

$$\Lambda \sim N^\alpha \quad \text{for } 1/2 < \alpha \leq 1$$

**Lemma:** The potential function

$$\phi(\xi, \eta) = \lim_{N \rightarrow \infty} \frac{1}{N} T(N^\alpha \xi, N^\alpha \eta)$$

for fast-growing sandpiles  $\phi$  is linear inside periodic patches.

**Proof:** Proof as before.



# Directed Triangular lattice showing $\alpha = 1$

Color Code: 0 1 2

N= 3760

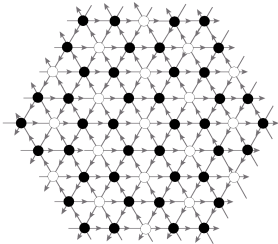


Figure: Periodic  
background: Filled  
circle=1, unfilled=2

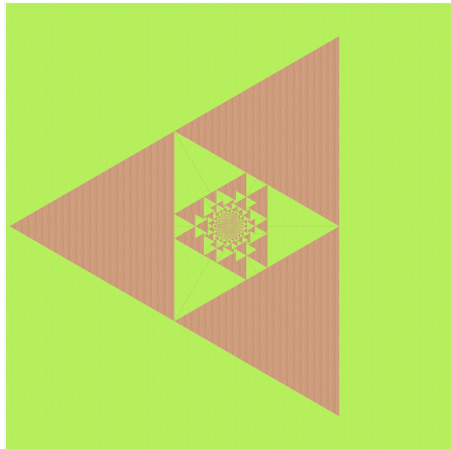
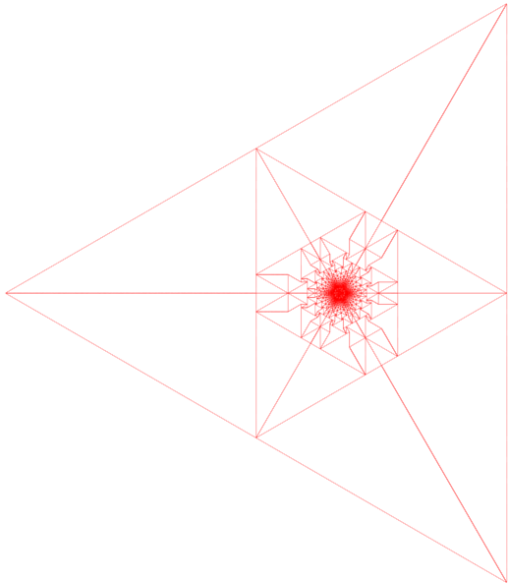
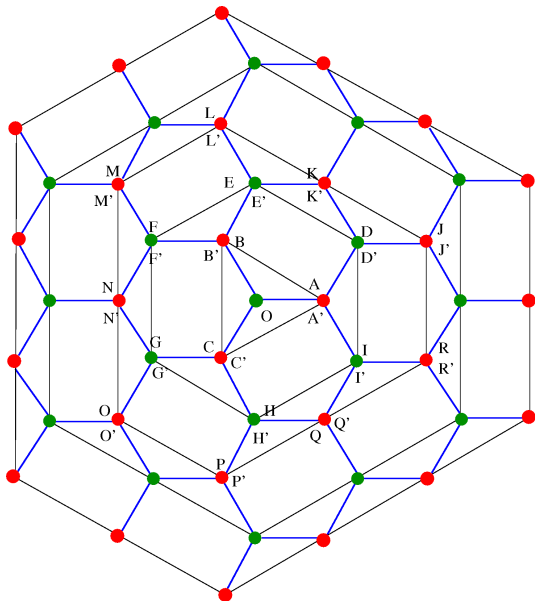


Figure: Pattern with  $N = 1000$



# The adjacency graphs



# Characterizing the triangular non-compact pattern

Analysis is similar to the earlier characterization, actually **simpler**.

The potential function in different patches is given by

$$\phi_P = a_P \xi + b_P \eta + f_P$$

$a_P$  and  $b_P$  are determined by matching slope discontinuity to line charge densities.

Then,  $f_P$  satisfies a Laplace's equation on the adjacency graph.



# Robustness of the pattern

The arguments only depend on the existence of only two types of patches, and straight line boundaries.

These can be found ( by trial and error) in other cases also.  
Then the asymptotic pattern is **identical**.

Some examples:

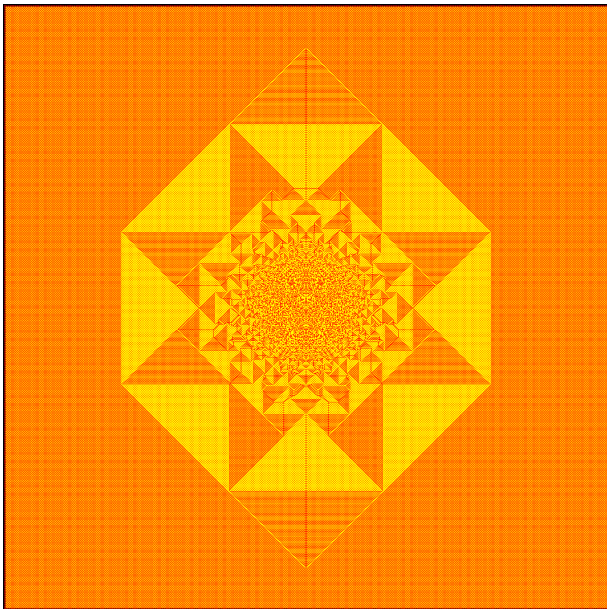


Figure: F-lattice with background density  $5/8$

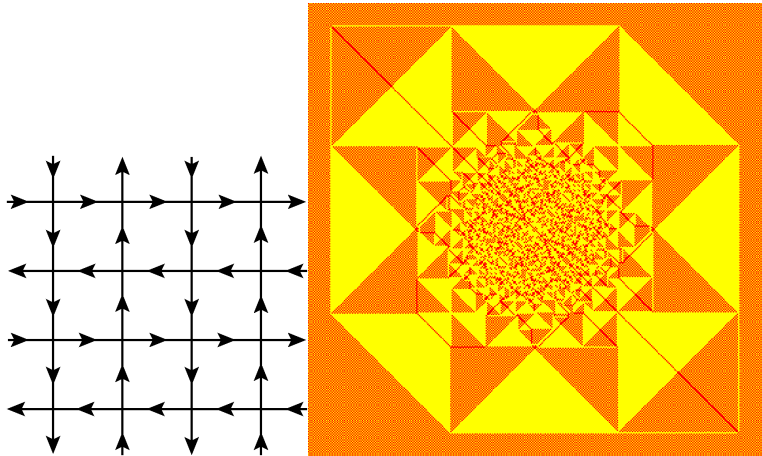


Figure: Manhattan lattice, with initial density  $1/2$ , and 120,000 particles

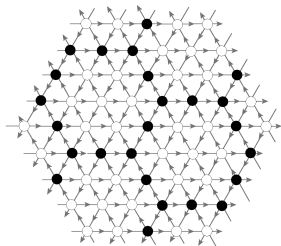


Figure: Periodic  
background: Filled  
circle=1, unfilled=2

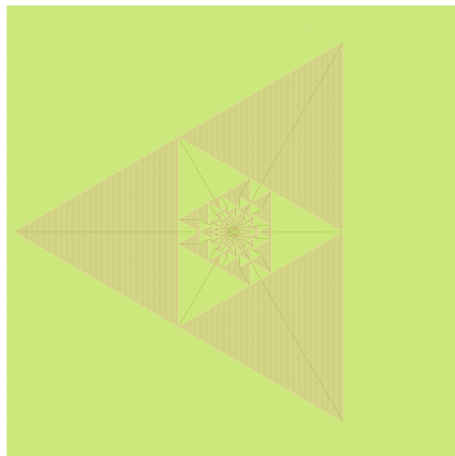


Figure: Pattern with  $N = 1000$

Color Code: 0 1 2

N=500

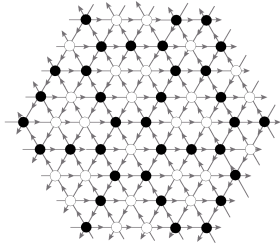


Figure: Periodic  
background: Filled  
circle=1, unfilled=2

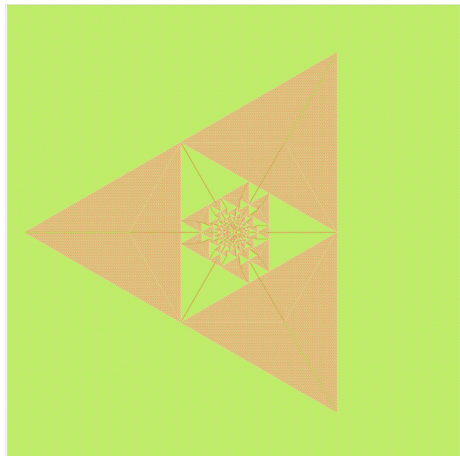


Figure: Pattern with  $N = 1000$

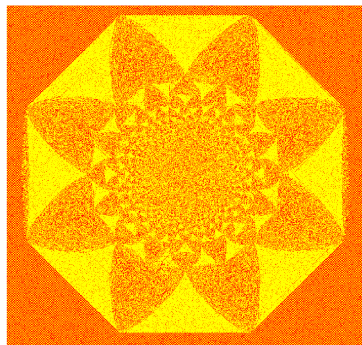
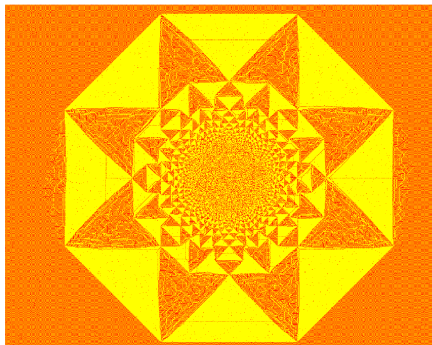


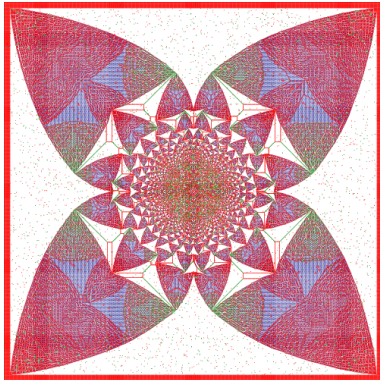
Figure: (a) 1% noise (b) 10%

Noise in the initial particle distribution.

# Noise in the background

Color Code: 0 1 2 3

N=256k



Color Code: 0 1 2 3

N=256k

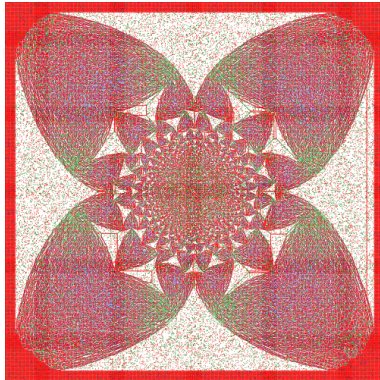


Figure: (a) 1% noise (b) 10%

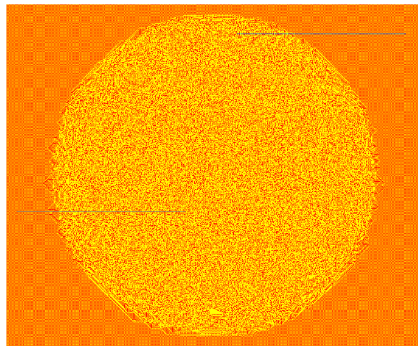
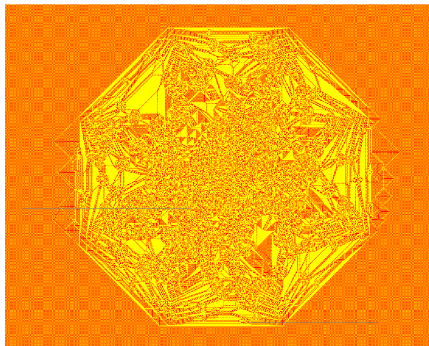


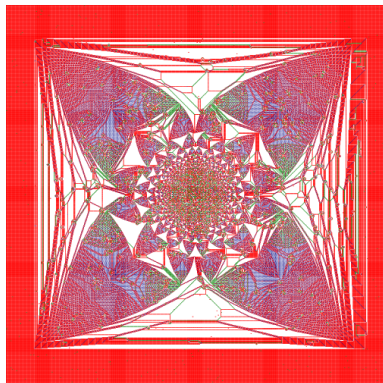
Figure: (a) 0.1% noise (b) 1%

Noise in the relaxation rule.



# Random broken edges

Color Code: 0 1 2 3    N=    150000



Color Code: 0 1 2 3    N=    100000

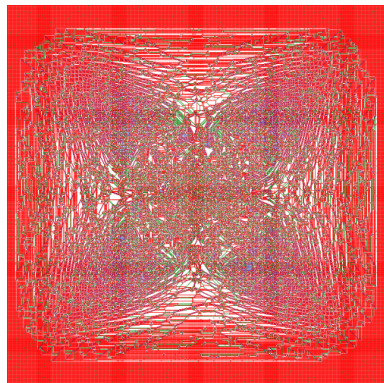


Figure: (a) 1% noise (b) 10%

# A variational formulation

The general principle is called the (lazy man's)

'Least Action Principle': The actual pattern is the stable pattern reached by minimum number of toppling.

$$\nabla^2 \phi = +\delta\rho - \delta(\xi, \eta)$$

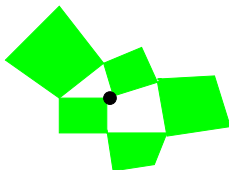
Proof is trivial for abelian models: If a site is unstable, it will not stabilize, until toppled. Order of toppling does not matter.

# Formulation as an electrostatics problem

We have  $\nabla^2\phi = +\delta\rho - \delta(\xi, \eta)$

Positive point charge +1 at origin, and unit negative charge of areal density 1

Can we distribute the negative charge in such a way that the net potential is piecewise-quadratic, and exactly zero far away?

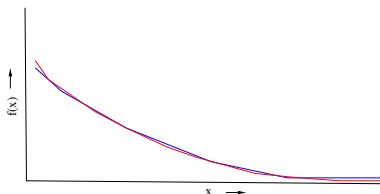


The answer, presumably unique, is the observed pattern on the F-lattice.

Other backgrounds have more choices of charge densities .

# Discrete Quadratic Approximants

Example of discrete approximants:



**Figure:** Approximate  $f(x)$  by piece-wise linear functions with integer slopes

The best “discrete approximant” to a given smooth function.

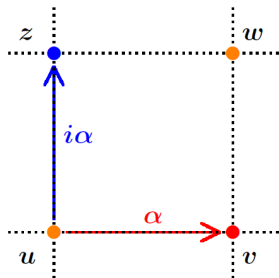
# An iterative formulation

- ▶ Start with a trial pattern.
- ▶ Determine the corresponding  $\phi(\xi, \eta)$
- ▶ Determine the “best” piece-wise quadratic approximants to  $\phi(\xi, \eta)$  using the given set of quadratic functions  $\phi_P$ .
- ▶ The correspond charge density is piece-wise constant. Remove singularities at boundaries.
- ▶ Determine corresponding potential  $\phi^{(1)}(\xi, \eta)$ .
- ▶ Iterate

If the process converges, we get the asymptotic pattern.

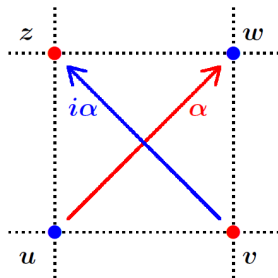
## Discrete Analytic Functions

Functions defined only on discrete points in the complex  $z$ - plane.



$$\frac{F(z) - F(u)}{z - u} = \frac{F(v) - F(u)}{v - u} \quad (1^{st})$$

$$F(z) - F(u) = i(F(v) - F(u))$$



$$\frac{F(z) - F(v)}{z - v} = \frac{F(w) - F(u)}{w - u} \quad (2^{nd})$$

$$F(z) - F(v) = i(F(w) - F(u))$$

simple discrete analytic functions are constant,  $z, z^2, z^3, z^4 - z\bar{z}, \dots$

Define DA function  $F_{1/2}(z)$ , which varies as  $\sqrt{z}$  for large  $|z|$ , and  $F(0) = 0$

The function  $d(m, n) + ie(m, n)$  which characterizes the pattern for F-lattices is  $cF_{1/2}(m + in)$ .

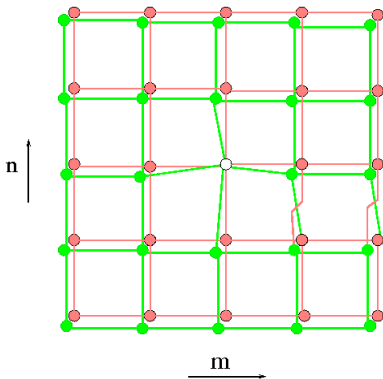


Figure: A discretized two sheeted Riemann surface for  $F_{1/2}(z)$

# Connection to Tropical Mathematics

Define

$$a \oplus b = \text{Max}[a, b]$$

$$a \otimes b = a + b$$

Then standard properties of usual addition and multiplication (commutative, identity, distributive ..) continue to hold.

Example:  $3 \oplus 5 \oplus 2 = 5$

$$3 \otimes 4 = 7$$

Tropical polynomials:  $a \otimes x \otimes x \oplus b \otimes x \oplus c$

Example:  $x \otimes x \oplus 2 \otimes x \oplus 5 = \text{Max}[2x, x + 2, 5]$ .

Fundamental theorem of tropical algebra.

A piecewise -linear convex function can be represented as a tropical polynomial.

Hence useful for describing the function  $\phi(\xi, \eta)$ .



# Summary

- ▶ A model of proportionate growth
- ▶ Quantitatively characterized a large class of patterns with only two types of patches.
- ▶ Additional symmetries.
- ▶ Characterized patterns with multiple sources and sinks. Also determined the growth rates
- ▶ Analyzed a large class of patterns with  $\Lambda > \sqrt{N}$  and quantitatively characterized some such patterns.

# References

- ▶ Pattern formation in growing sandpiles.  
Deepak Dhar, Tridib Sadhu and Samarth chandra, *Euro. Phys. Lett.*,(2008).
- ▶ Pattern formation in growing sandpiles with multiple sources and sinks.  
Tridib Sadhu and Deepak Dhar,*J. Stat. Phys.*,(2009).
- ▶ Pattern formation in fast growing sandpiles.  
In preparation.

Thank you

# Conformal transformation

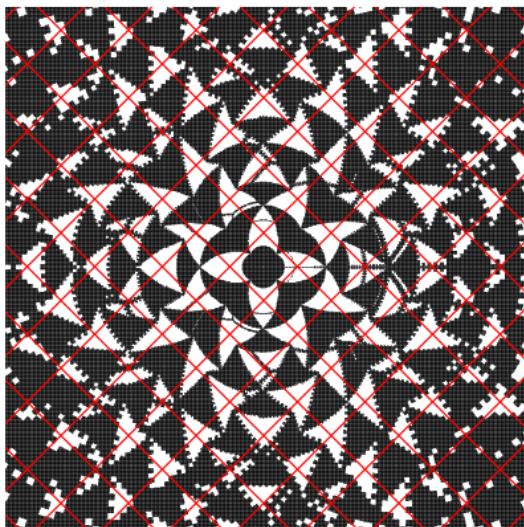


Figure:  $z' = 1/z^2$  picture of the original picture.

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**Proof:** Taylor expand  $\phi(\xi, \eta)$  inside a patch about  $(\xi_0, \eta_0)$ .

$$\phi(\xi_0 + \Delta\xi, \eta_0 + \Delta\eta) = \phi(\xi_0, \eta_0) + d\Delta\xi + e\Delta\eta + a_2\Delta\xi^2 + \dots + K(\Delta\xi)^3 + \dots$$

In terms of toppling number function  $T(X, Y)$  this becomes

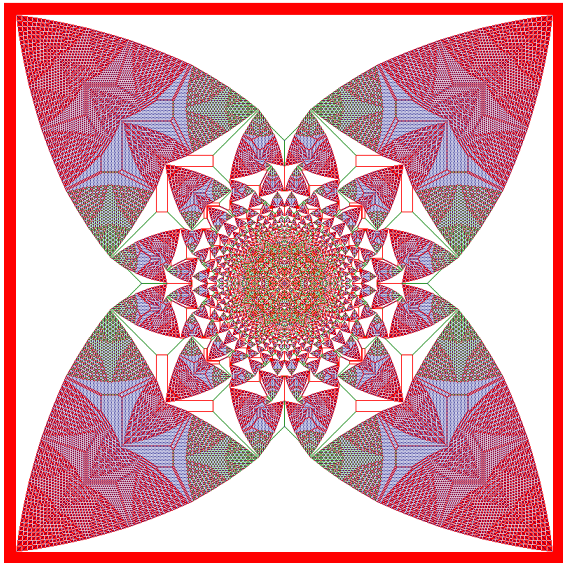
$$T(X_0 + \Delta X, Y_0 + \delta Y) \simeq T(X_0, Y_0) + d\sqrt{N}\Delta X + e\sqrt{N}\Delta Y + a_2\Delta X^2 + \dots + \frac{K(\Delta X)^3}{\sqrt{N}} + \dots$$

Since  $T$  is always an integer, it would jump by 1 at separations  $N^{1/6}$ , causing many defect lines. Hence  $K = 0$ . [← main stream](#)



Color Code: 0 1 2 3

N=250k



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